



---

---

# **Recent progress and future prospects for high energy density experimental physics**

---

---

**Warren W. Hsing**  
**Lawrence Livermore National Laboratory**

**American Physical Society**  
**Division of Plasma Physics**  
**Orlando, FL**  
**Nov 11-15, 2002**

\*This work was performed under the auspices of the U. S. Department of Energy  
by Lawrence Livermore National Laboratory under contract No. W-7405-ENG-48.

# Summary



- 
- **Significant advances in high energy density physics have occurred over the last six years**
  - **The ability to make precise measurements in new regimes allows comparion with models**
    - *Hugoniot equation-of-state*
    - *Materials science at high pressure*
    - *Hydrodynamics*
    - *Radiation transport*
  - **New facilities will expand access to high energy density regimes**

**Regimes of high energy density are typically associated with material energy density  $\geq 1$  MBar**

---



- **Energy density and pressure have the same units**

$$\text{Energy density} = \frac{\text{Energy}}{\text{Volume}}$$

$$\text{Pressure} = \frac{\text{Force}}{\text{Area}} = \frac{\text{Force} \times \text{distance}}{\text{Area} \times \text{distance}} = \frac{\text{Energy}}{\text{Volume}}$$

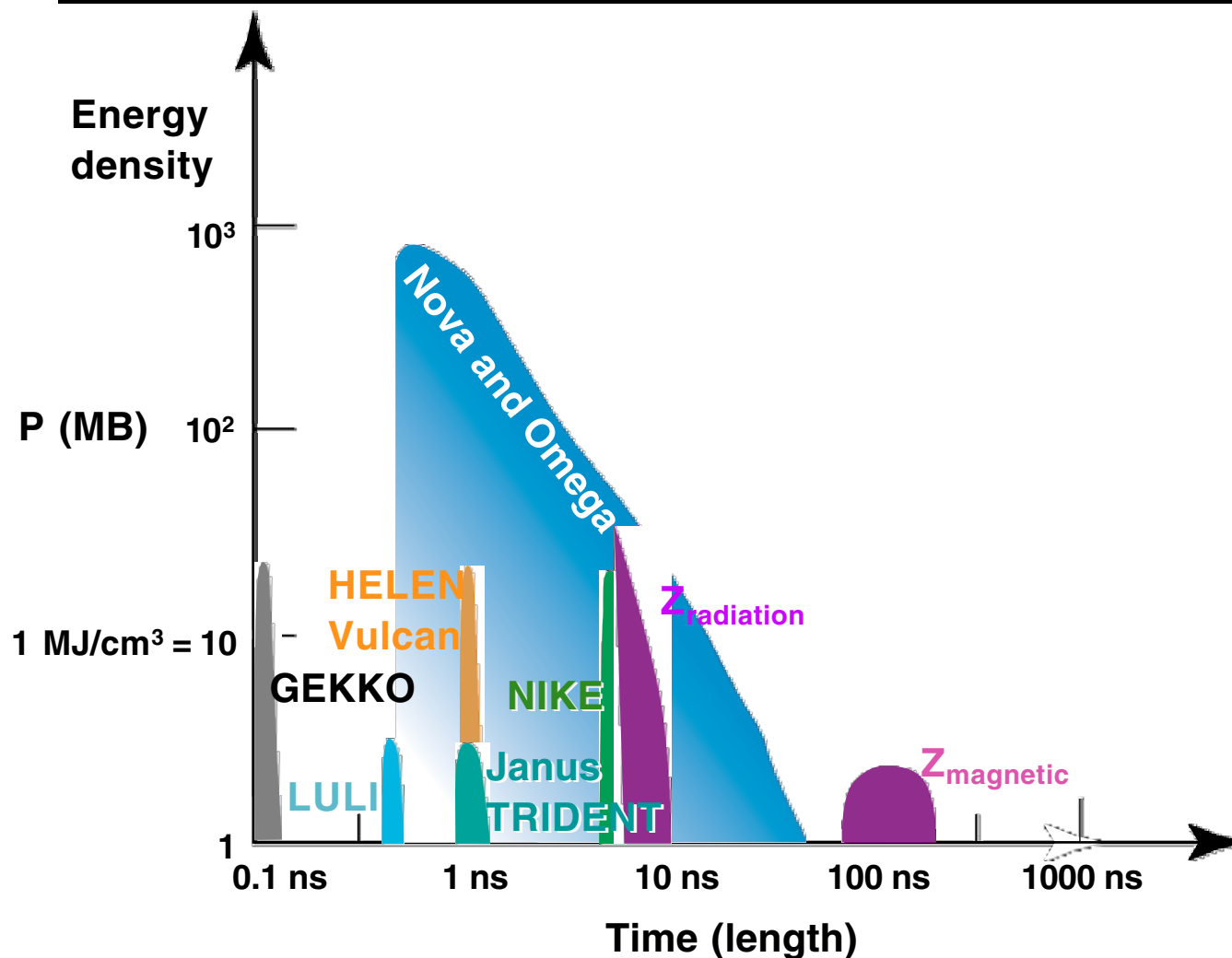
- **$\sim 1$  Mbar is the energy density required to compress material**

$$\text{Energy density} = \frac{\text{Energy}}{\text{Volume}} = \frac{\text{Energy of Bohr atom}}{(\text{Bohr diameter})^3}$$

**$\sim$  few MBar**

**Bulk modulus  $\sim 1/\text{compressibility} \sim 1$  MBars**

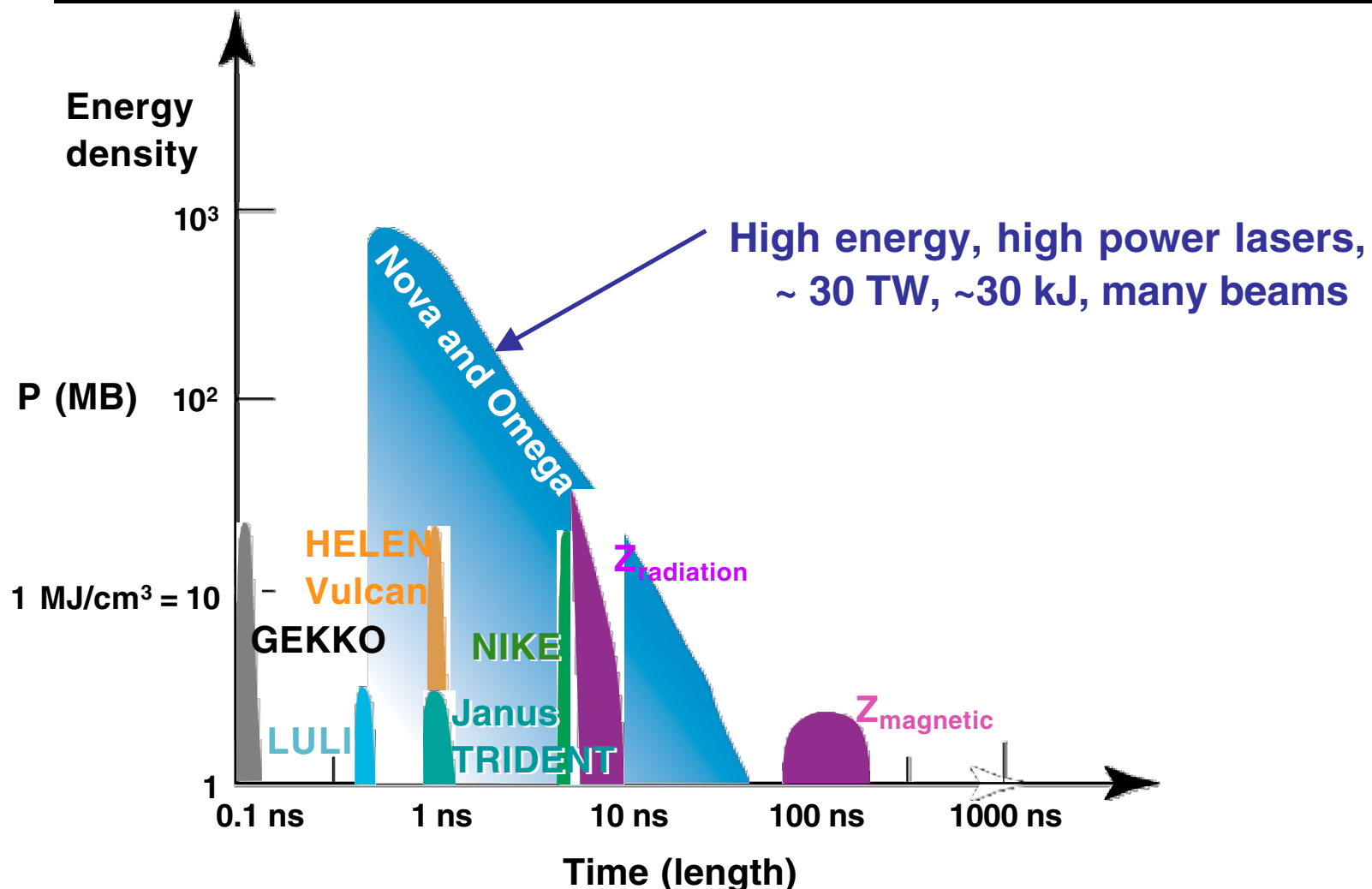
# Laser and pulsed power facilities access high energy density regimes



Most of these facilities were built for ICF  
They can also be used for high energy density physics

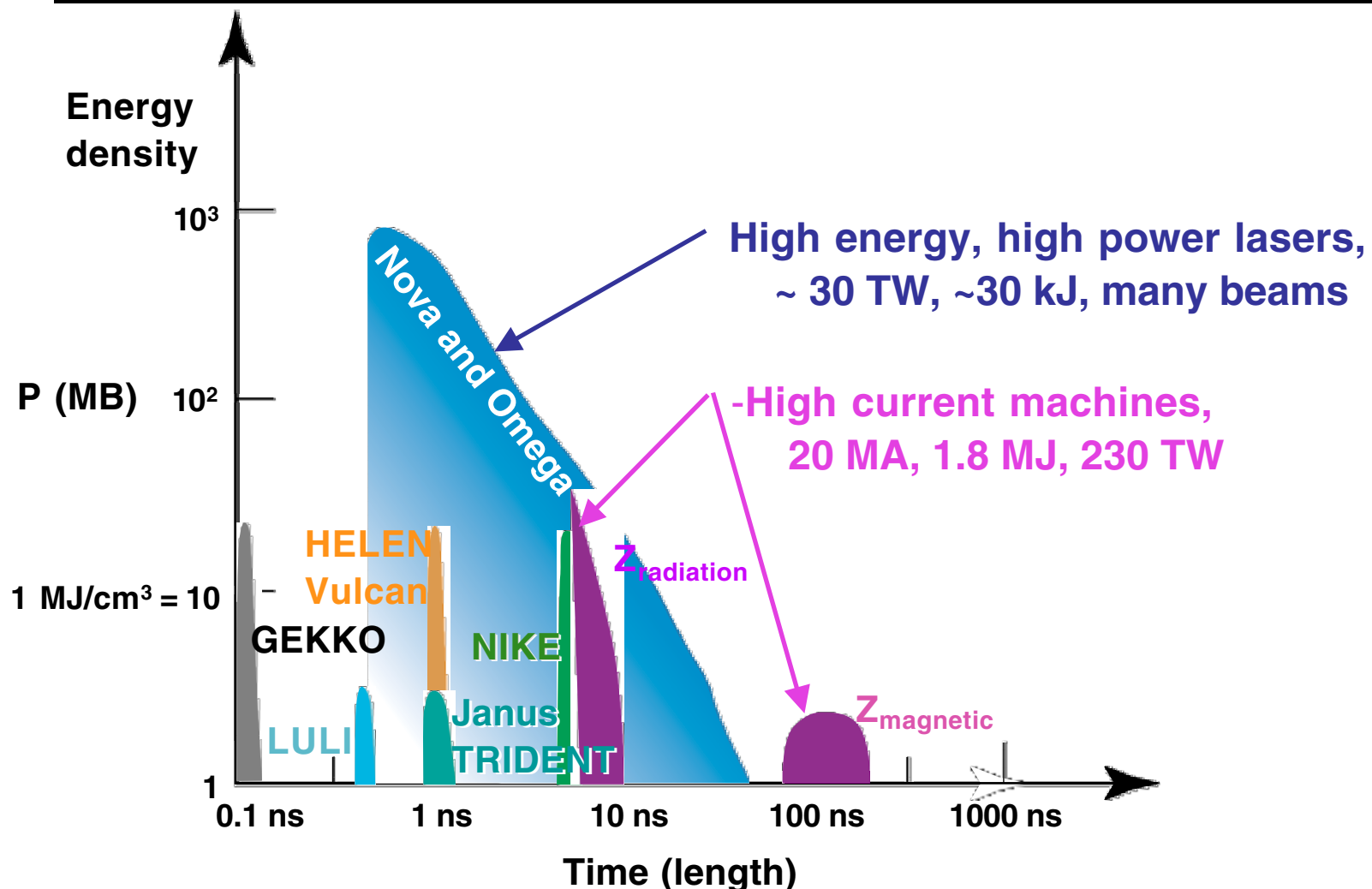


# Laser and pulsed power facilities access high energy density regimes



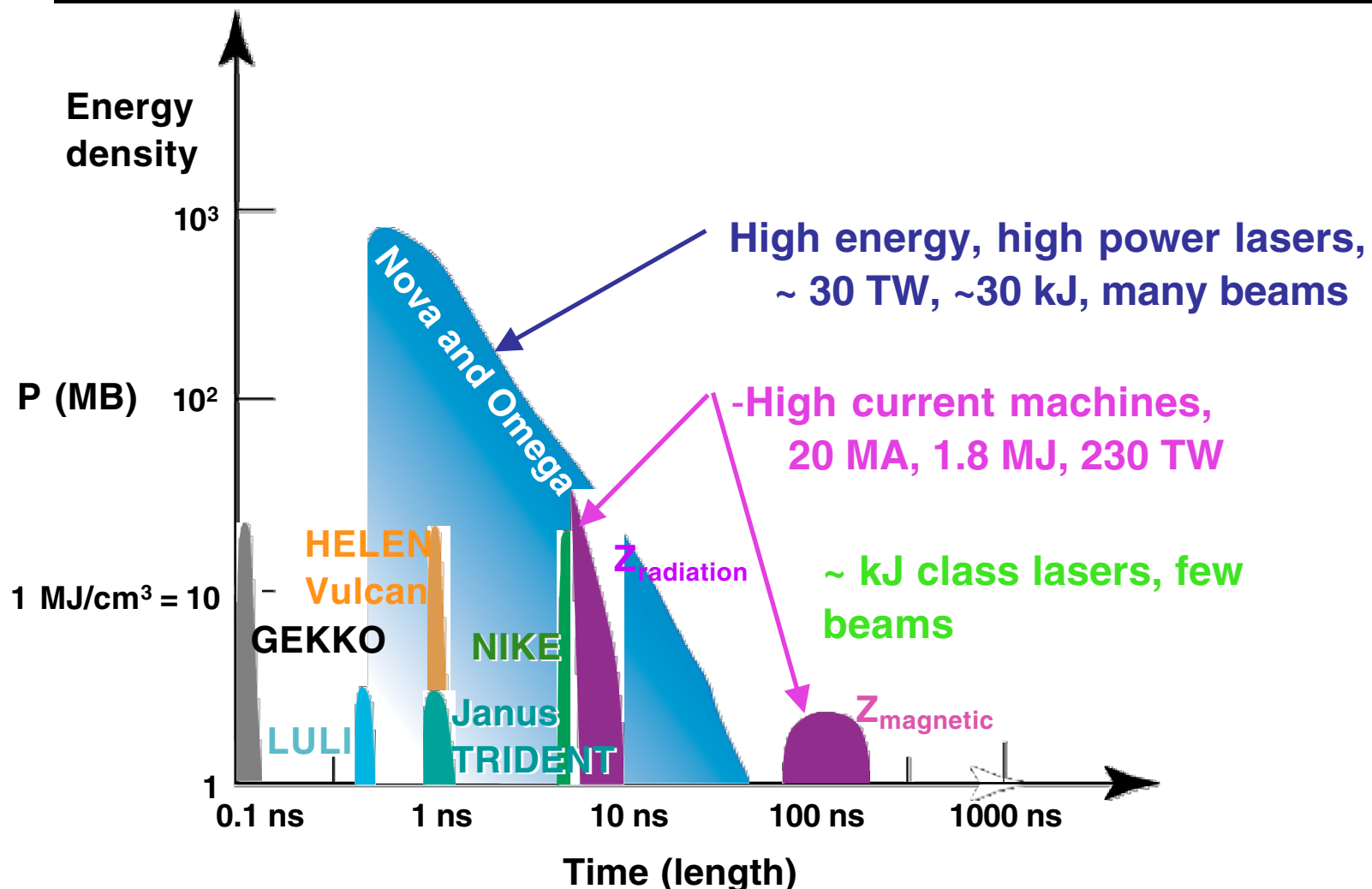
Most of these facilities were built for ICF  
They can also be used for high energy density physics

# Laser and pulsed power facilities access high energy density regimes



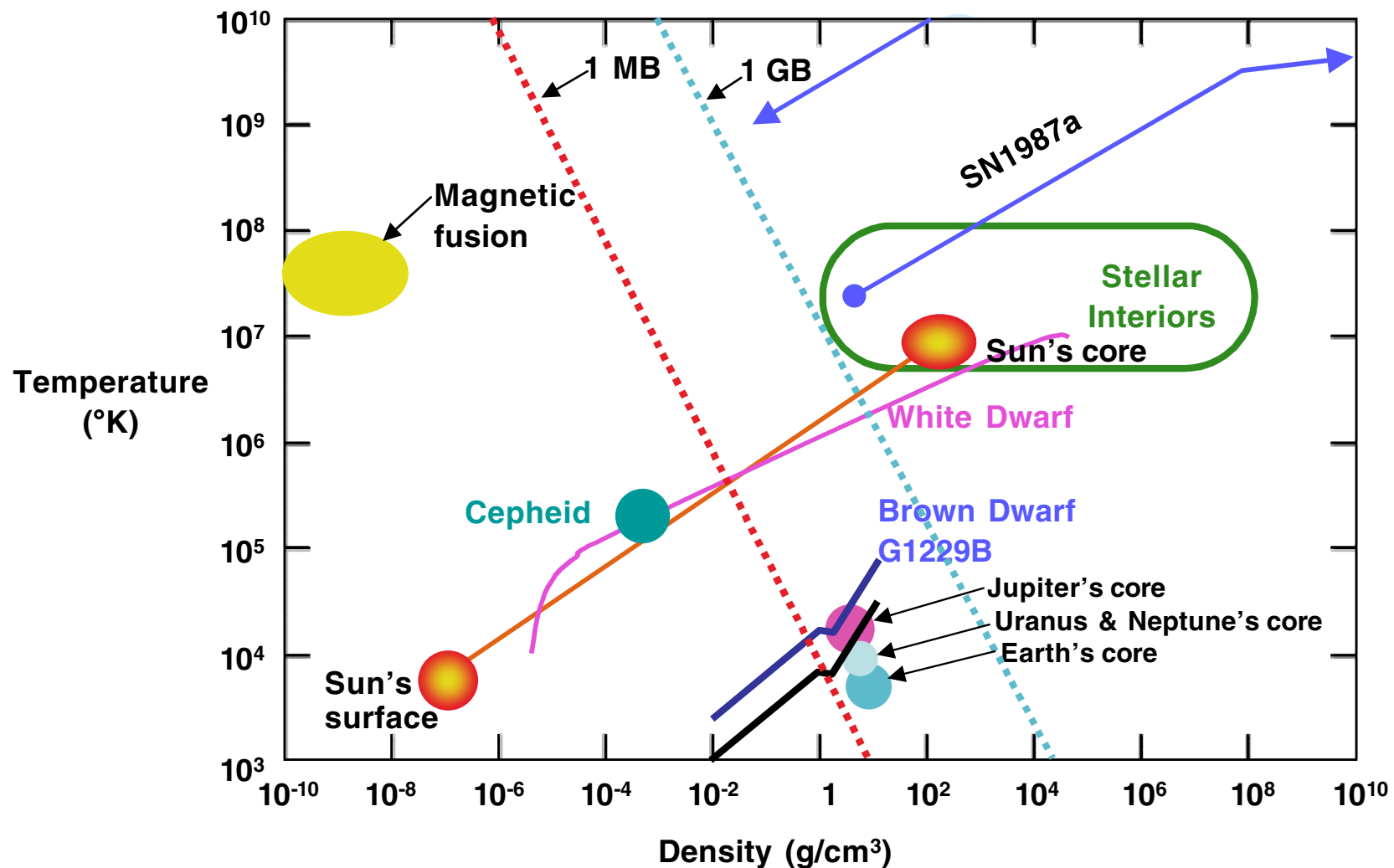
Most of these facilities were built for ICF  
They can also be used for high energy density physics

# Laser and pulsed power facilities access high energy density regimes

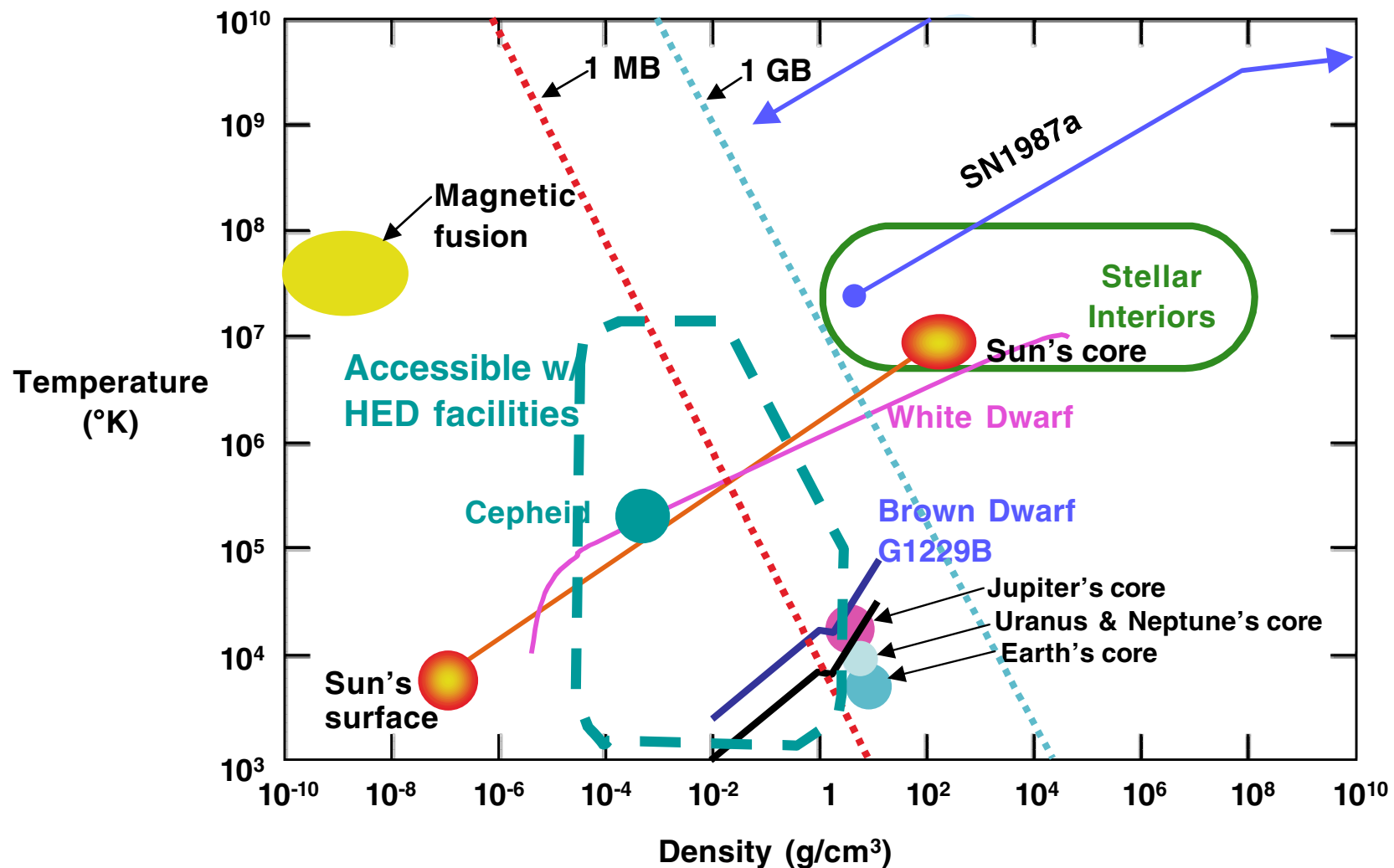


Most of these facilities were built for ICF  
They can also be used for high energy density physics

# High energy density conditions exists in planets and stars



# High energy density conditions exists in planets and stars

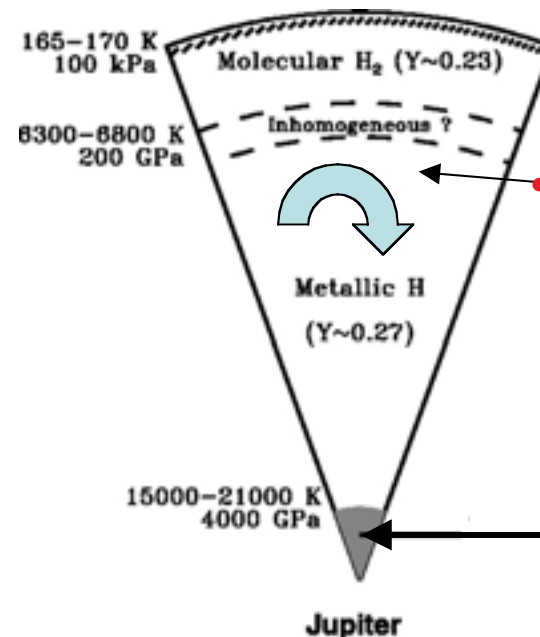
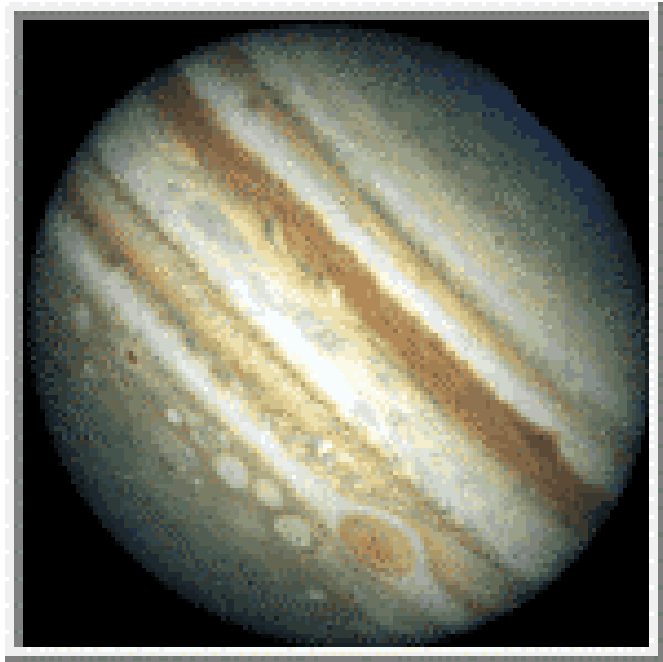


# Outline



- **Significant advances in high energy density physics**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport and atomic physics*
- **Future directions**

# Uncertainties in the equation-of-state of H and He have a large impact on models of the giant planets

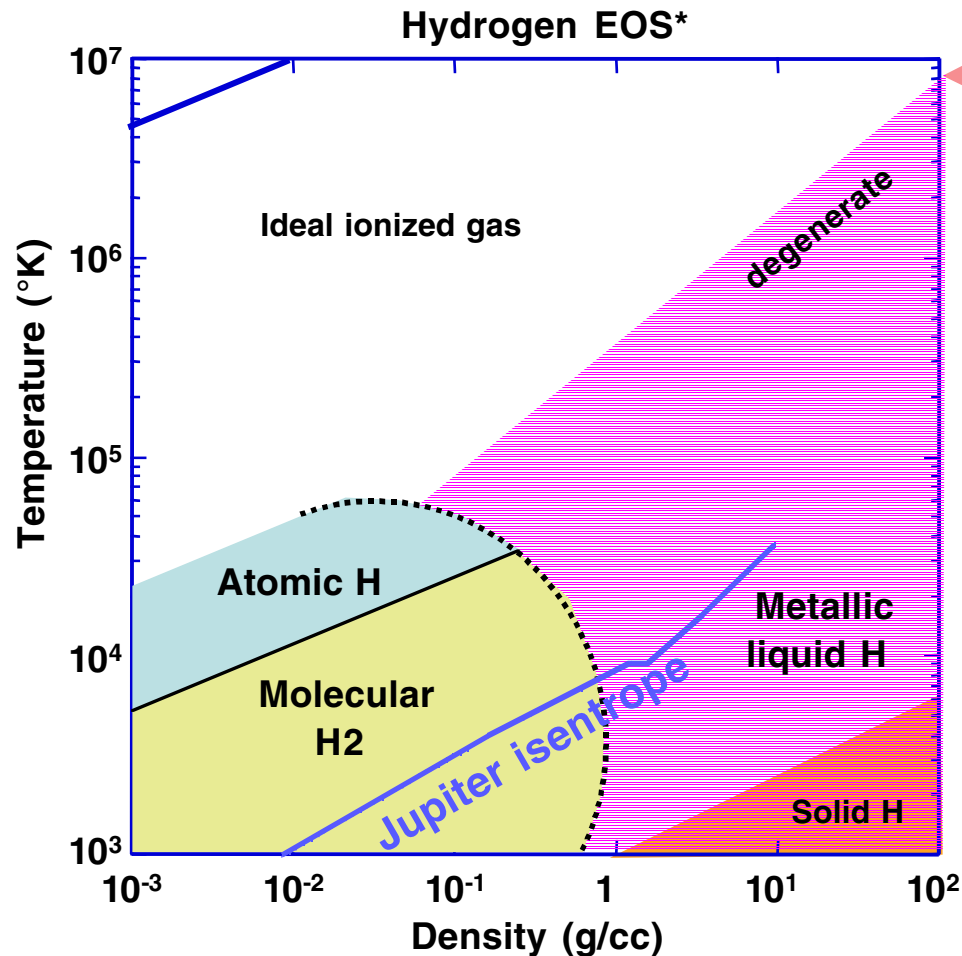


**1st order phase transition to metallic hydrogen?**  
affects thermal evolution, He abundance, magnetic field

**Core : solid or molten?**

- **Models of Jupiter must match a limited set of measurements**
  - *Gravitational field (radius, mass)*
  - *Surface conditions (T, luminosity, spectra)*
- **Hydrogen EOS affects models**
  - *Density structure*
  - *Core*

# Accurate models of hydrogen EOS are difficult in these regimes

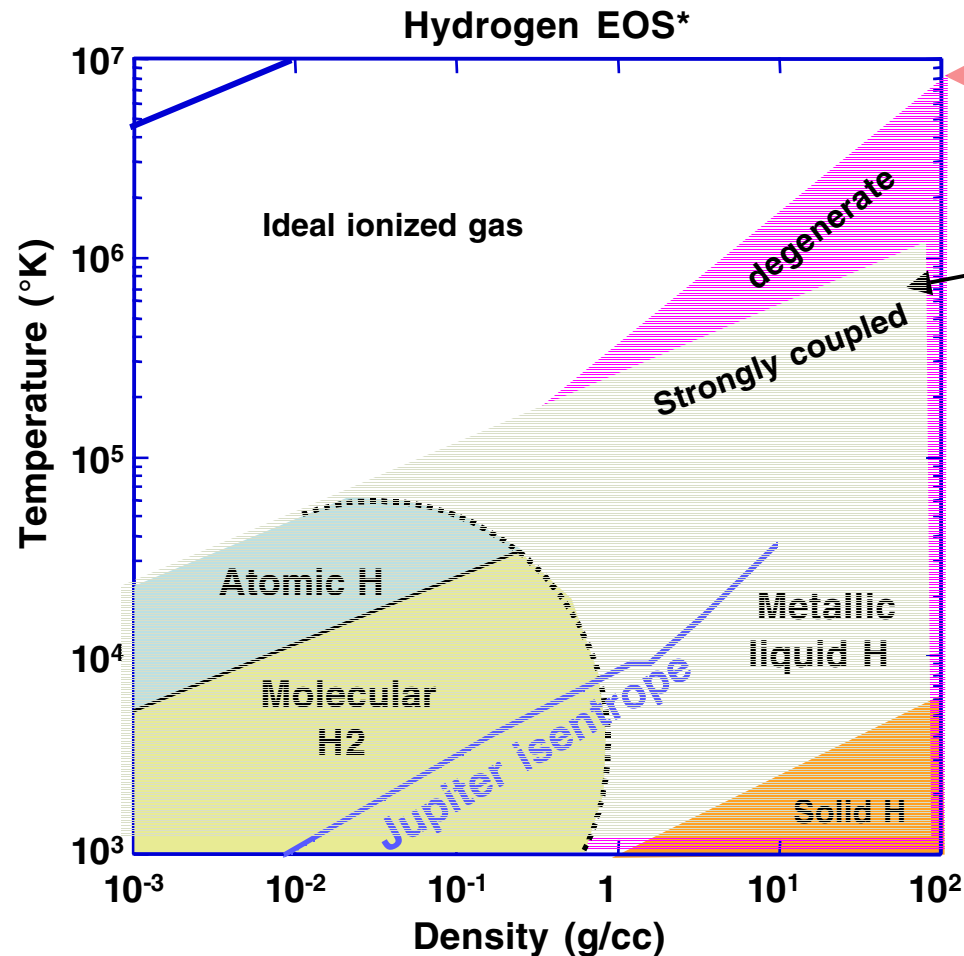


## Degeneracy

- $kT \sim E \text{ (Fermi)} = p^2/2m_e$ ,
- Particle correlations become important



# Accurate models of hydrogen EOS are difficult in these regimes



## Degeneracy

- $kT \sim E \text{ (Fermi)} = p^2/2m_e$ ,
- Particle correlations become

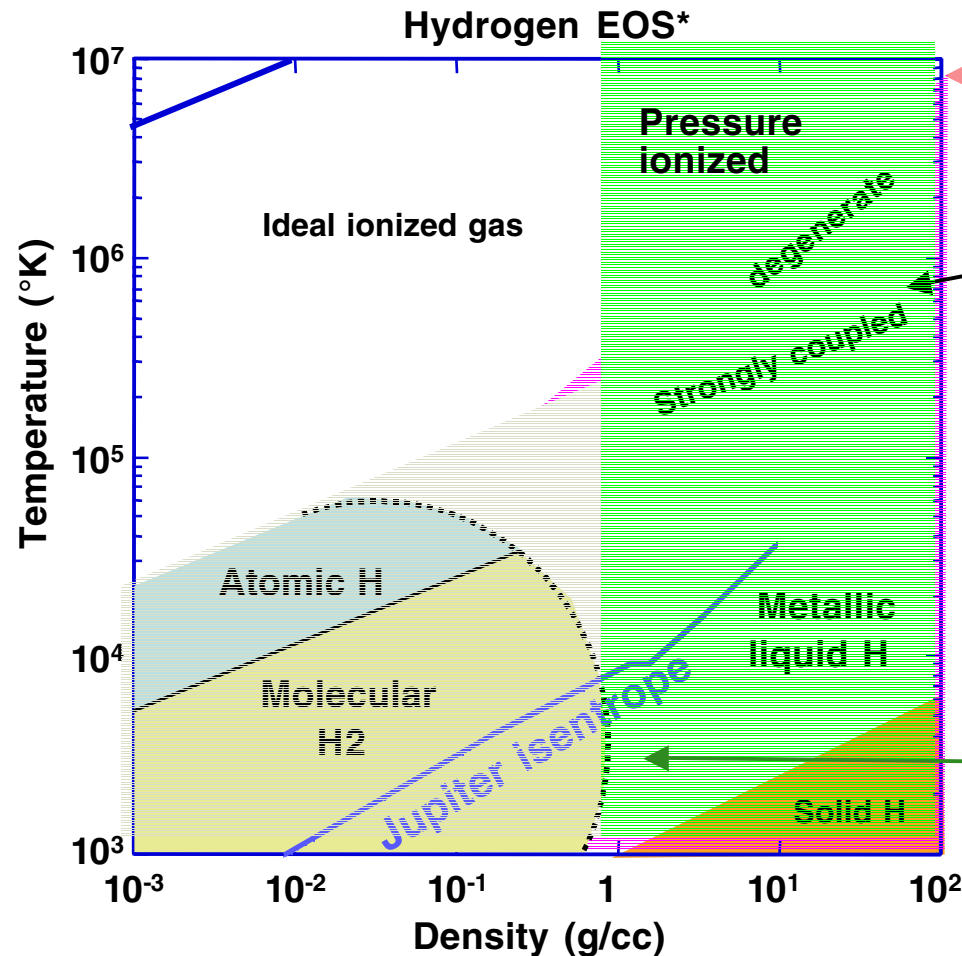
## Strongly coupled plasmas

$$\Gamma \equiv \frac{Z^2 e^2 / a}{kT} = \frac{\text{Coulomb}}{\text{Thermal}}$$

$$\propto Z^{5/6} n_e^{1/3} / T$$

- Long range correlations become important

# Accurate models of hydrogen EOS are difficult in these regimes



## Degeneracy

- $kT \sim E \text{ (Fermi)} = p^2/2m_e$ ,
- Particle correlations become

## Strongly coupled plasmas

$$\Gamma \equiv \frac{Z^2 e^2 / a}{kT} = \frac{\text{Coulomb}}{\text{Thermal}}$$

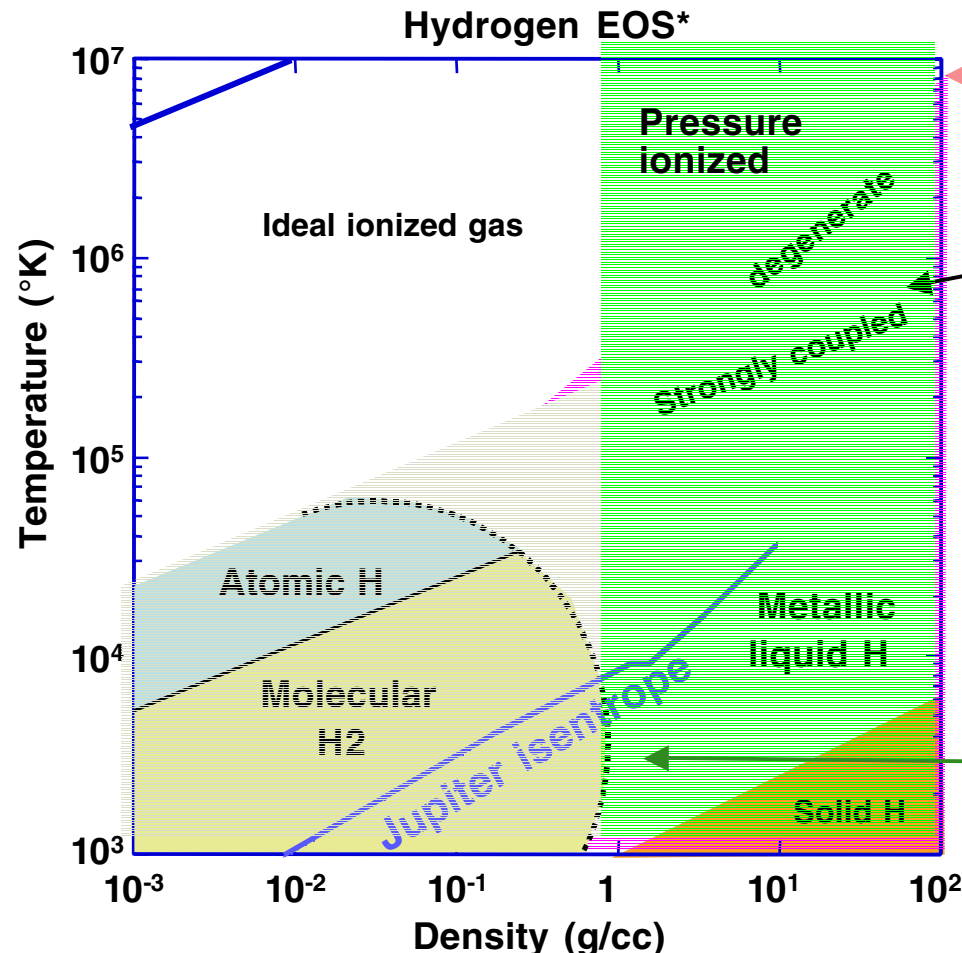
$$\propto Z^{5/6} n_e^{1/3} / T$$

- Long range correlations become important

## Pressure ionization

- Bohr radius  $\sim$  inter-ion spacing
- Ionization potentials are depressed

# Accurate models of hydrogen EOS are difficult in these regimes



## Degeneracy

- $kT \sim E \text{ (Fermi)} = p^2/2m_e$ ,
- Particle correlations become

## Strongly coupled plasmas

$$\Gamma \equiv \frac{Z^2 e^2 / a}{kT} = \frac{\text{Coulomb}}{\text{Thermal}}$$

$$\propto Z^{5/6} n_e^{1/3} / T$$

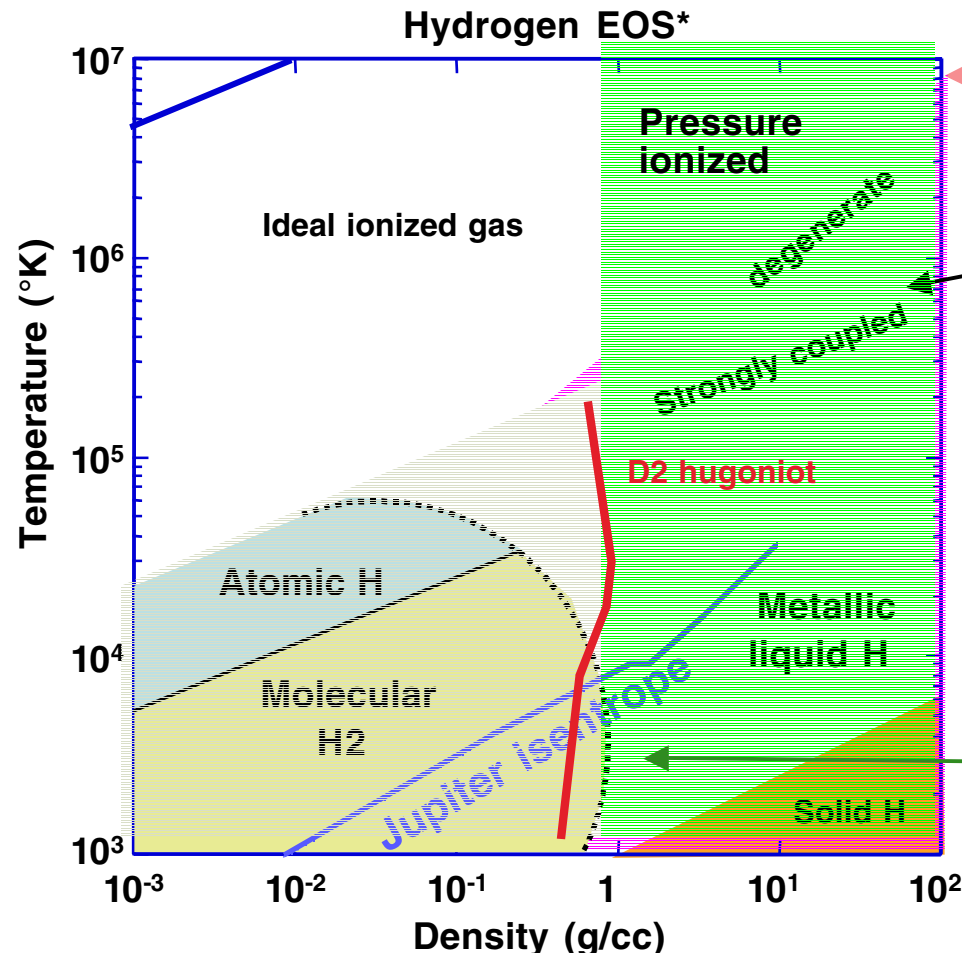
- Long range correlations become important

## Pressure ionization

- Bohr radius  $\sim$  inter-ion spacing
- Ionization potentials are depressed

**Isentrope of Jupiter is strongly coupled and degenerate, and crosses the molecular and pressure ionized regimes**

# Accurate models of hydrogen EOS are difficult in these regimes



## Degeneracy

- $kT \sim E \text{ (Fermi)} = p^2/2m_e$ ,
- Particle correlations become

## Strongly coupled plasmas

$$\Gamma \equiv \frac{Z^2 e^2 / a}{kT} = \frac{\text{Coulomb}}{\text{Thermal}}$$

$$\propto Z^{5/6} n_e^{1/3} / T$$

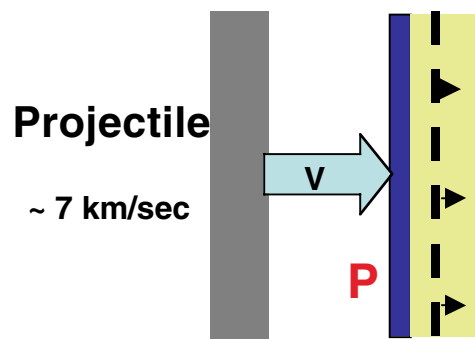
- Long range correlations become important

## Pressure ionization

- Bohr radius  $\sim$  inter-ion spacing
- Ionization potentials are depressed

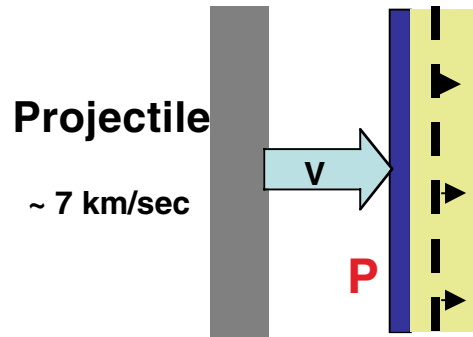
**Isentrope of Jupiter is strongly coupled and degenerate, and crosses the molecular and pressure ionized regimes**

# EOS data for extreme pressures are most accessible using shocks



$P < .2$  Mbars in D2

# EOS data for extreme pressures are most accessible using shocks



$P < .2 \text{ Mbars in D2}$

Across a shock, conservation equations:

Mass:  $\rho_0/\rho = 1 - u_p/U_s$

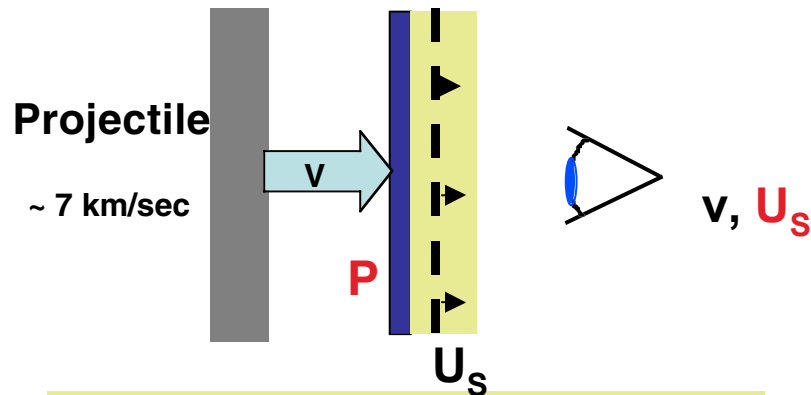
Momentum:  $P - P_0 = \rho_0 U_s u_p$

Energy:  $E - E_0 = .5 (P + P_0)(V_0 - V)$

Three equations in **5 unknowns**

$\Rightarrow$  Velocity is easiest to measure

# EOS data for extreme pressures are most accessible using shocks



Impedance match  
Need reference EOS  $v \rightarrow P$

$P < .2$  Mbars in D2

Across a shock, conservation equations:

Mass:  $\rho_0/\rho = 1 - u_p/U_s$

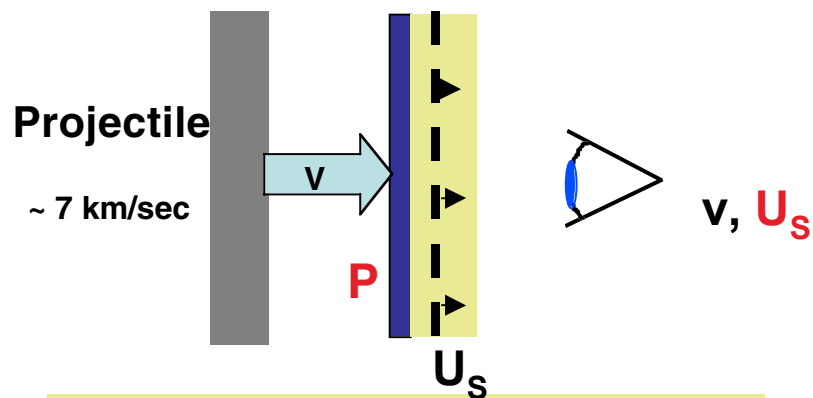
Momentum:  $P - P_0 = \rho_0 U_s u_p$

Energy:  $E - E_0 = .5 (P + P_0)(V_0 - V)$

Three equations in **5 unknowns**

$\Rightarrow$  Velocity is easiest to measure

# EOS data for extreme pressures are most accessible using shocks



Impedance match  
Need reference EOS  $v \rightarrow P$

$P < .2$  Mbars in D2

Across a shock, conservation equations:

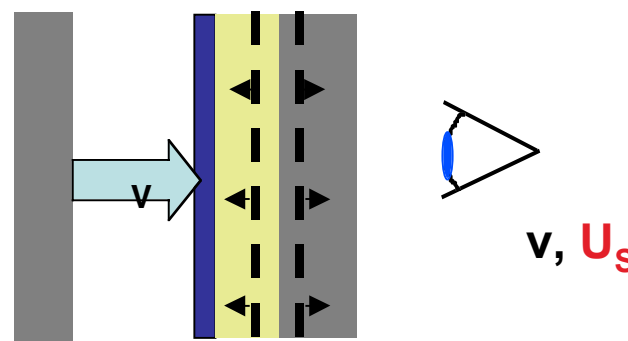
Mass:  $\rho_0/\rho = 1 - u_p/U_s$

Momentum:  $P - P_0 = \rho_0 U_s u_p$

Energy:  $E - E_0 = .5 (P + P_0)(V_0 - V)$

Three equations in **5 unknowns**

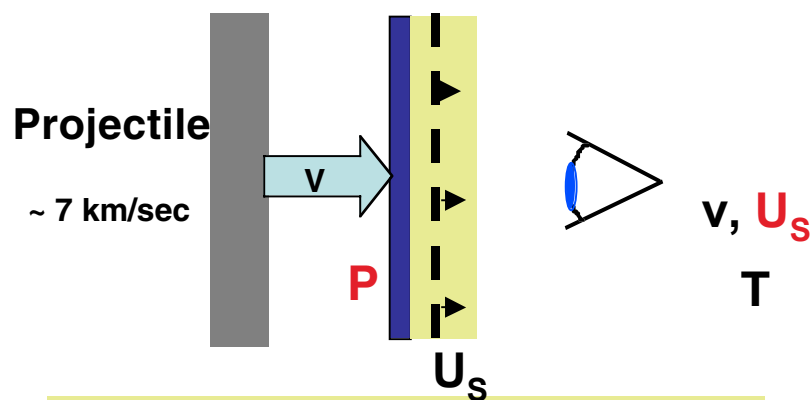
$\Rightarrow$  Velocity is easiest to measure



Double shock allows  $P < .8$  Mbars



# EOS data for extreme pressures are most accessible using shocks



Impedance match  
Need reference EOS  $v \rightarrow P$

$P < .2$  Mbars in D2

Across a shock, conservation equations:

Mass:  $\rho_0/\rho = 1 - u_p/U_s$

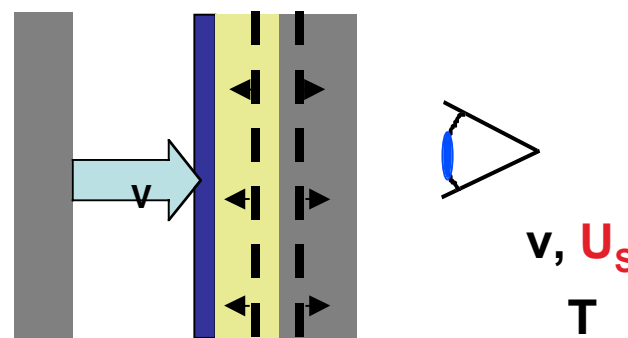
Momentum:  $P - P_0 = \rho_0 U_s u_p$

Energy:  $E - E_0 = .5 (P + P_0)(V_0 - V)$

Three equations in **5 unknowns**

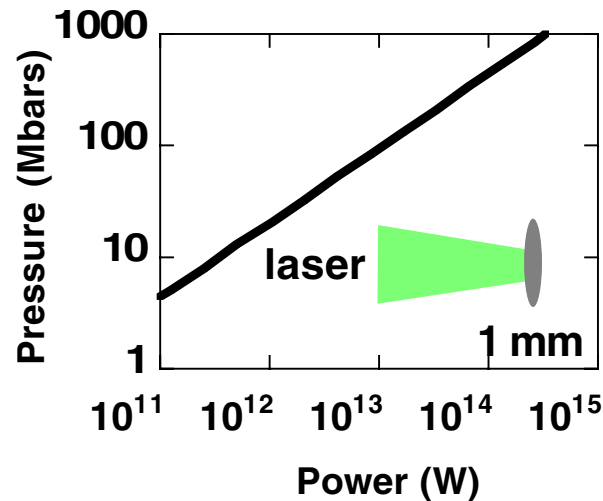
$\Rightarrow$  Velocity is easiest to measure

Measure temperature to get  
additional data for EOS models

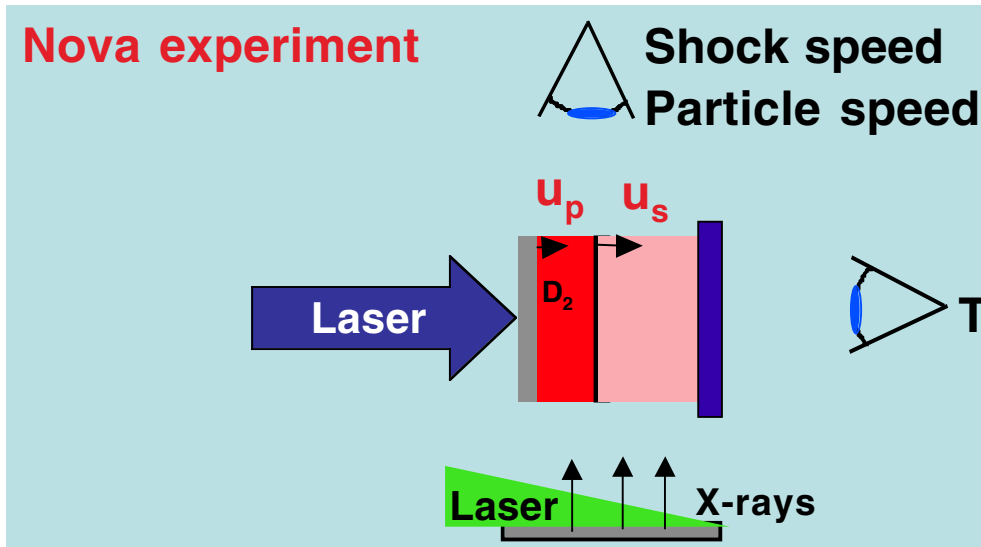


Double shock allows  $P < .8$  Mbars

# Lasers generate high pressure shocks through ablation of material



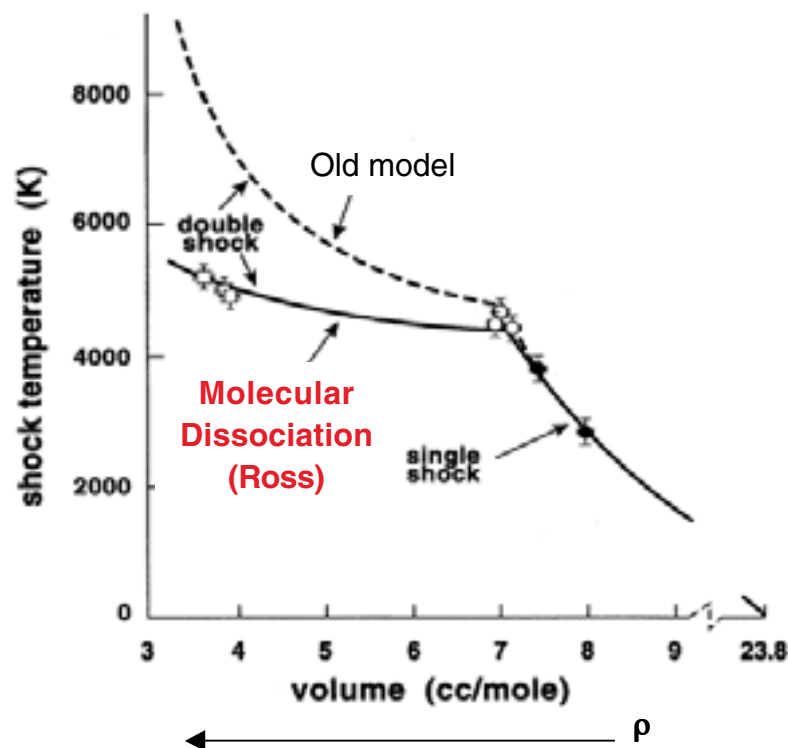
- Potential to get ultra high pressures  $P \sim I^{2/3}$
- Potential to measure 2 quantities directly
  - Particle velocity  $u_p$  and shock velocity  $u_s$
  - *Absolute - no reference material EOS is required*
- Making a precise measurement is difficult



# Shock temperature measurements in deuterium led to a model that predicted a softer EOS



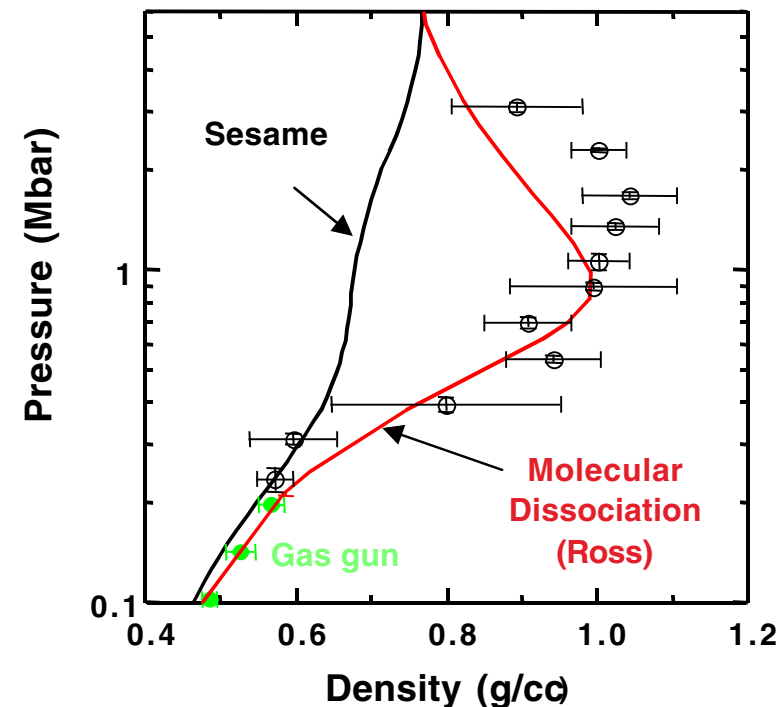
**New dissociation model required explain to gas gun data**



**Implies large compression on Hugoniot**

Holmes, PRB, 52(1995); Ross, J Chem Phys, 79 (1983)

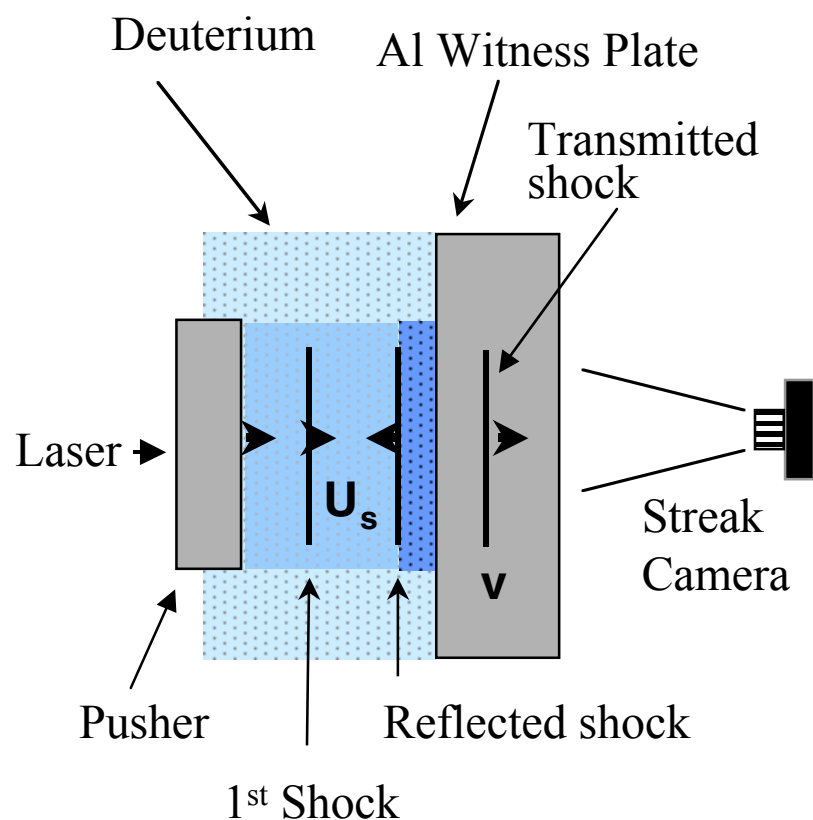
**Deuterium Hugoniot**



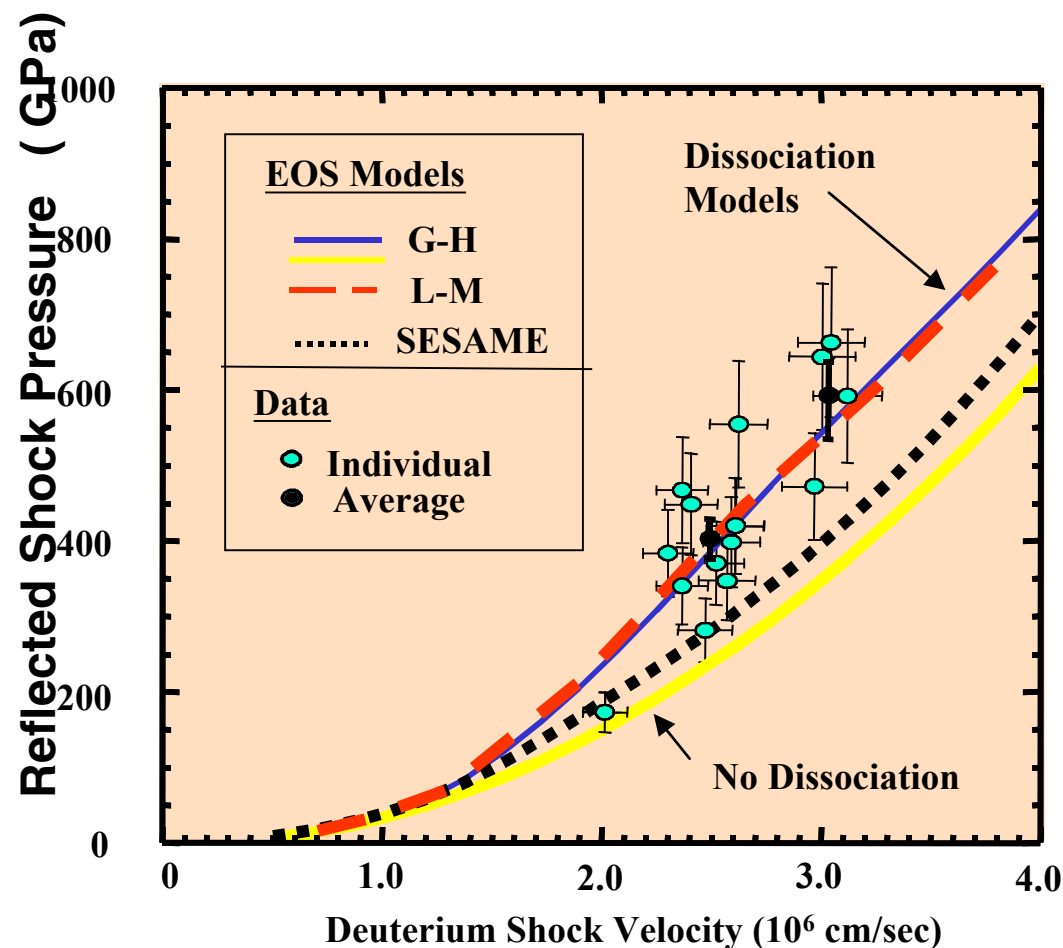
**Nova absolute shock Hugoniot data revealed 50% greater compressibility of hydrogen**

Nellis, J. Chem Phys, 79 (1983); Da Silva, PRL, 78 (1997)

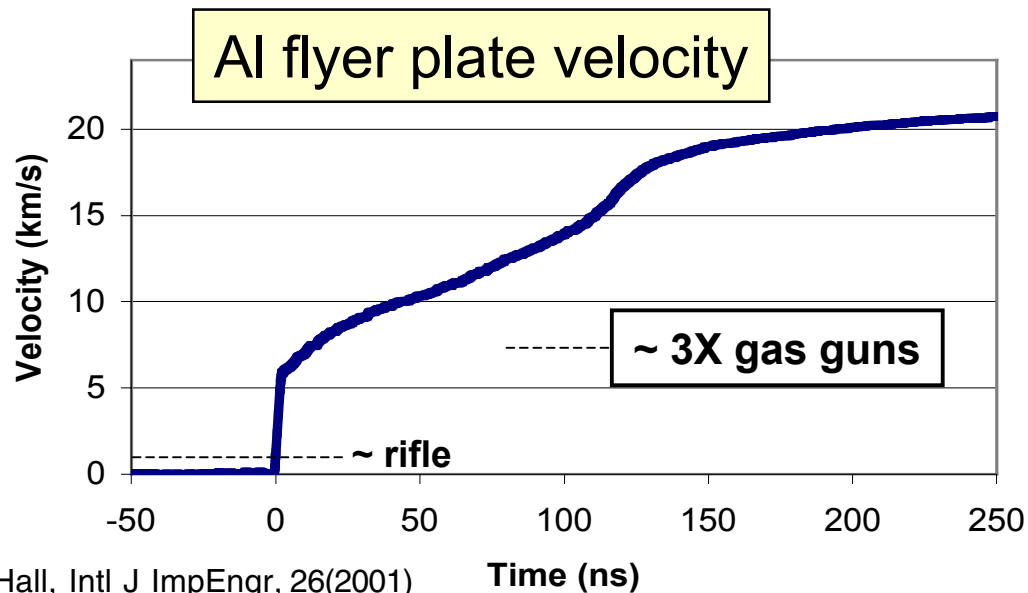
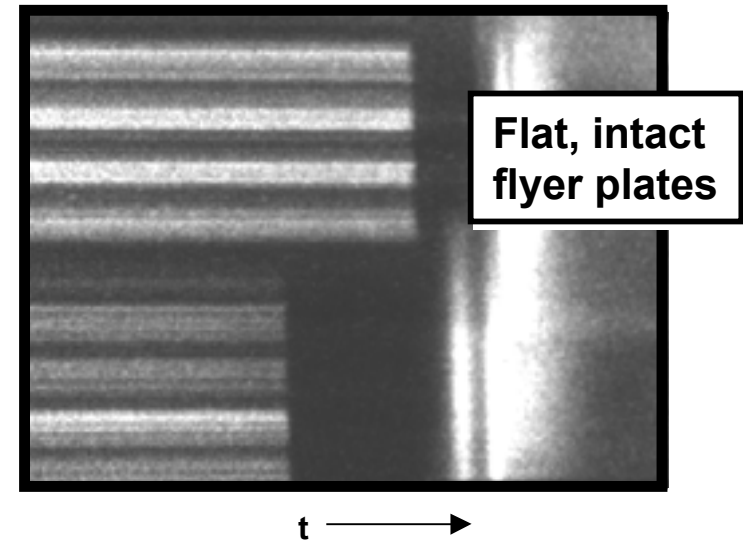
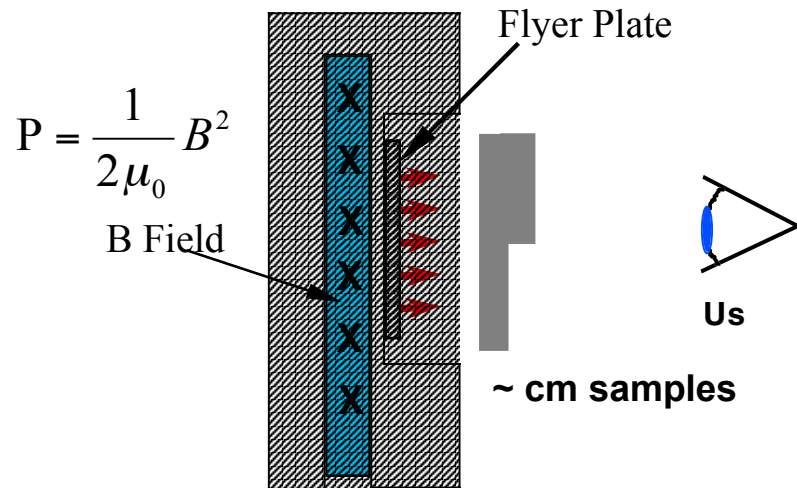
# NRL Reflected Shock Experiments are Consistent with High Compressibility EOS Models



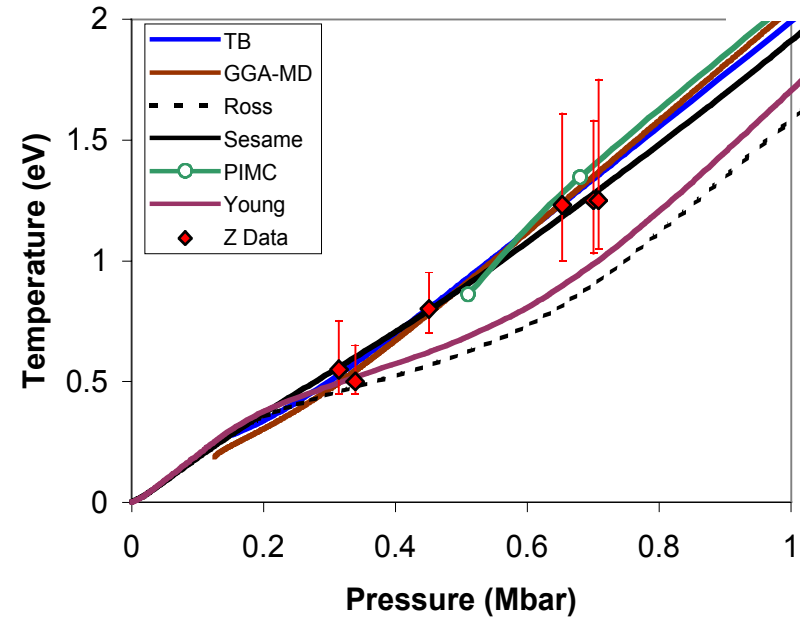
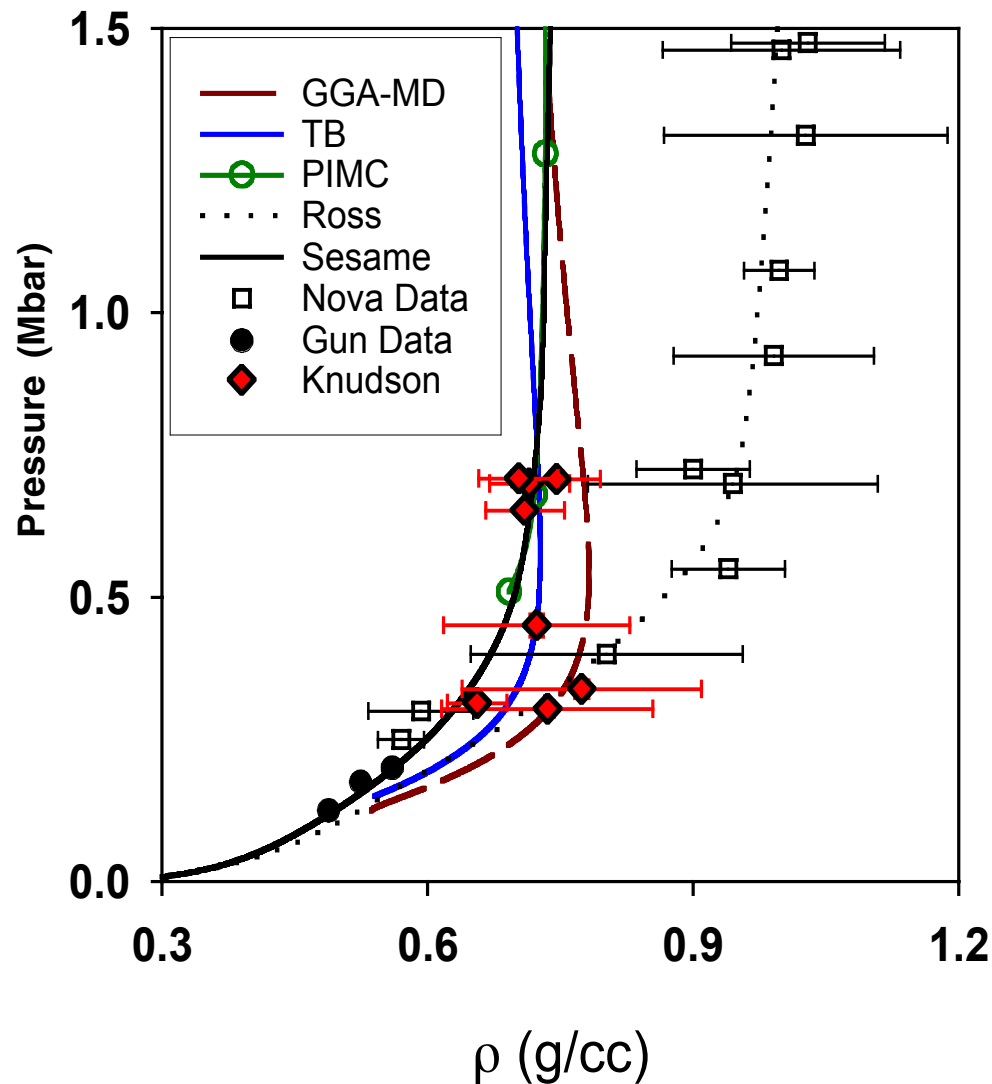
Measure  $U_s$  in D2  
Measure  $v$  in Al  $\rightarrow$  P



# Magnetic driven ultra-high velocity flyer plates were developed for EOS measurements on Z

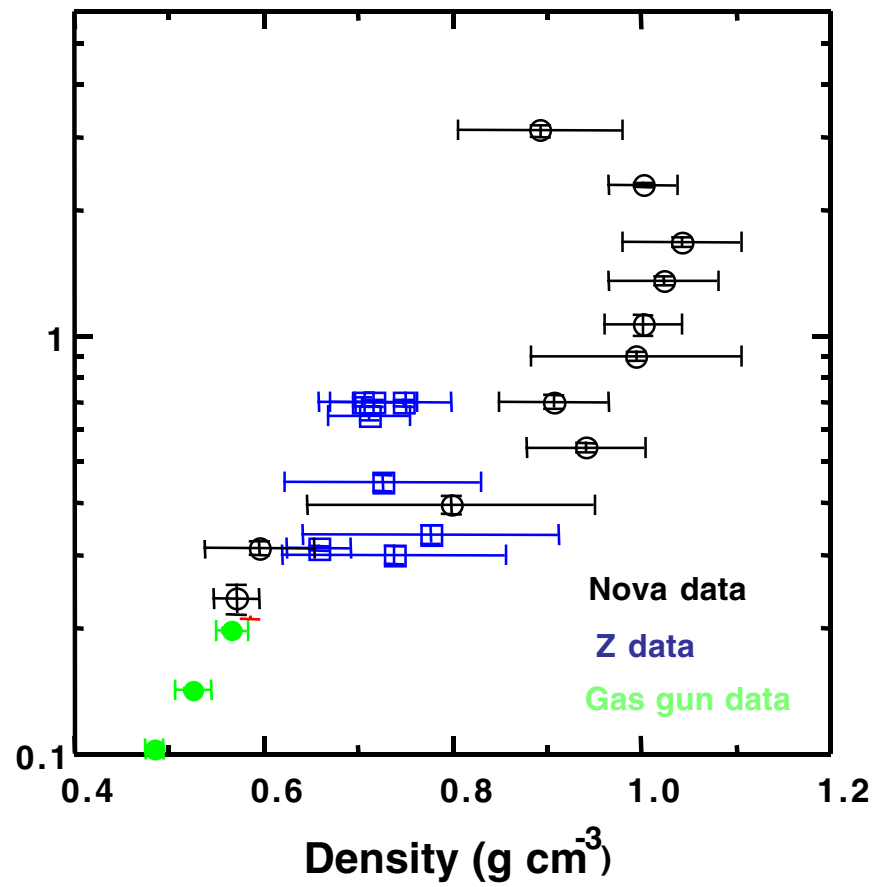


# D<sub>2</sub> EOS data obtained on Z suggest a stiff response in agreement with Sesame and ab-initio models

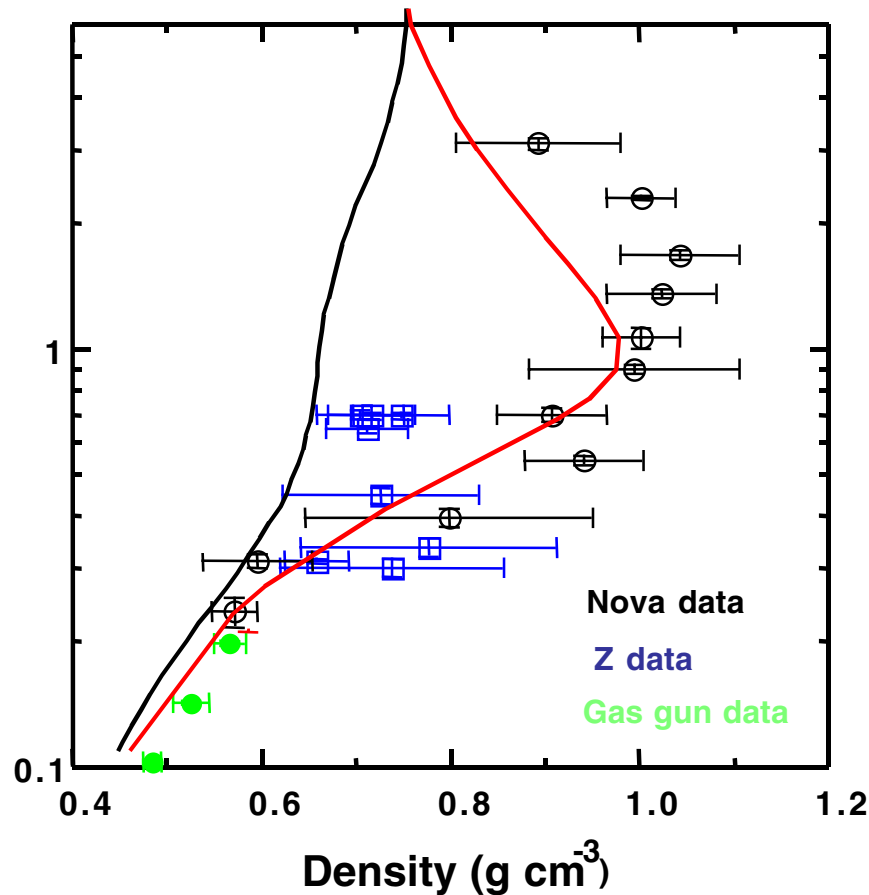


Temperature measurements concured with Sesame

## Models also have large differences in these regimes



# Models also have large differences in these regimes

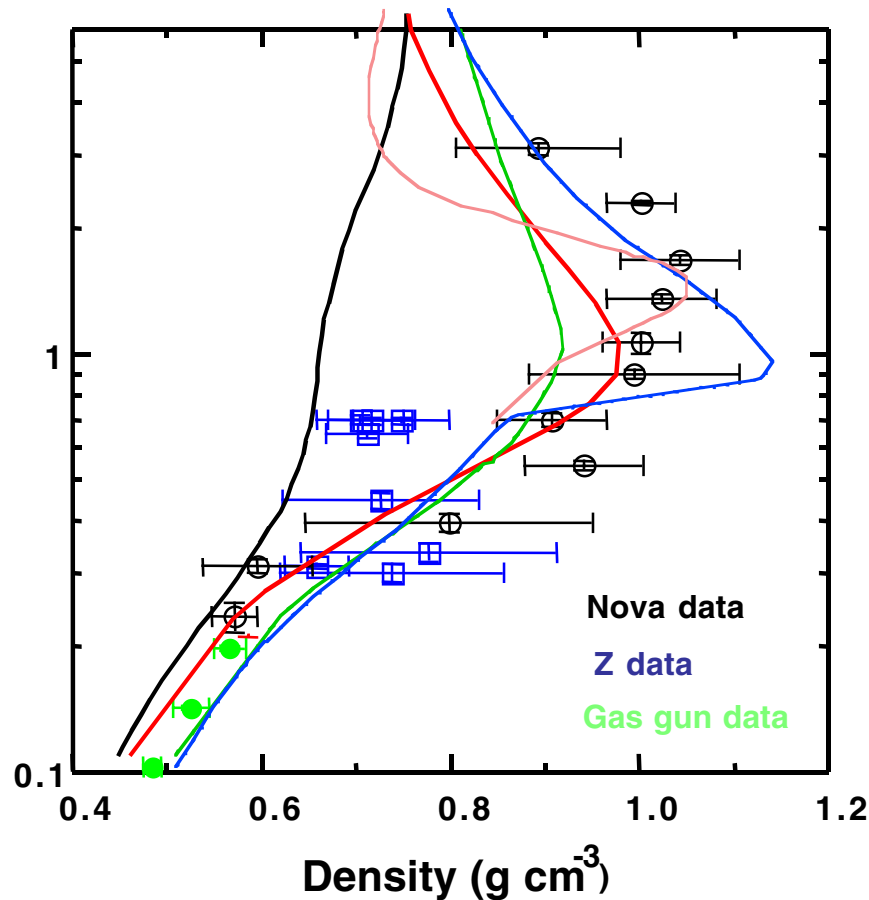


## Statistical mechanics models

— SESAME  
— Dissociation (Ross)



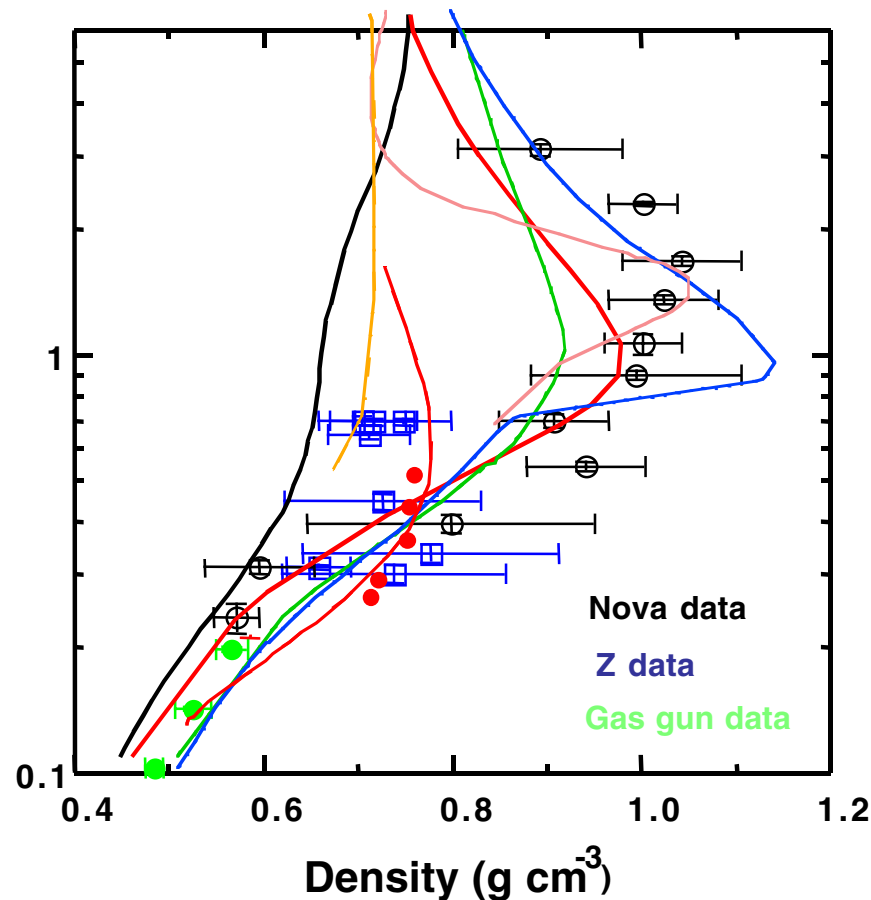
# Models also have large differences in these regimes



## Statistical mechanics models

- SESAME**
- Dissociation (Ross)**
- Free energy: Redmer**
- Free energy: Saumon-Chabrier**
- Activity expansion: Rogers**

# Models also have large differences in these regimes



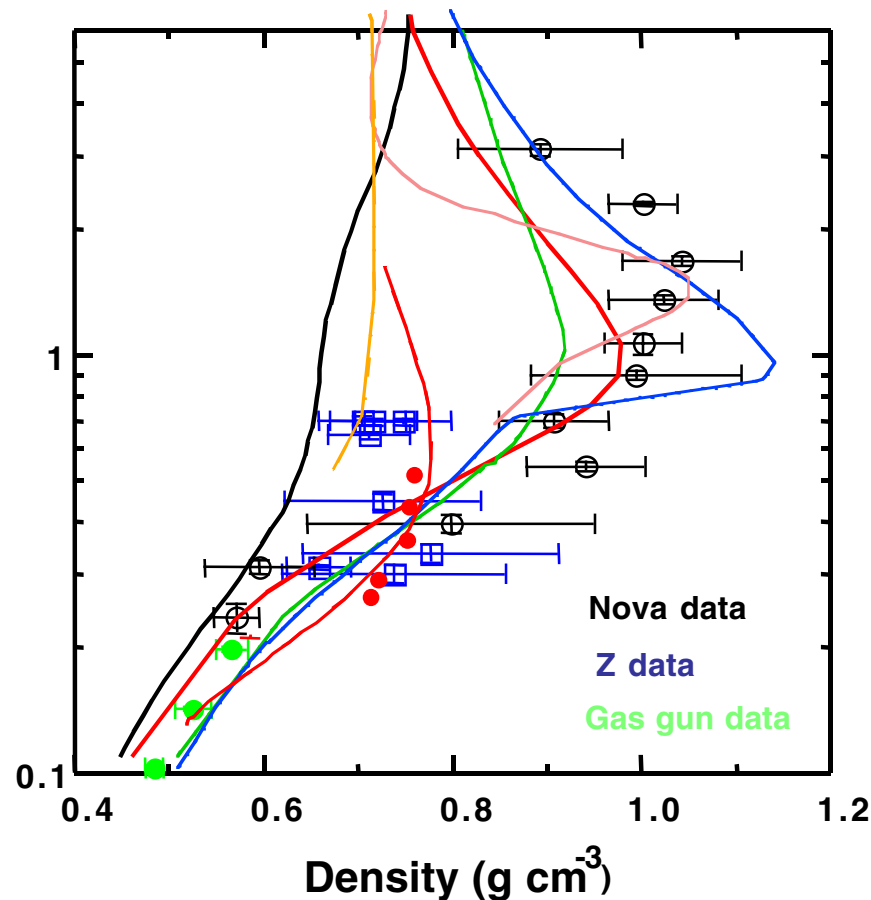
## Statistical mechanics models

- SESAME
- Dissociation (Ross)
- Free energy: Redmer
- Free energy: Saumon-Chabrier
- Activity expansion: Rogers

## Particle simulations

- Path Integral MC: Militzer
- Molecular dynamics: Lenosky
- Molecular dynamics : Galli

# Models also have large differences in these regimes



## Statistical mechanics models

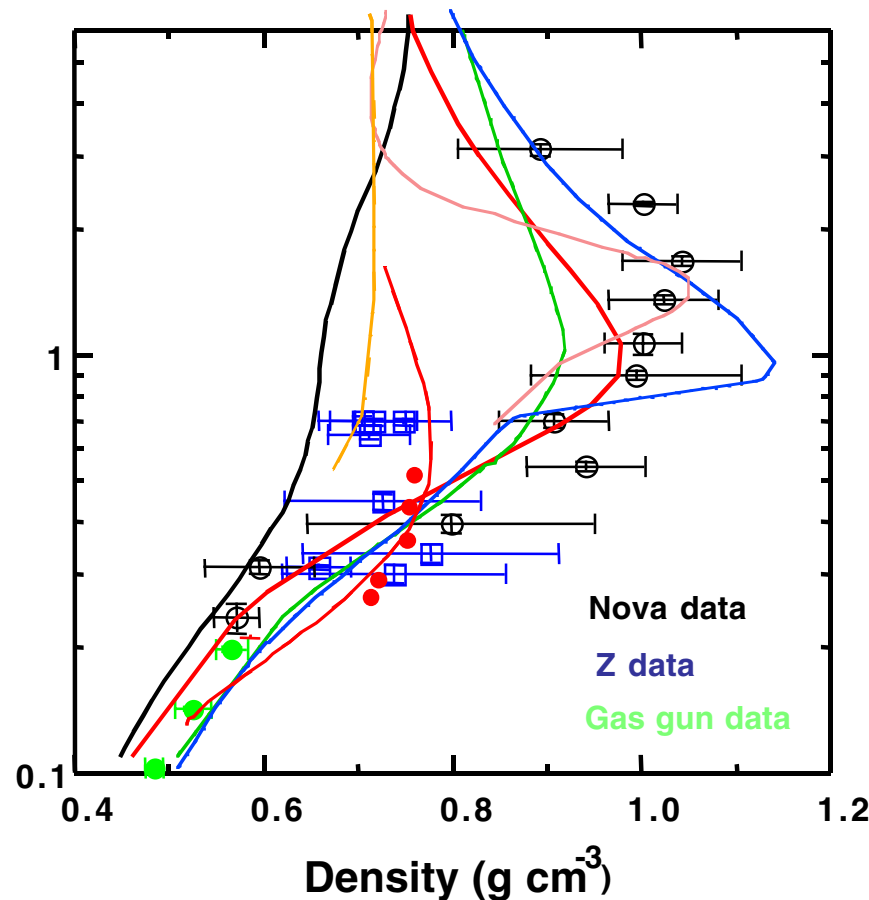
- SESAME
- Dissociation (Ross)
- Free energy: Redmer
- Free energy: Saumon-Chabrier
- Activity expansion: Rogers

## Particle simulations

- Path Integral MC: Militzer
- Molecular dynamics: Lenosky
- Molecular dynamics : Galli

- This is still very much work in progress: scientific method at work!  
*– Stimulated large amount of theoretical work*

# Models also have large differences in these regimes



## Statistical mechanics models

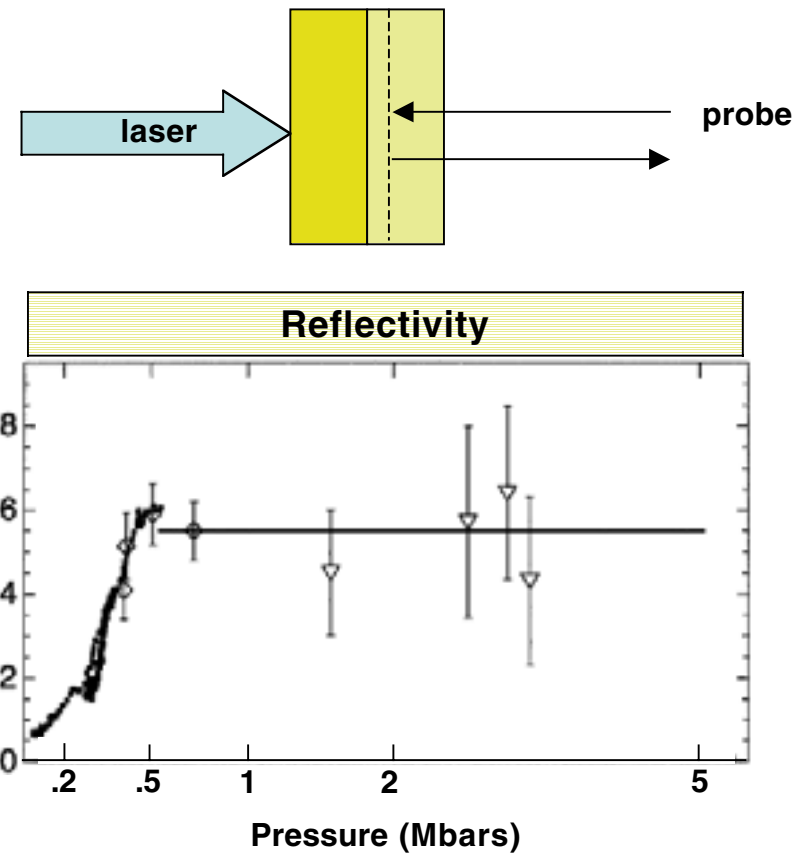
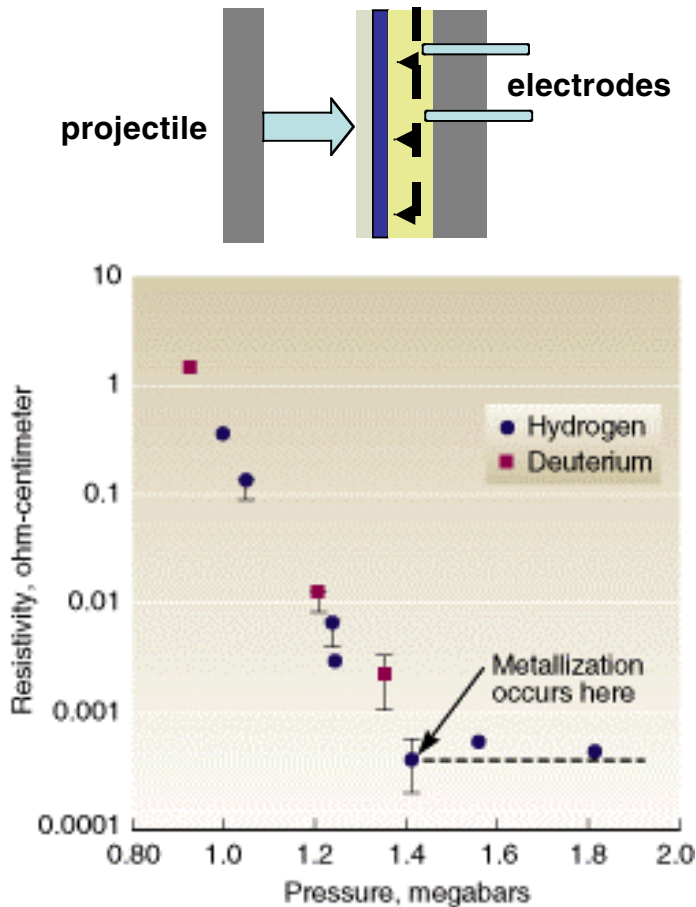
- SESAME
- Dissociation (Ross)
- Free energy: Redmer
- Free energy: Saumon-Chabrier
- Activity expansion: Rogers

## Particle simulations

- Path Integral MC: Militzer
- Molecular dynamics: Lenosky
- Molecular dynamics : Galli

- This is still very much work in progress: scientific method at work!
  - *Stimulated large amount of theoretical work*
- Only way to get high pressure Hugoniot EOS

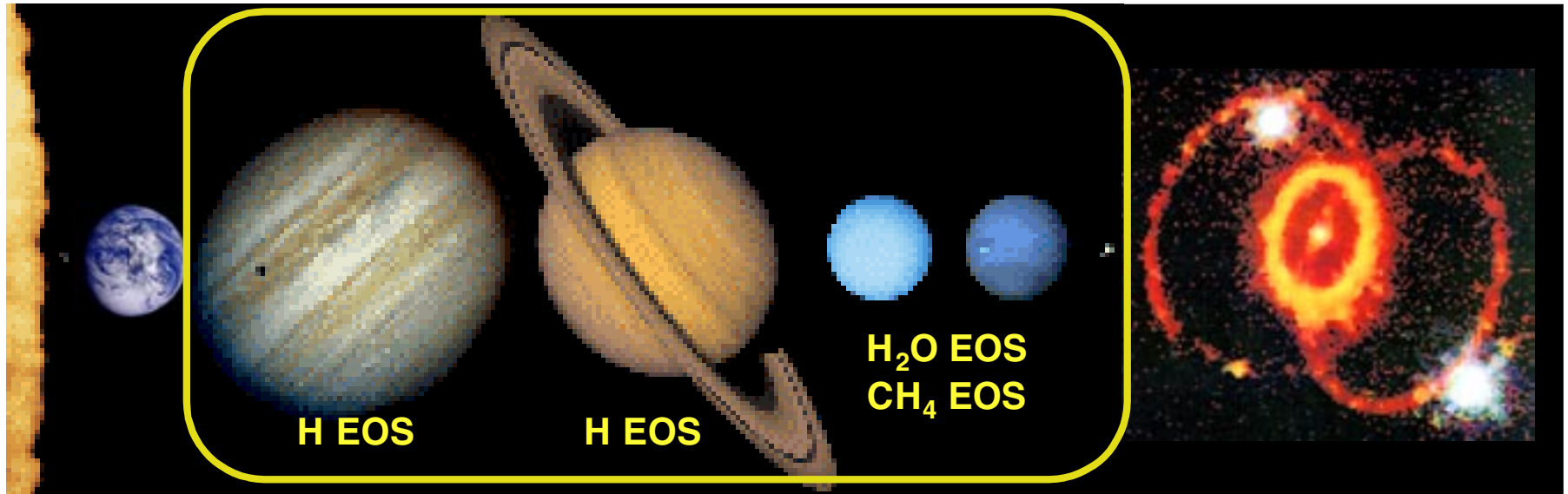
# Measurements show a continuous transition from insulating to metallic state



As shock pressure increases, D2 goes from transparent to a metallic reflector

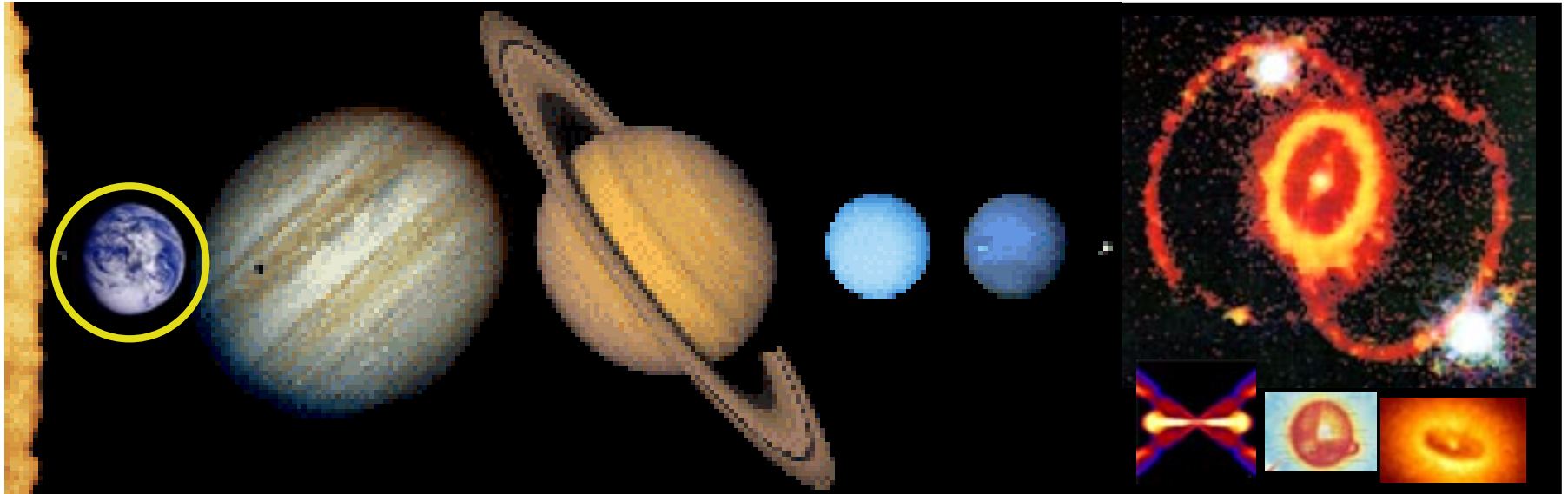
- **Models of Jupiter and giant planets are being reevaluated**

# Outline



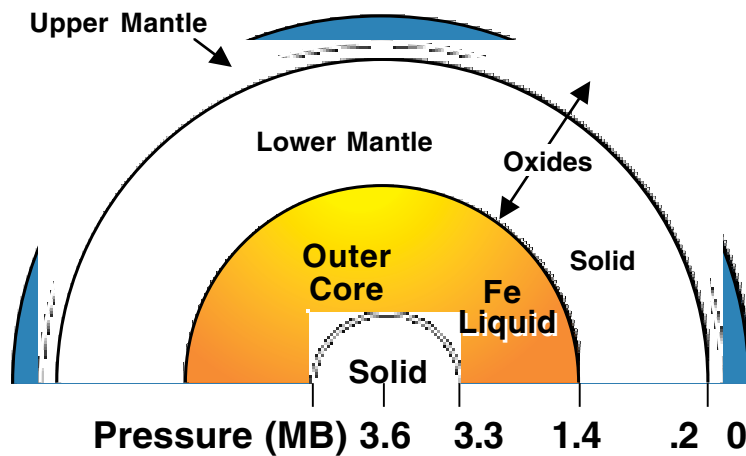
- **Significant advances in high energy density physics**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport and atomic physics*
- **Future directions**

# Outline

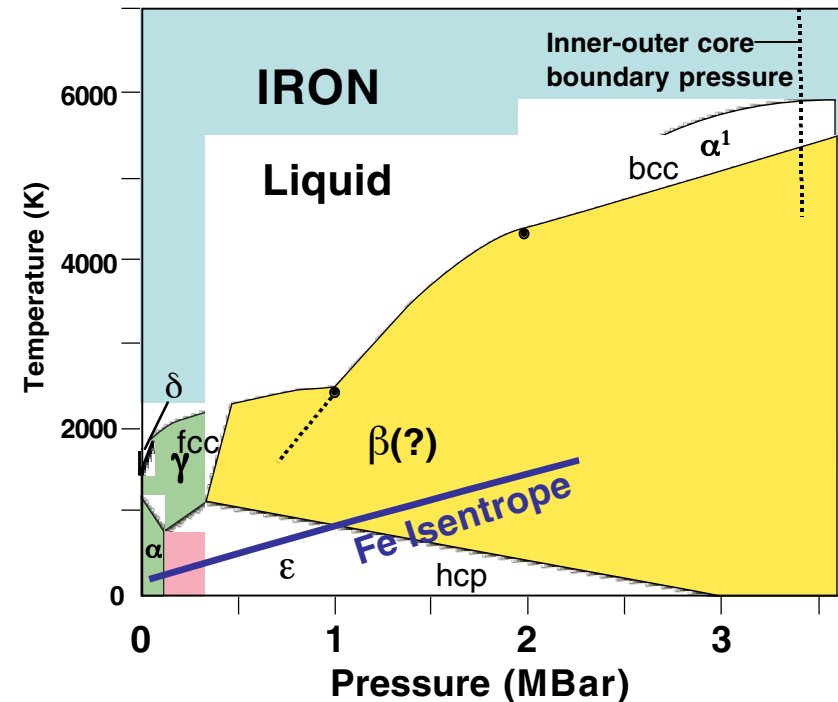


- **Significant advances in high energy density physics**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport and atomic physics*
- **Future directions**

# A new area of high energy density physics is the study of matter in the solid state under high pressure



- Earth's core contains solid Fe surrounded by liquid Fe
- Fe phase  $\leftrightarrow$  earth's magnetic field



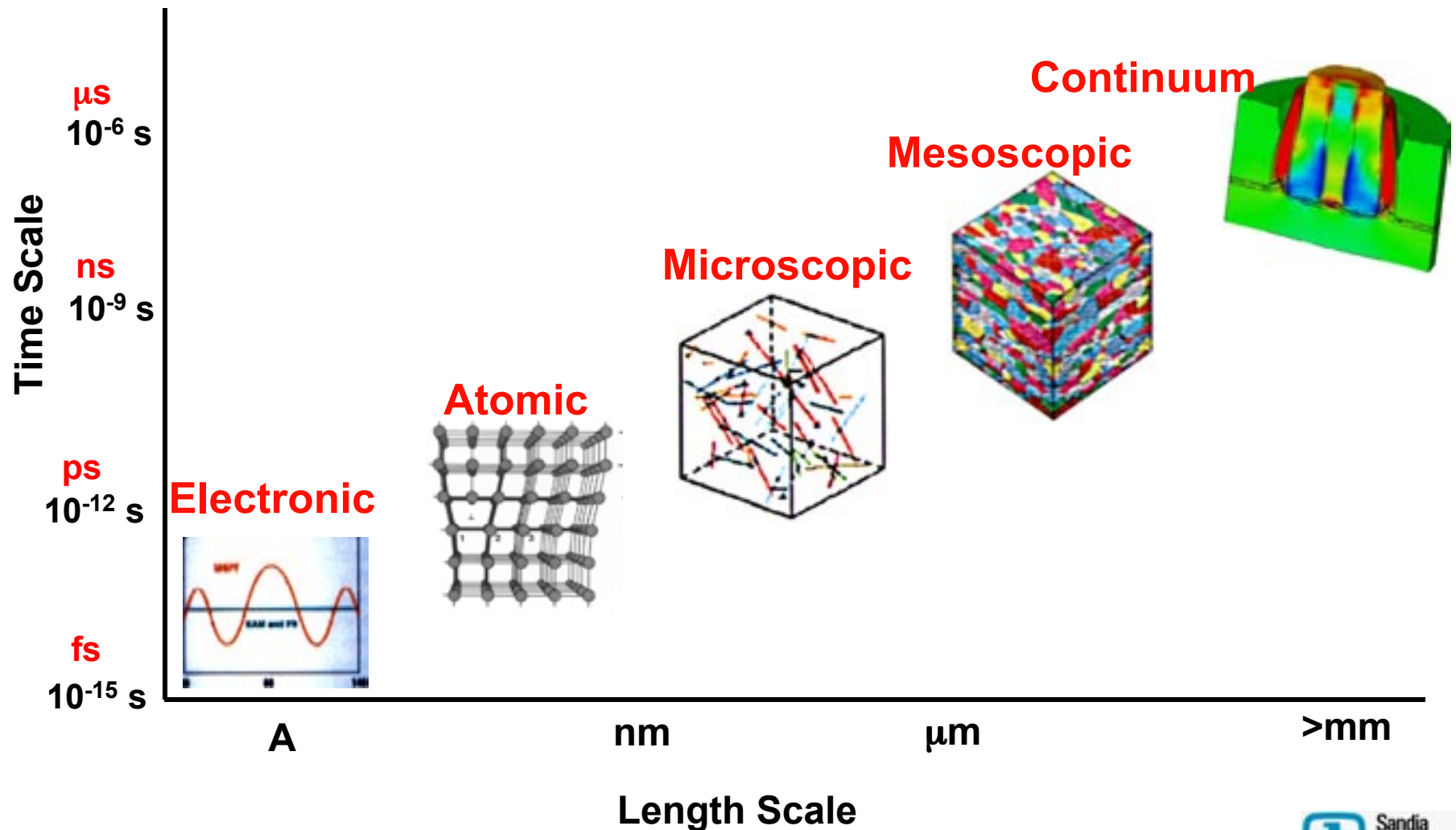
Anderson, O. L. *Science* 278 (1997)

- High-pressure phase boundaries and structure ?
- Mechanical properties at high pressure?

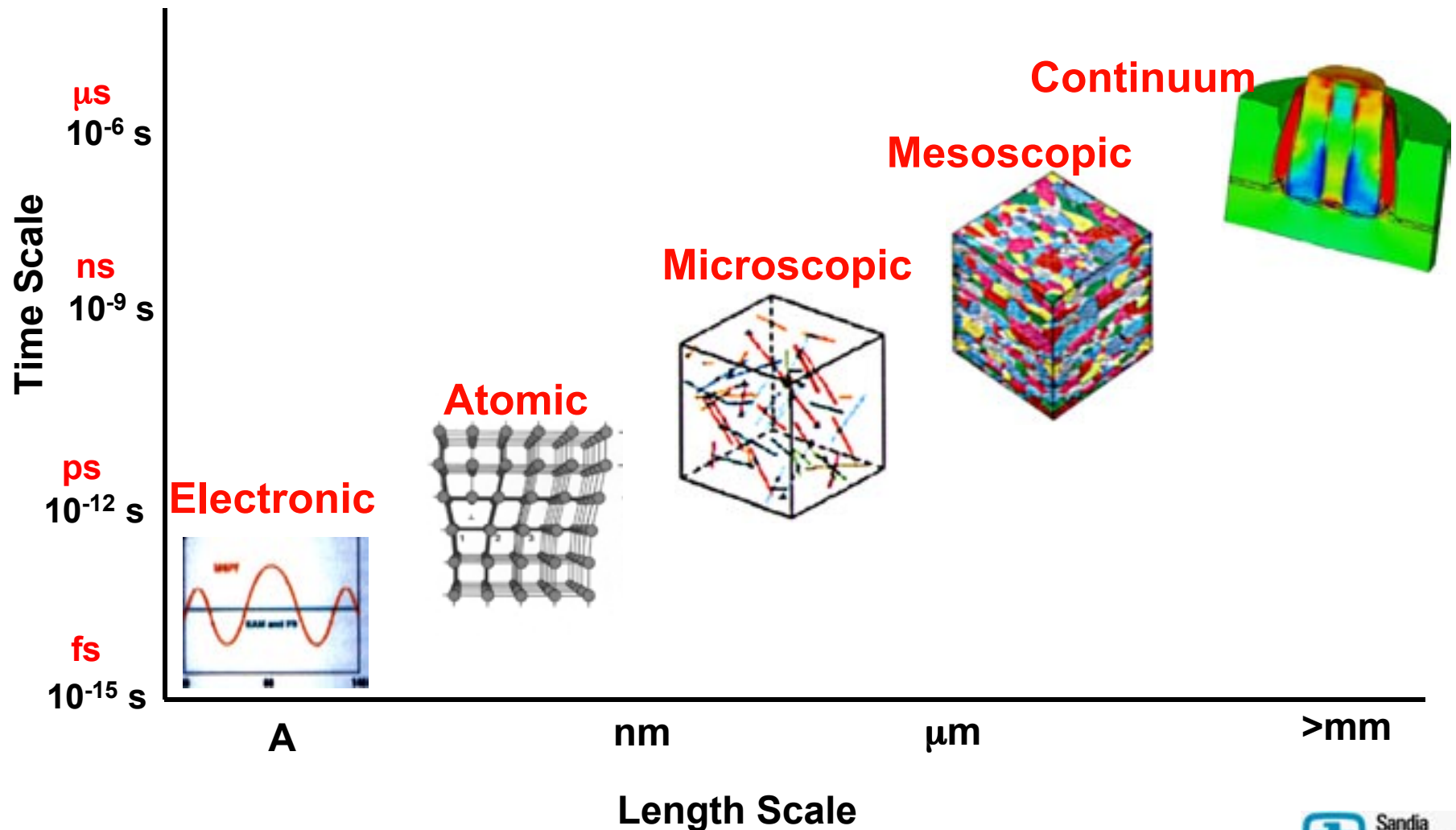
Compression along the isentropic can access high pressure solid states



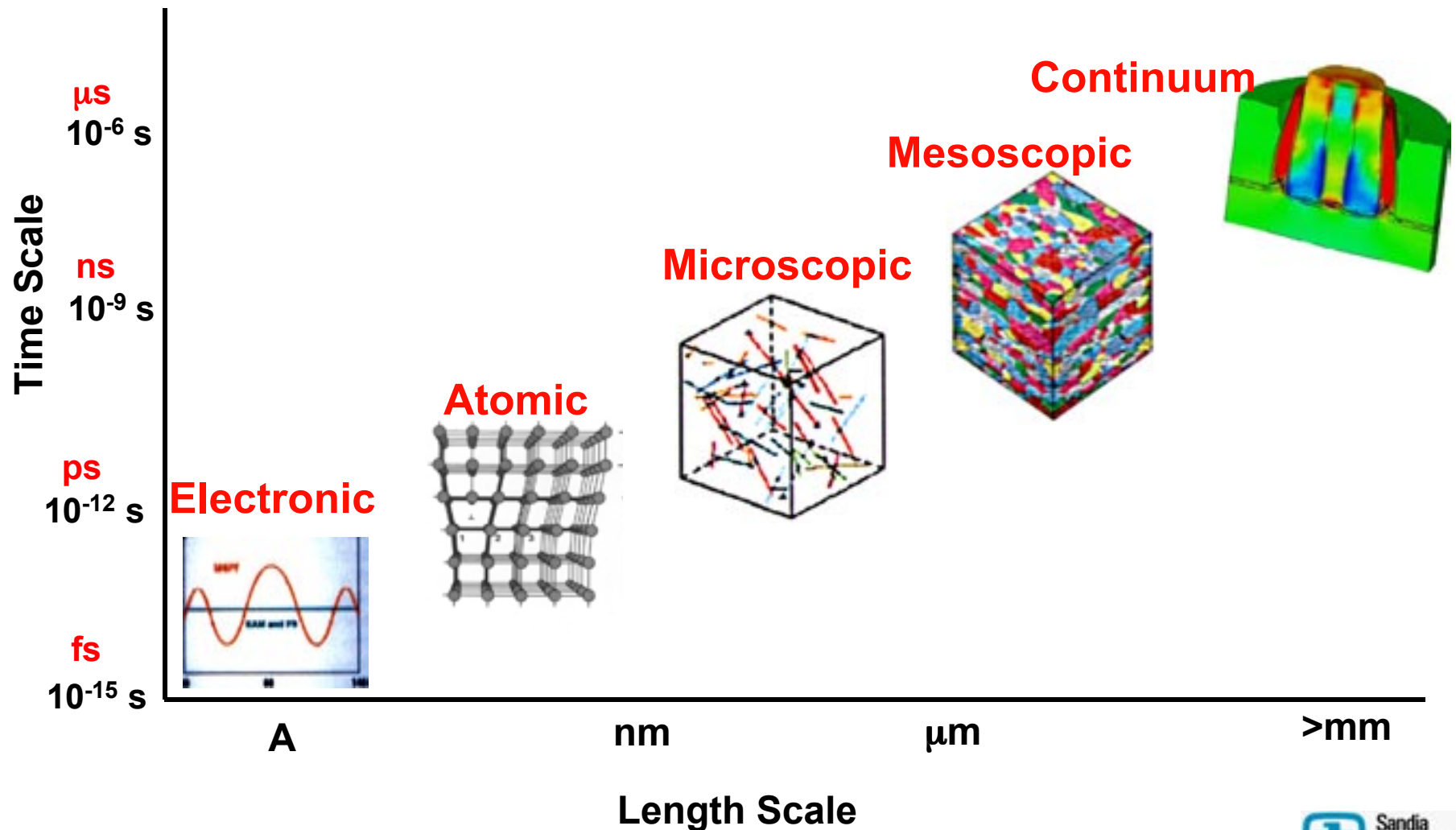
# Material response to compression in solids is complex and occurs on different scales



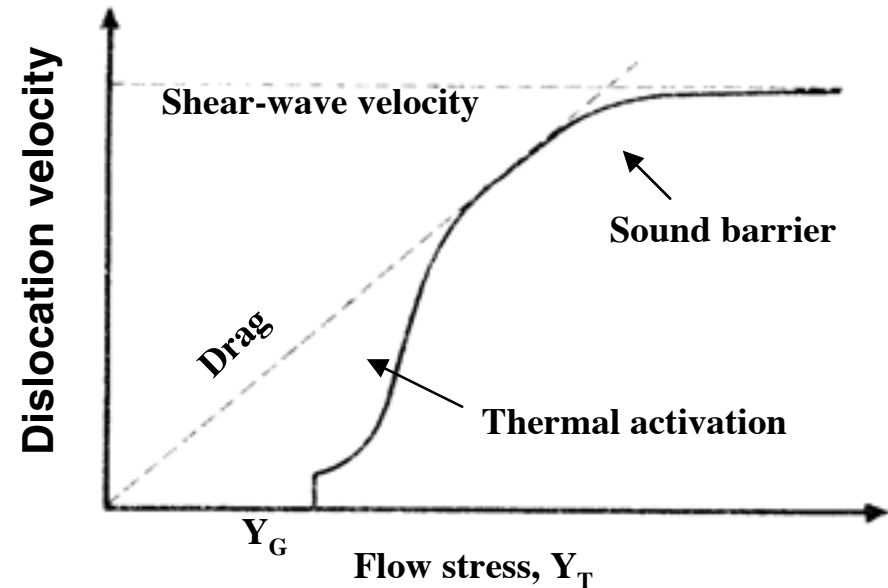
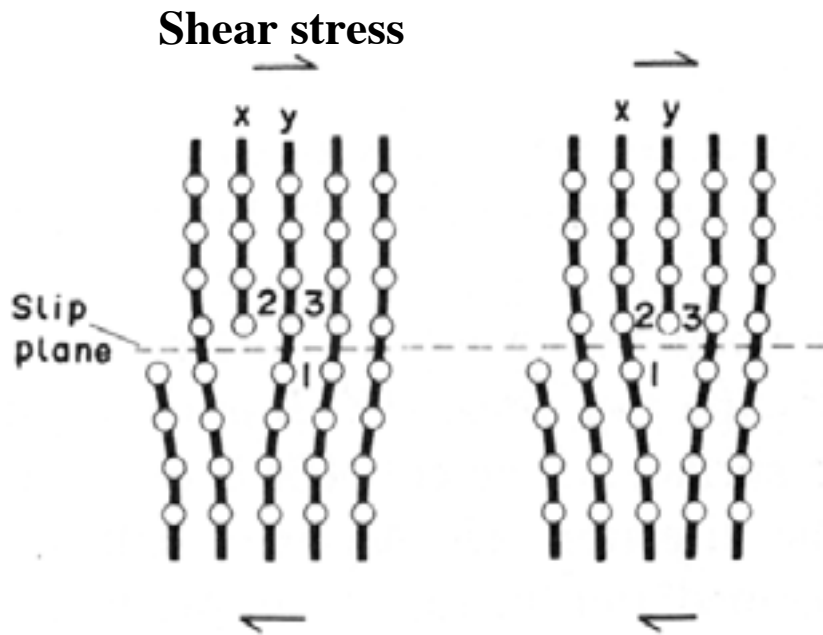
# Material response to compression in solids is complex and occurs on different scales



# Material response to compression in solids is complex and occurs on different scales



# Strength is characterized by a material's resistance to dislocation transport



M.A. Meyers, Dynamic Behavior of Materials  
(Wiley, 1994), pp. 358-361

Shear stress

Deformation  $\sim$  strain rate  $\sim$  ave. disloc. vel.

$$\dot{\epsilon} = \rho_m b \bar{v}_d$$

Dislocation density

Three categories of deformation are relevant :

1. Thermally activated dislocation transport
2. Dislocation glide resisted by phonon drag
3. Dislocation "speed limit"

The physics mechanisms underlying solid-state deformation can be categorized according to dislocation velocity or strain rate

## Mechanical properties

We need data at high pressures and varying strain rates to compare with models of strength



### Steinberg-Guinan constitutive model

Work hardening  $\rightarrow$  Shear modulus

$$Y = Y_o f(\epsilon) \frac{G(P, T)}{G_o}$$

P enhancement  $\rightarrow$  T softening

$$G = G_o \left( 1 + \left( \frac{G'_P}{G_o} \right) \frac{P}{\eta^{1/3}} - \left( \frac{G'_T}{G_o} \right) (T - 300) \right)$$

- SG semiempirical model with data at low pressures and strain rates
- Strain rate independent

## Mechanical properties

We need data at high pressures and varying strain rates to compare with models of strength



### Steinberg-Guinan constitutive model

$$Y = Y_o f(\epsilon) \frac{G(P, T)}{G_o}$$

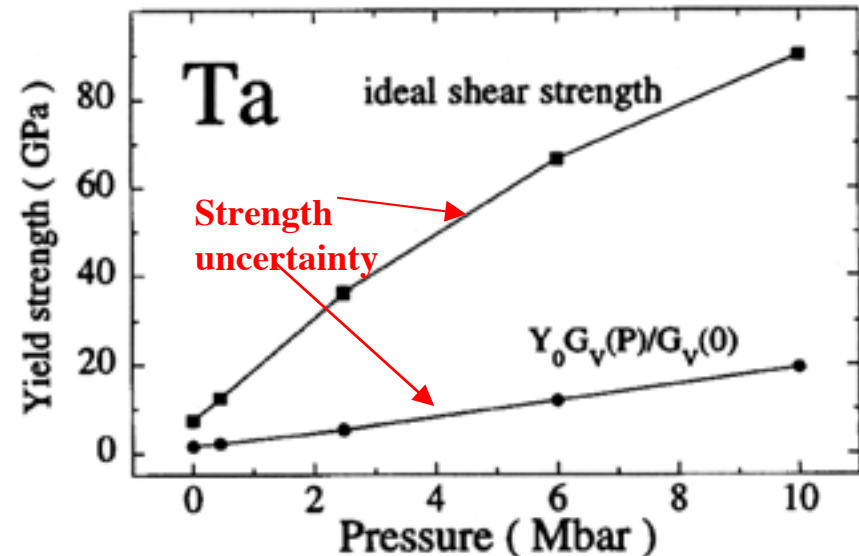
Work hardening  $\rightarrow f(\epsilon)$       Shear modulus  $\rightarrow G(P, T)$

$$G = G_o \left( 1 + \left( \frac{G'_P}{G_o} \right) \frac{P}{\eta^{1/3}} - \left( \frac{G'_T}{G_o} \right) (T - 300) \right)$$

P enhancement  $\rightarrow \frac{P}{\eta^{1/3}}$       T softening  $\rightarrow (T - 300)$

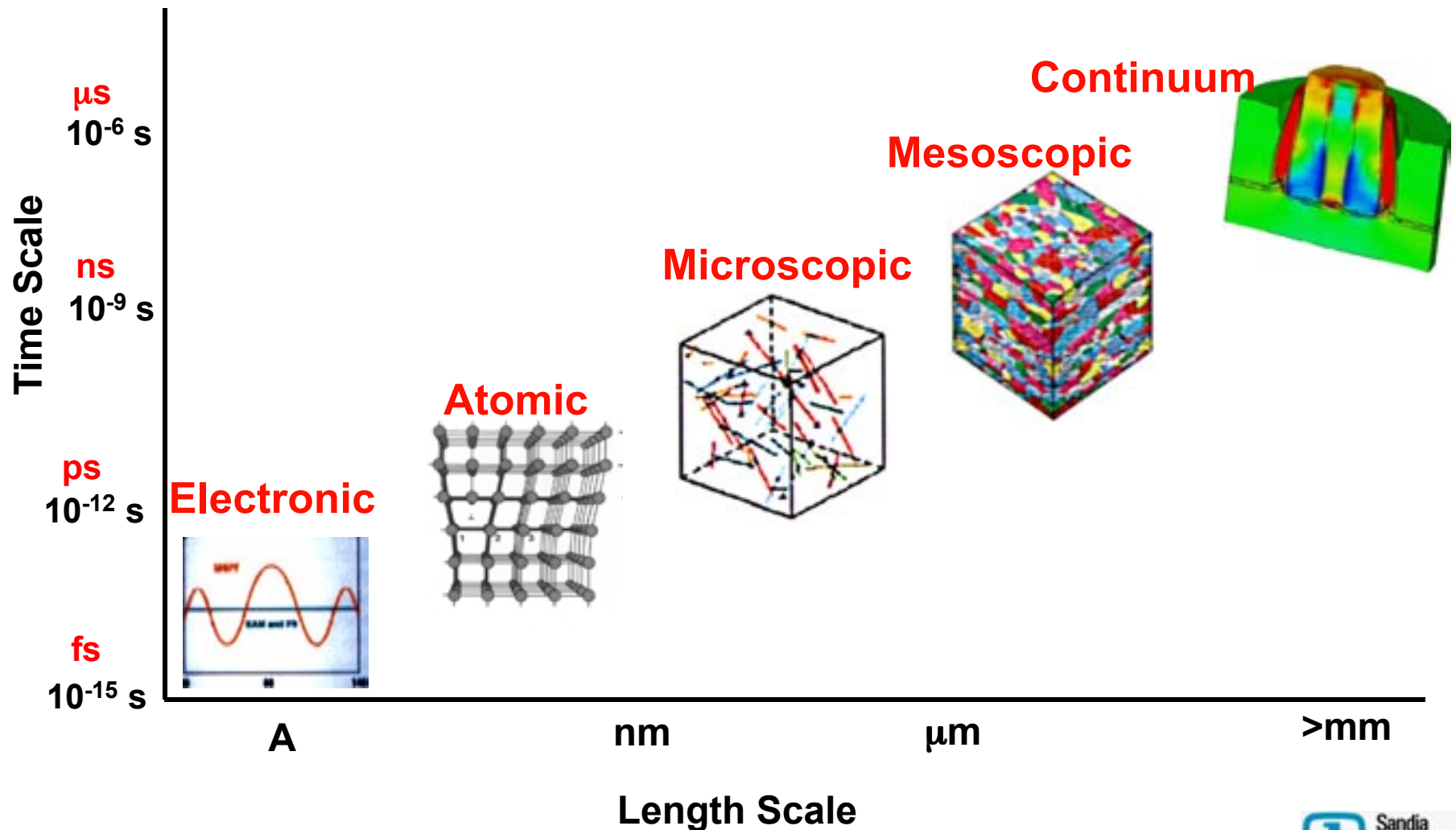
- SG semiempirical model with data at low pressures and strain rates
- Strain rate independent

### Ab-initio calculations



- At high pressure, there is an 4x difference in predictions of the strength

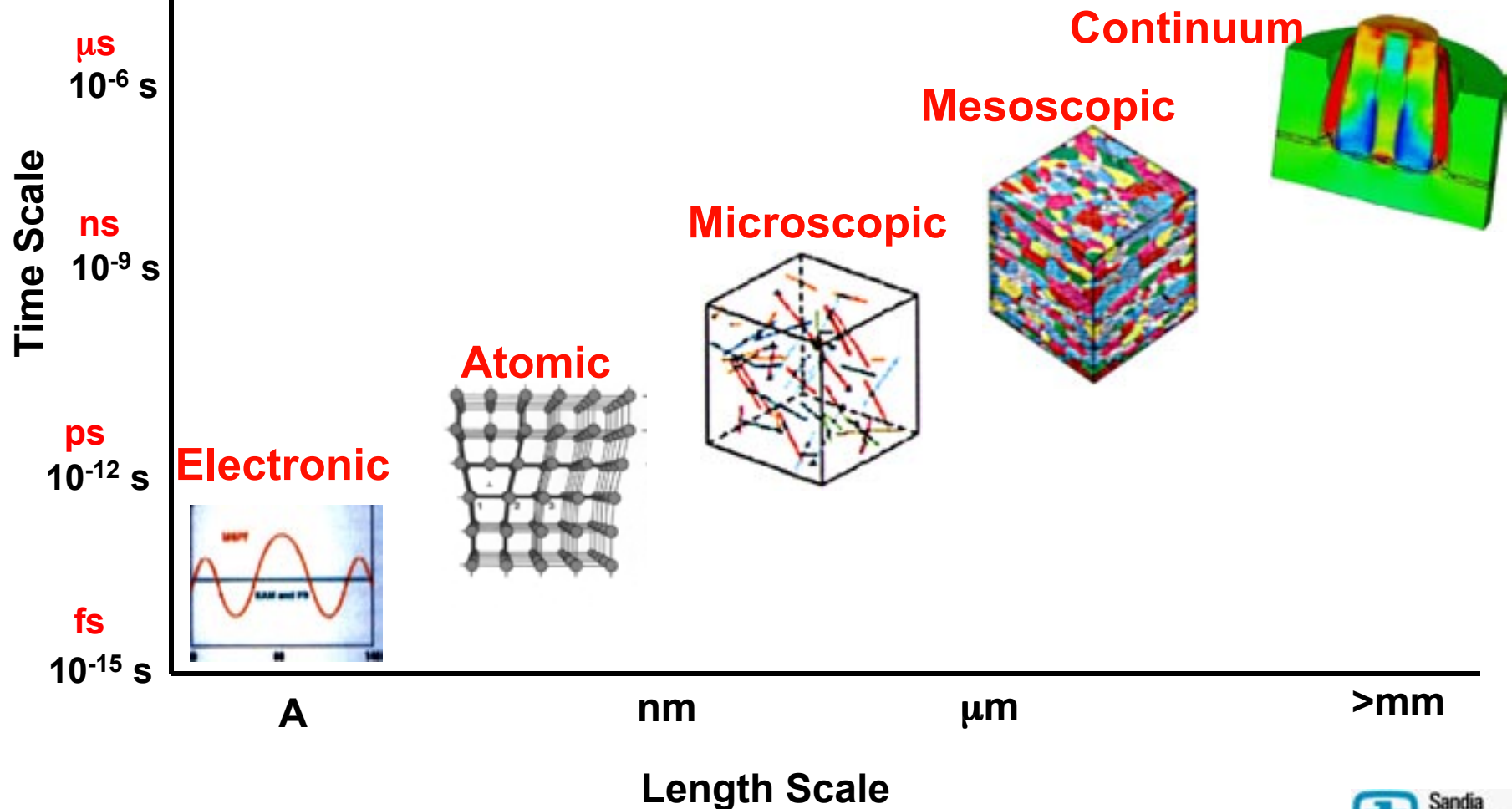
# Understanding mechanical properties requires measurements at each length scale





# Understanding mechanical properties requires measurements at each length scale

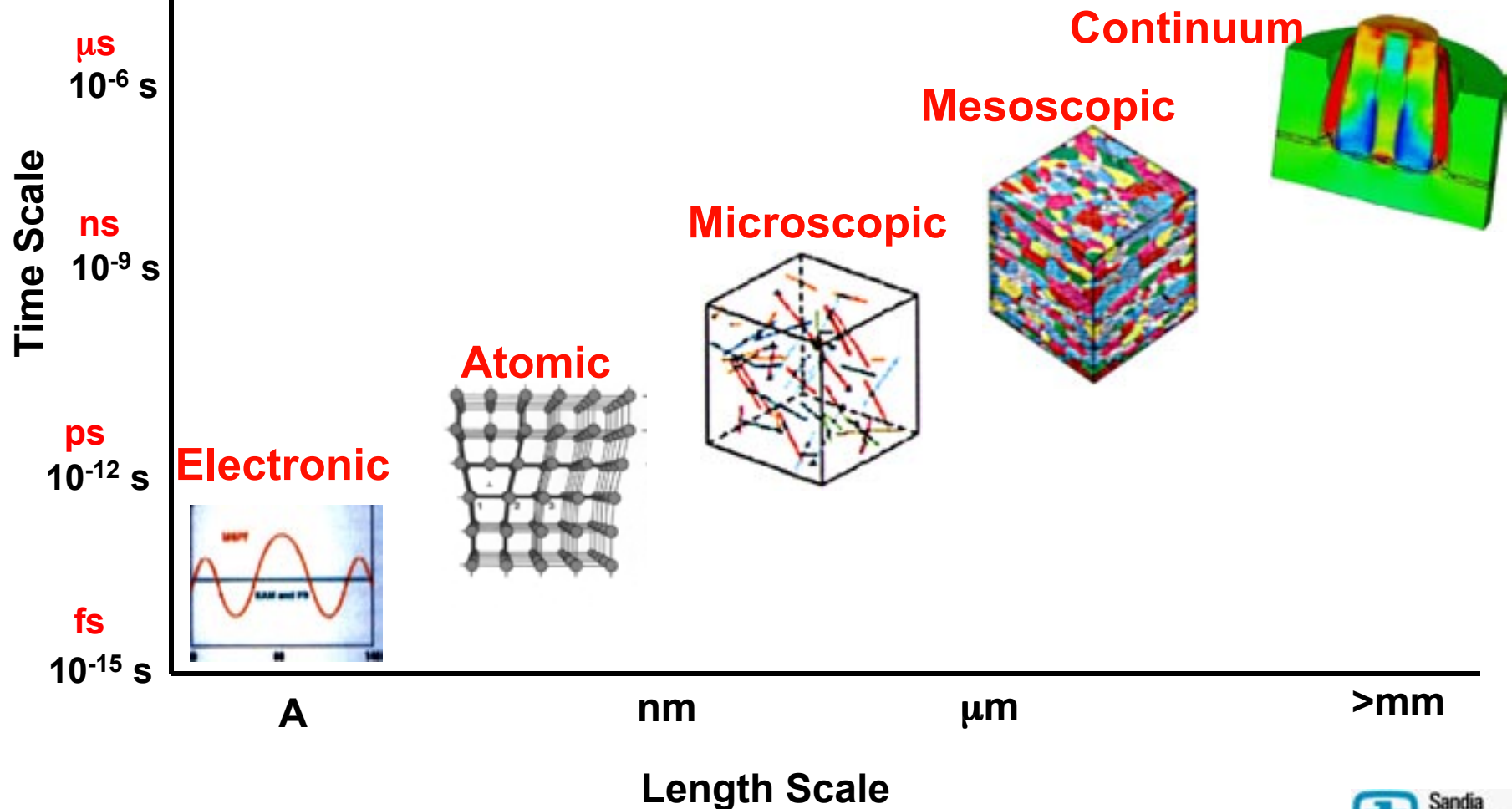
Diagnostics are being developed in each regime





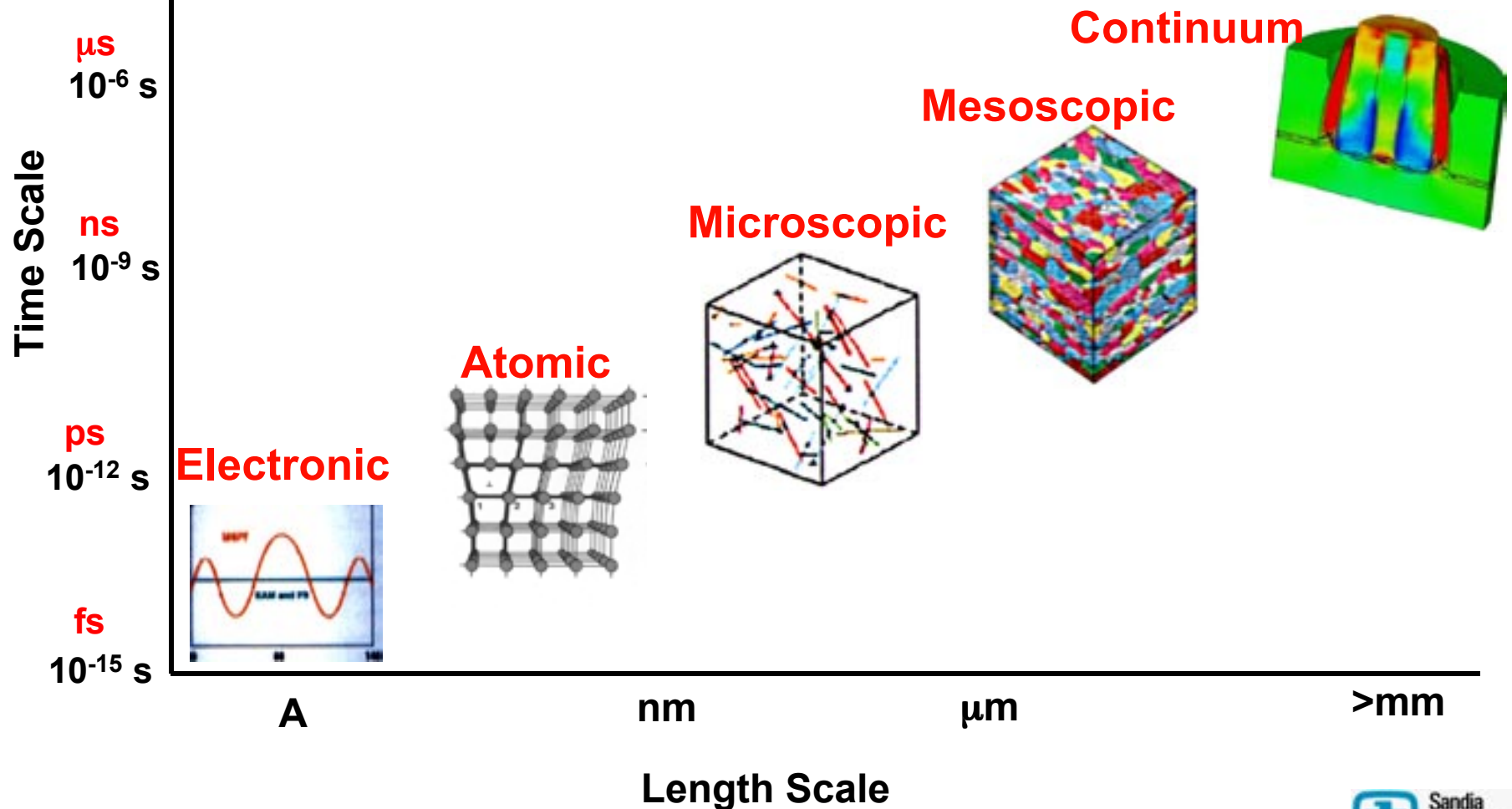
# Understanding mechanical properties requires measurements at each length scale

Diagnostics are being developed in each regime



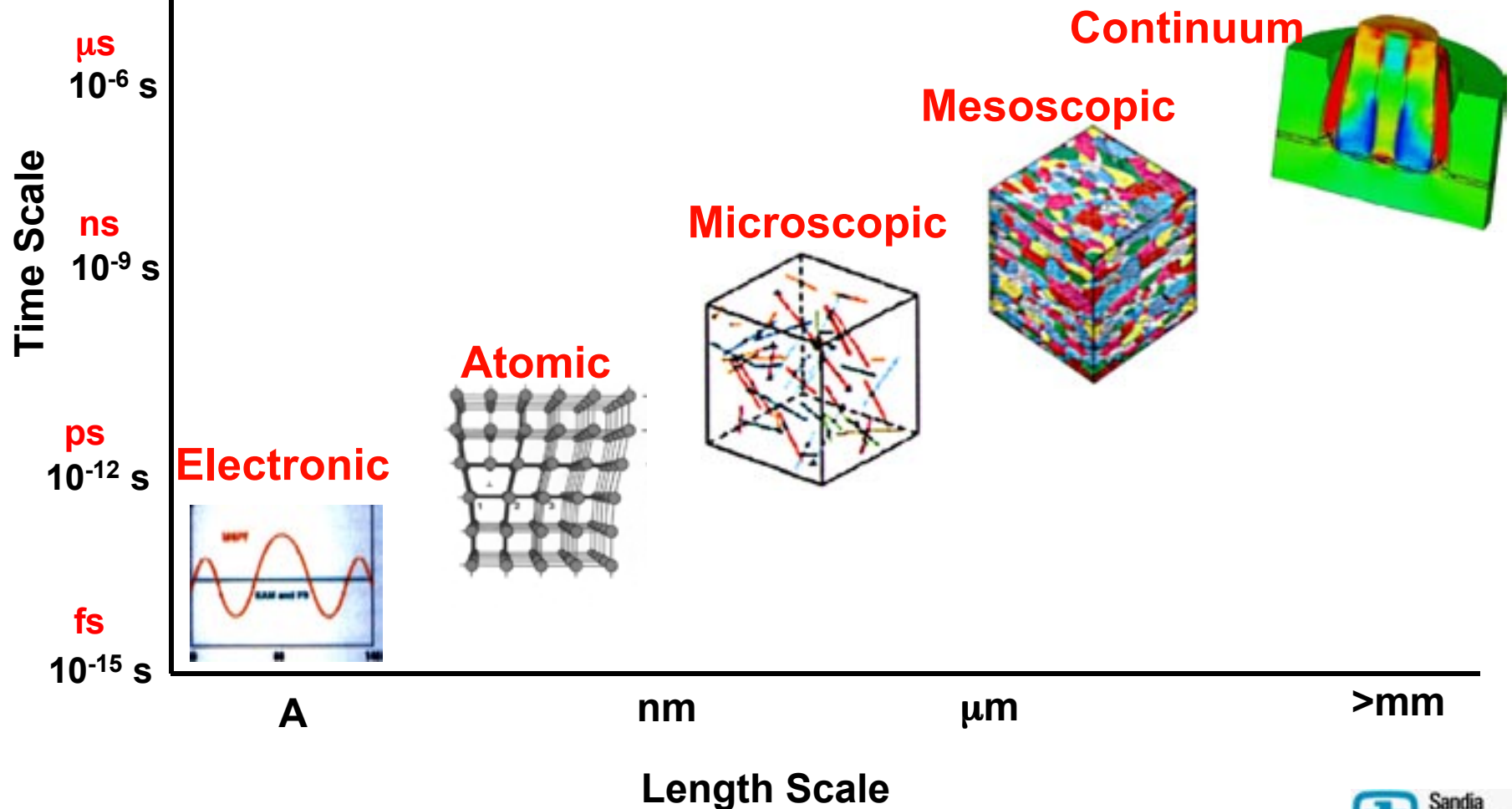
# Understanding mechanical properties requires measurements at each length scale

Diagnostics are being developed in each regime



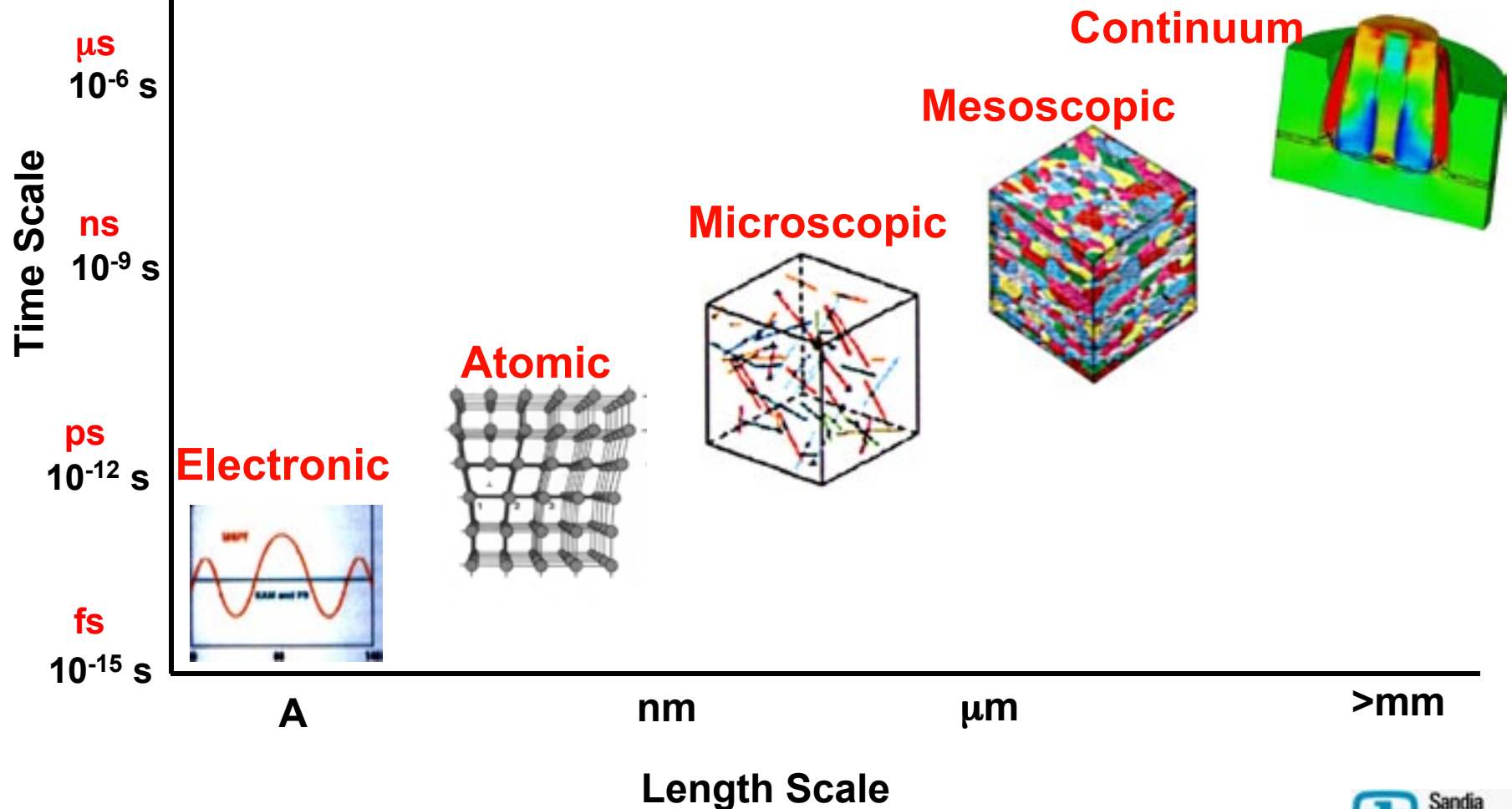
# Understanding mechanical properties requires measurements at each length scale

Diagnostics are being developed in each regime



# Understanding mechanical properties requires measurements at each length scale

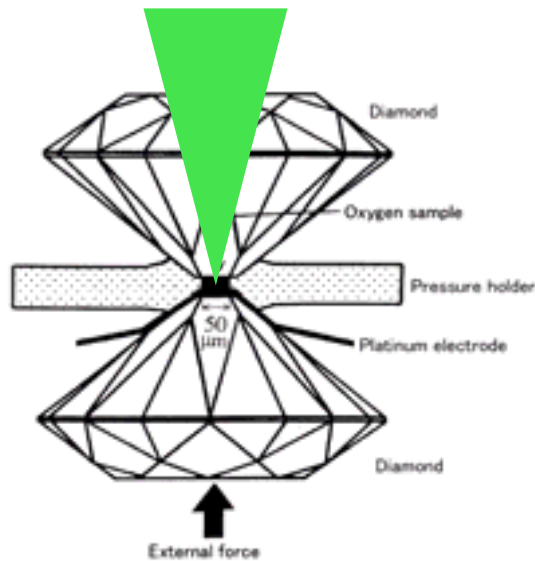
Diagnostics are being developed in each regime



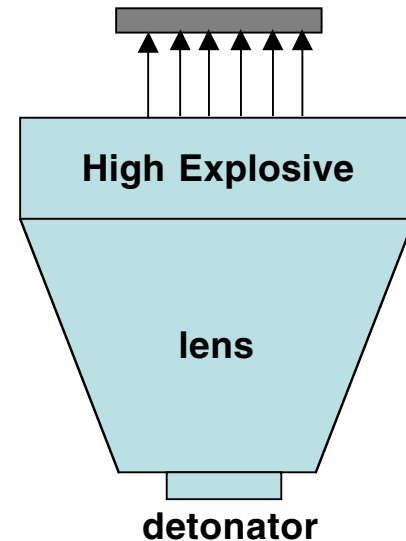
# It is difficult to get above ~ 1 Mbar isentropically



- **Laser heated DACs get to ~ 3 Mbars, 4000 K in small samples**



- **High explosives get to ~ 1 Mbars for strength measurements**



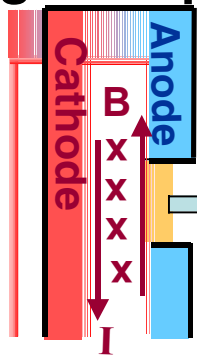
- **Isotherm, then isochor**
- **Thermodynamic properties (P, T) + structure on a synchrotron**

- **Isentropic**

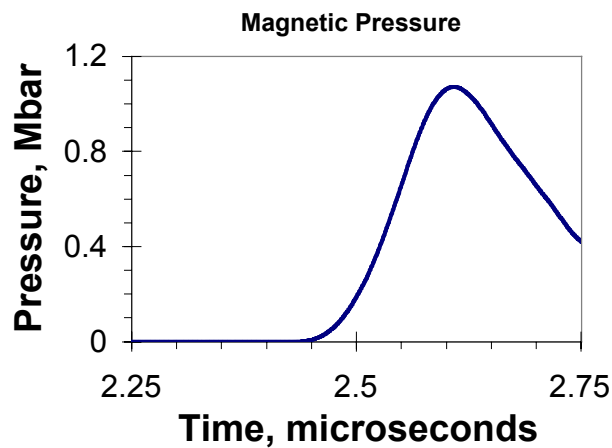
# Methods have been demonstrated to quasi-isentropically compress in the solid state to ~ 1 Mbar



## Magnetic pressure



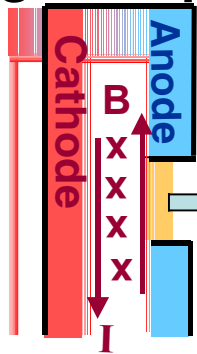
$$P = \frac{1}{2\mu_0} B^2$$



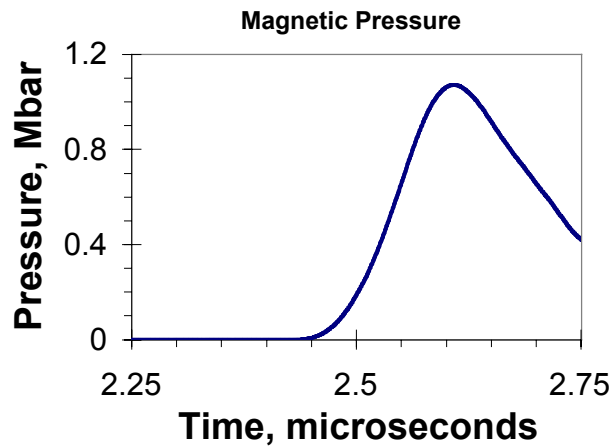
# Methods have been demonstrated to quasi-isentropically compress in the solid state to ~ 1 Mbar



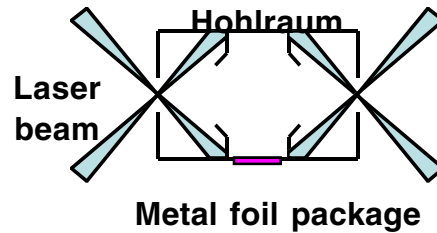
## Magnetic pressure



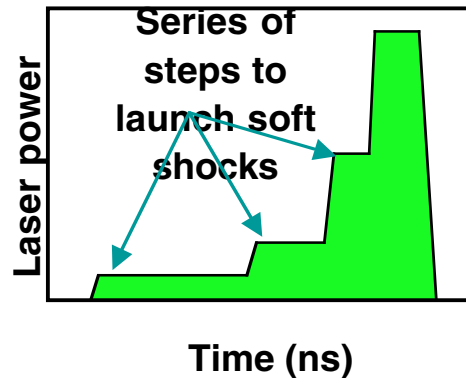
$$P = \frac{1}{2\mu_0} B^2$$



## Laser driven hohlraum



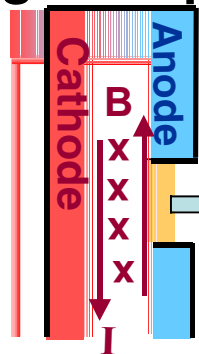
## Laser pulse shape



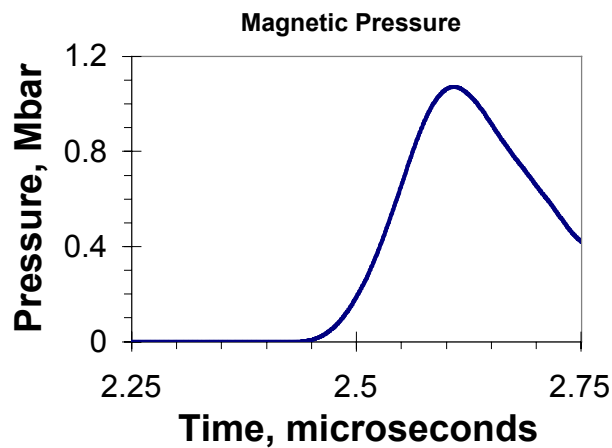
# Methods have been demonstrated to quasi-isentropically compress in the solid state to ~ 1 Mbar



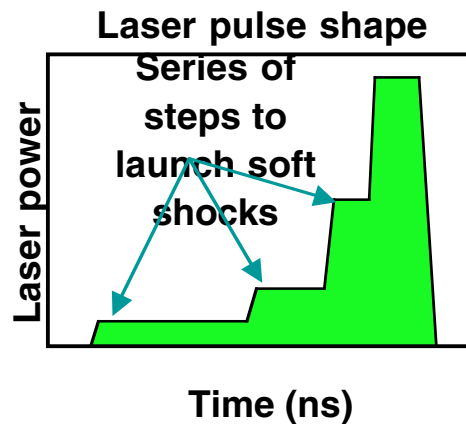
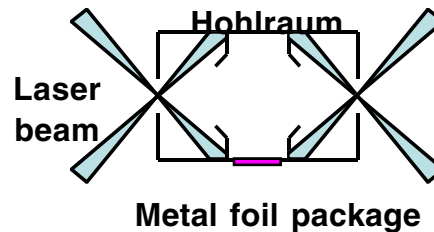
## Magnetic pressure



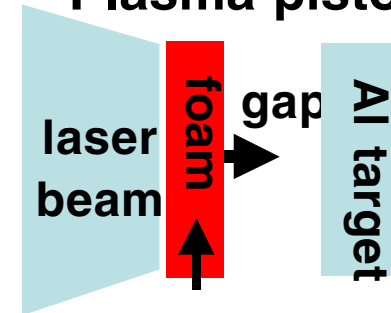
$$P = \frac{1}{2\mu_0} B^2$$



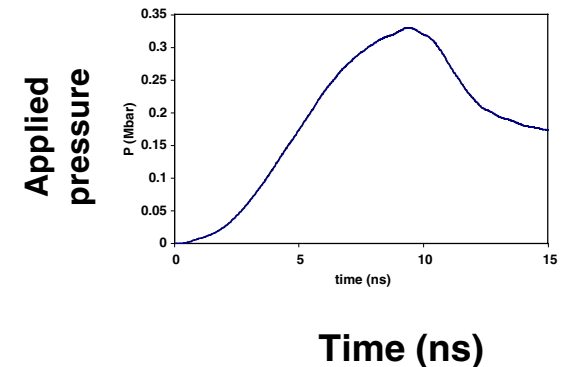
## Laser driven hohlraum



## Plasma piston

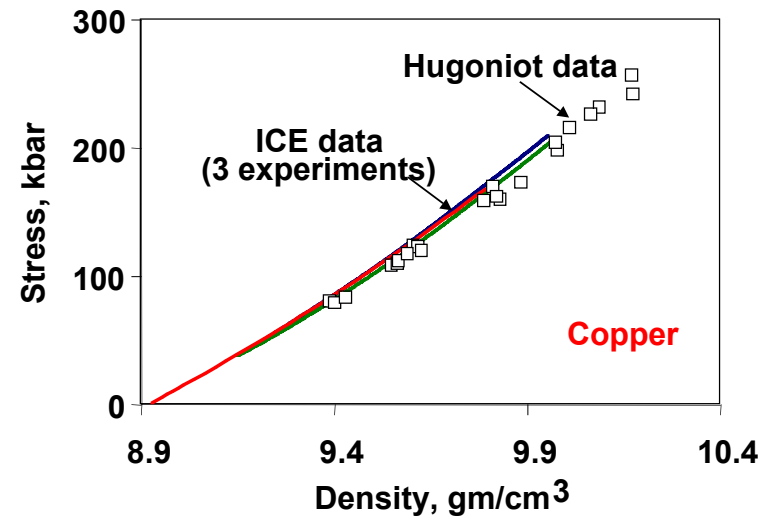
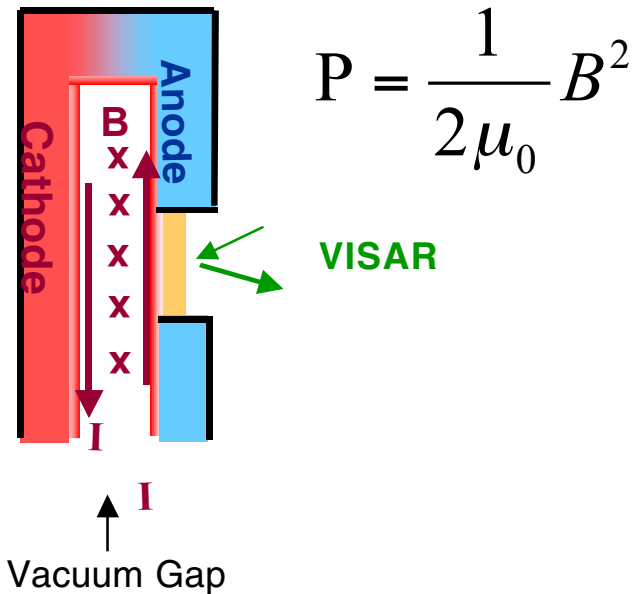


## unloading plasma reservoir

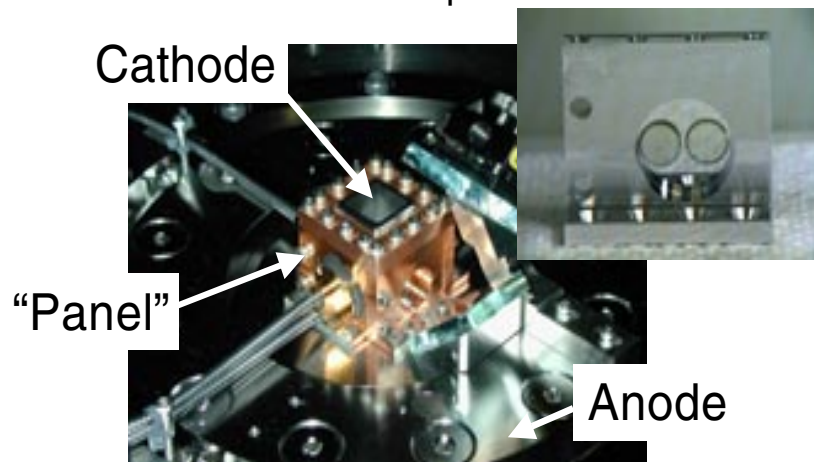




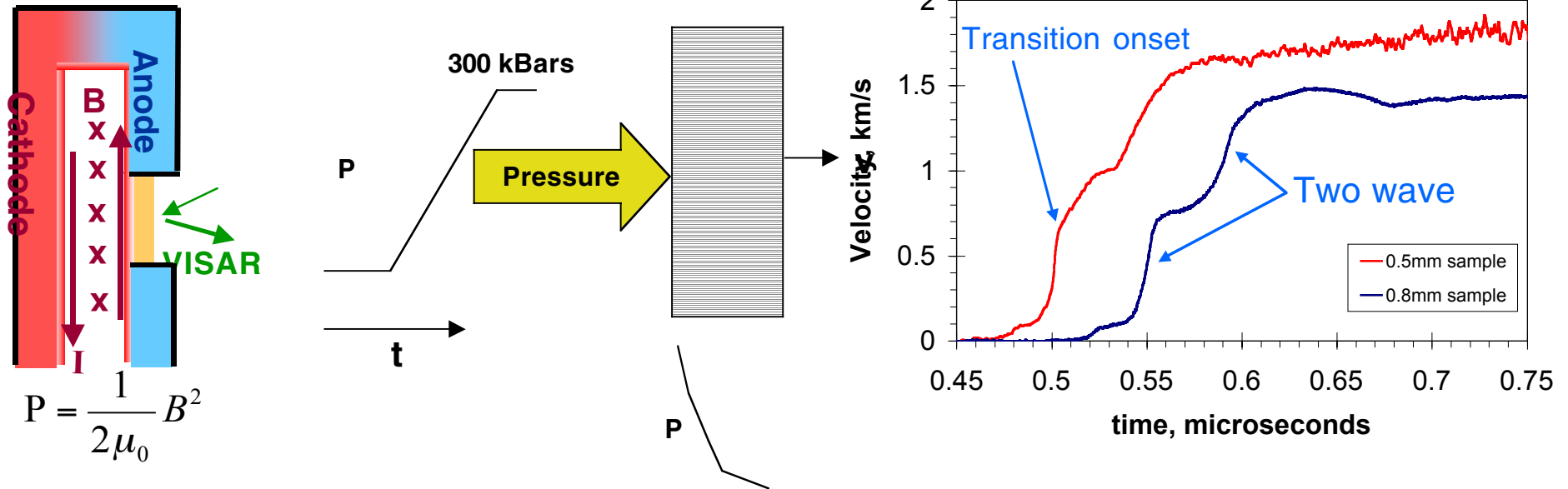
# Measurement along an isentrope provides continuous equation-of-state data



In comparison, a shock Hugoniot measurement provides 1 EOS datum per experiment



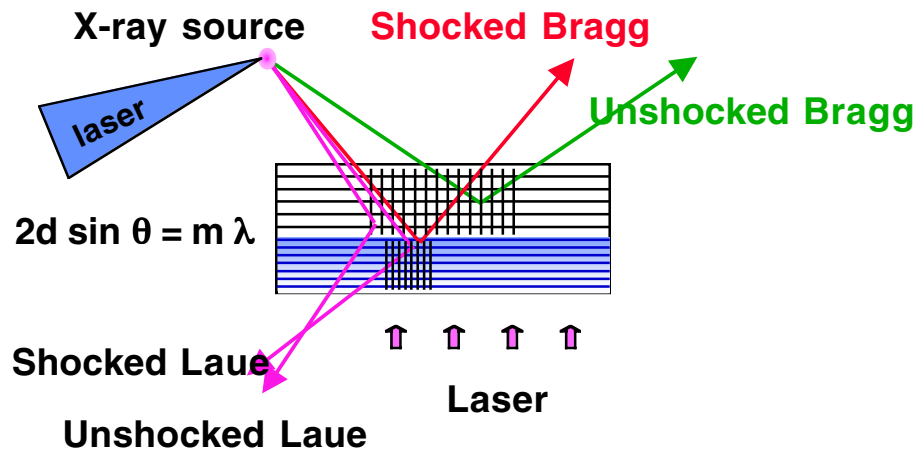
# A two wave structure is a signal of phase transition



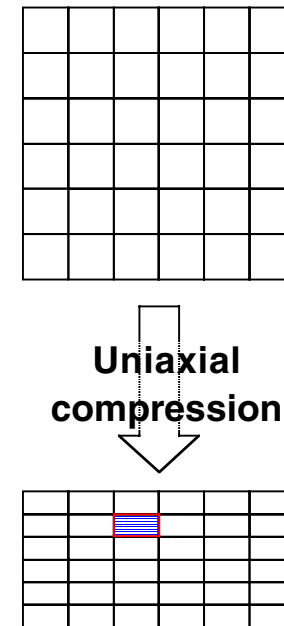
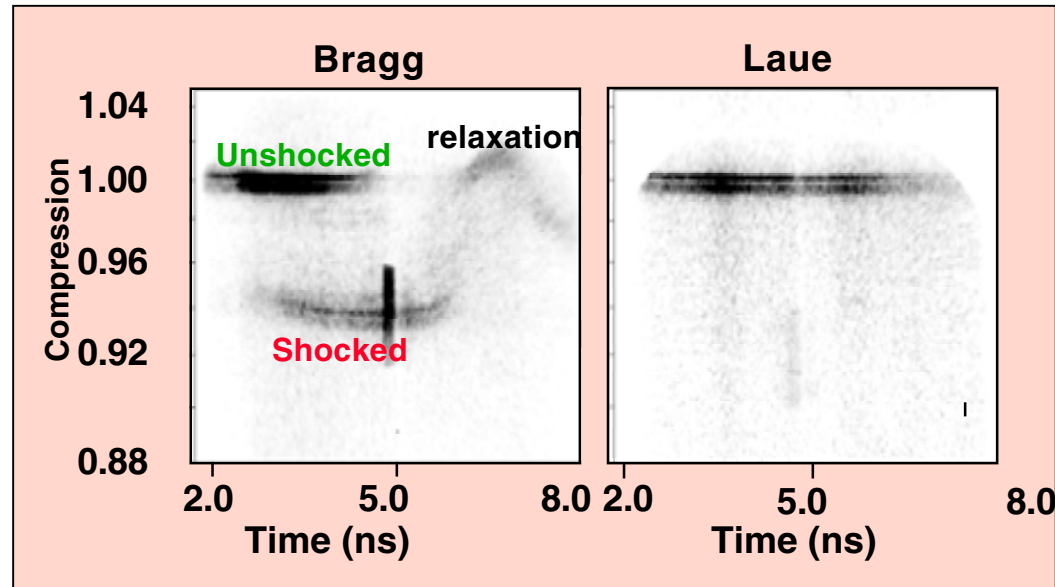
- The solid to solid (bcc-hcp) phase change was measured in Fe on Z
- Modeling allowed transition time to be determined

$$\tau \sim 40 \text{ ns}$$

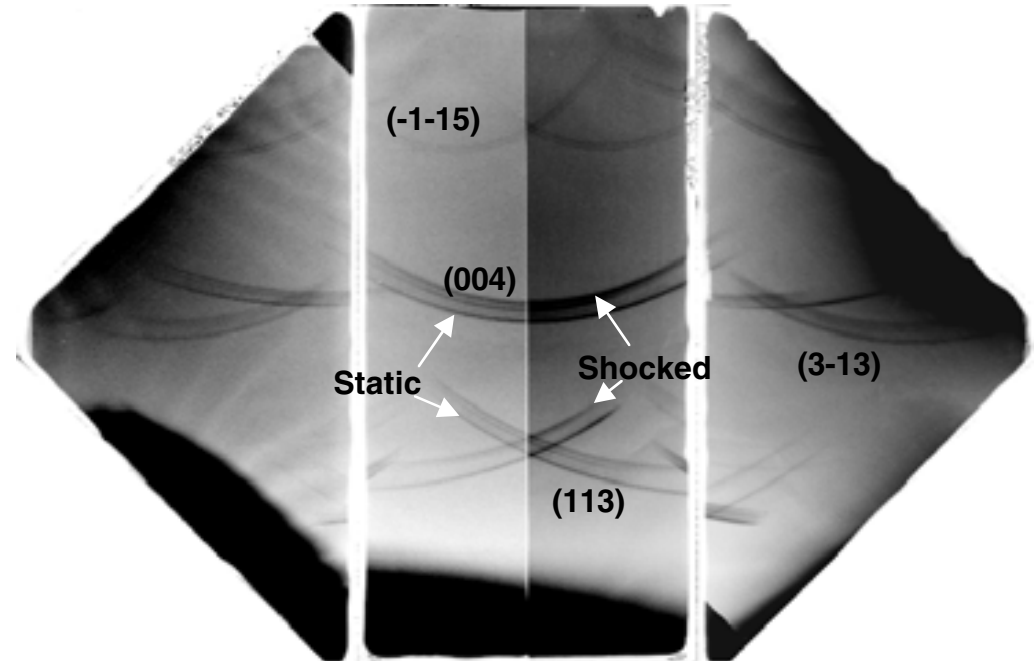
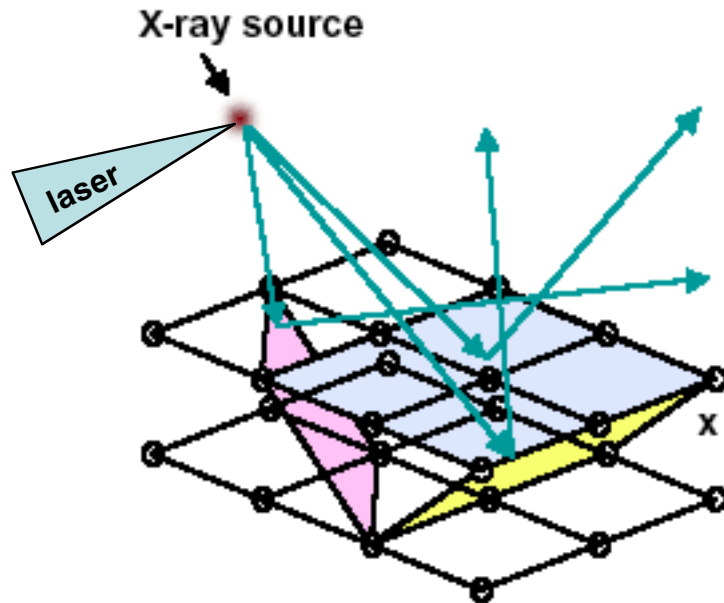
# X-ray diffraction has been demonstrated to dynamically measure the lattice structure



- Si responds uniaxially on a ns time scale to compression 2x over elastic limit

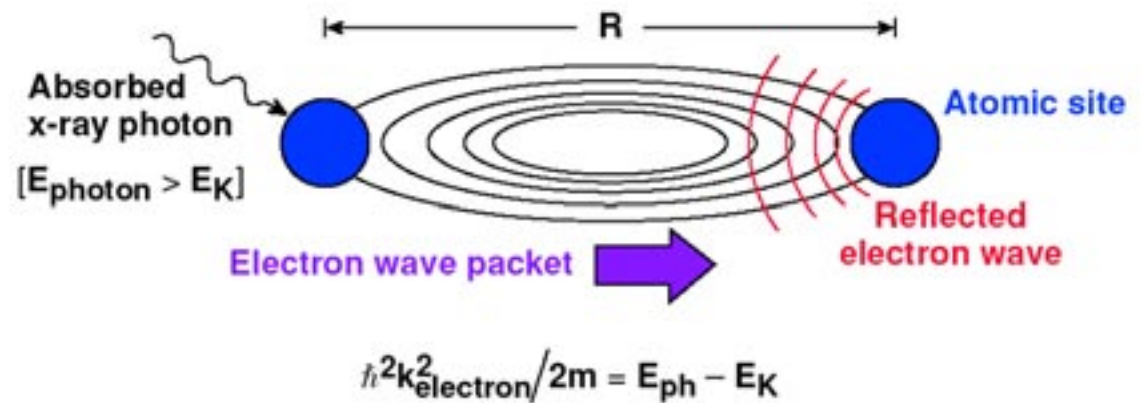
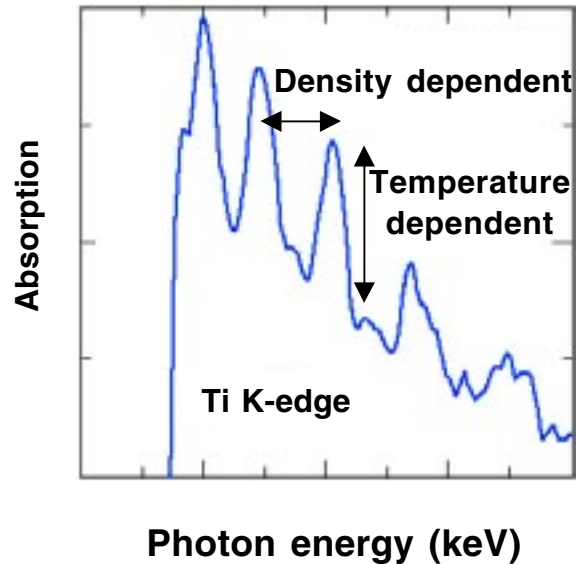


# The 3D lattice structure can be dynamically measured using a novel diffraction geometry



- Time-resolved X-ray source allows 3D lattice structure to be measured under dynamic conditions

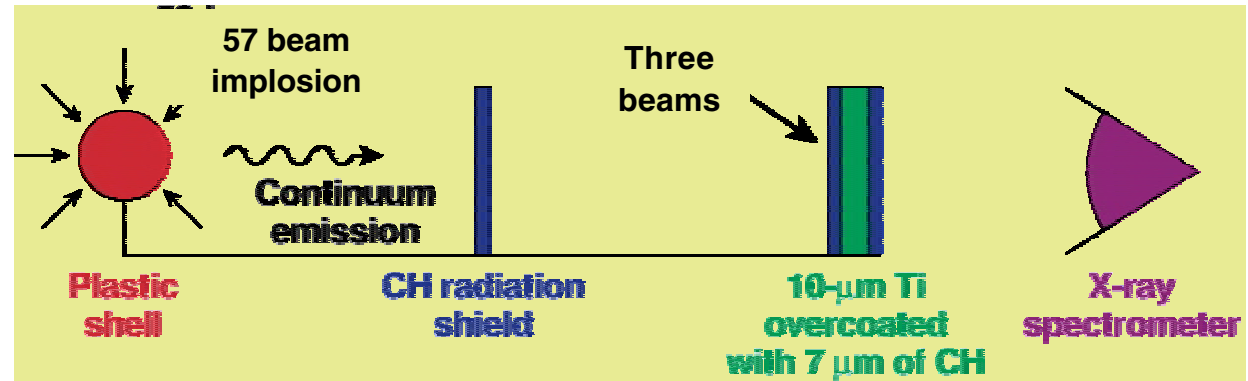
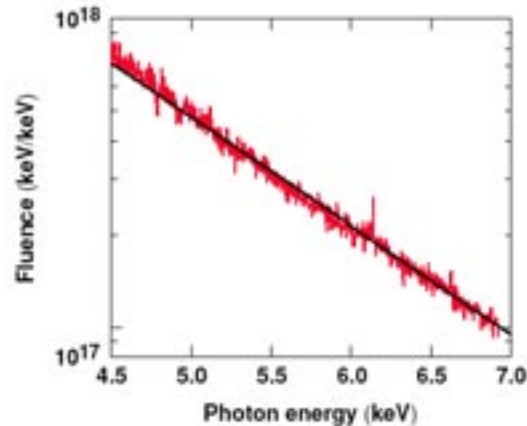
# Measurements of X-ray absorption near an edge provides information on density and temperature in a solid



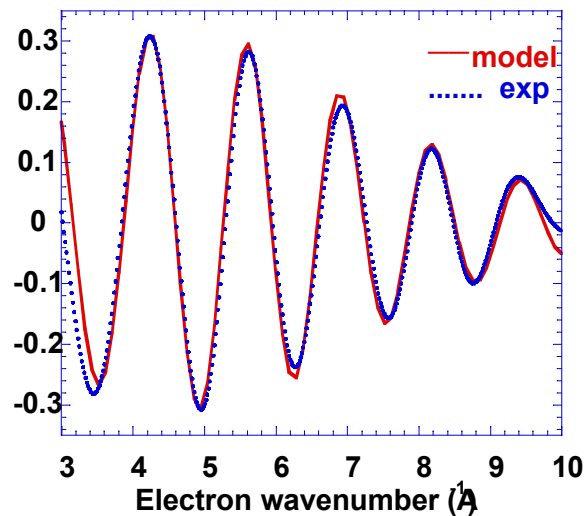
- EXAFS - interference of ejected electron waves with neighbors
  - Modulation  $\sim \sin(2kR) \Rightarrow$  density
  - Amplitude of oscillations  $\sim \exp[-2\sigma^2 k^2]$ 
    - $\sigma^2$  is the Debye- Waller factor
    - $\sigma^2 \sim f(T/\rho^{5/2})$

EXAFS technique developed on a synchrotron is now being applied on a laser

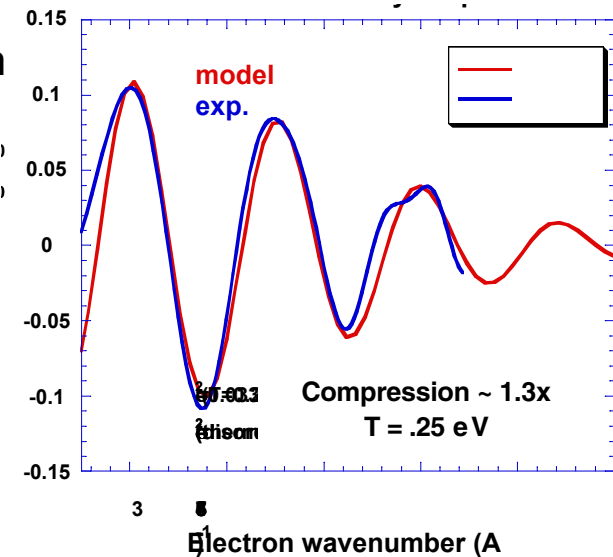
# EXAFS provides a way to measure T to 20-30% for isentropically compressed targets



Undriven target



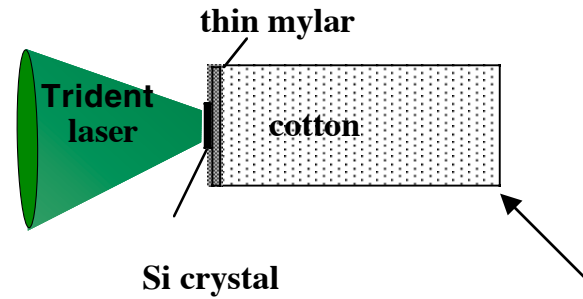
Driven target



- A capsule is imploded on the Omega laser to provide a continuum source for EXAFS diagnosis of a target

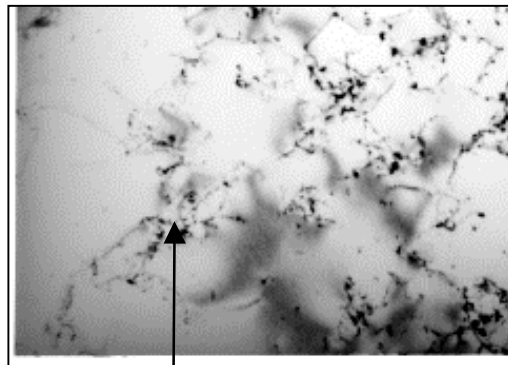
## Microscale

Recovery of compressed samples and metallurgical analysis show identical features to those shocked on a gas gun



A filled tube slow down and capture samples

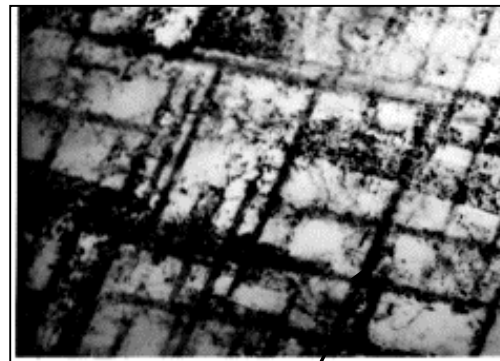
—Laser shocked samples have identical features to samples shocked on gas guns at the same pressure



40 J, Cu  
130 kbar

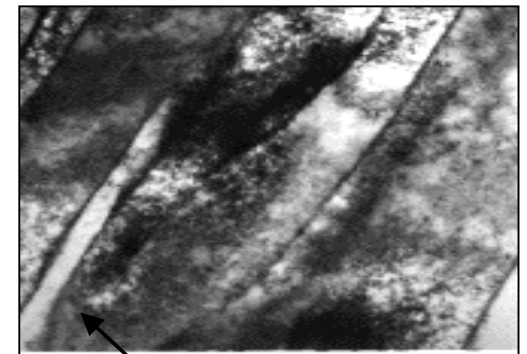
Dislocations

0.5  $\mu\text{m}$



205 J, Cu  
400 kbar

Micro-twins



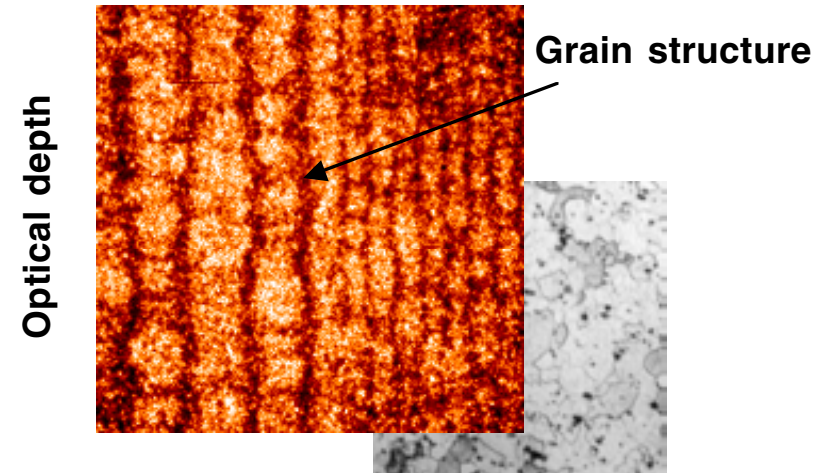
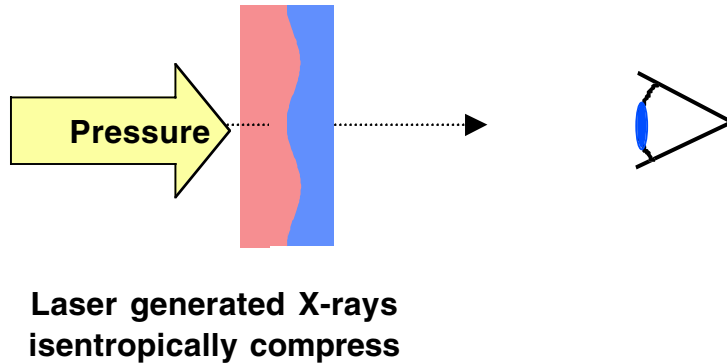
320 J, CH/Cu  
1.4 Mbar

Twinning



## Mesoscale and continuum

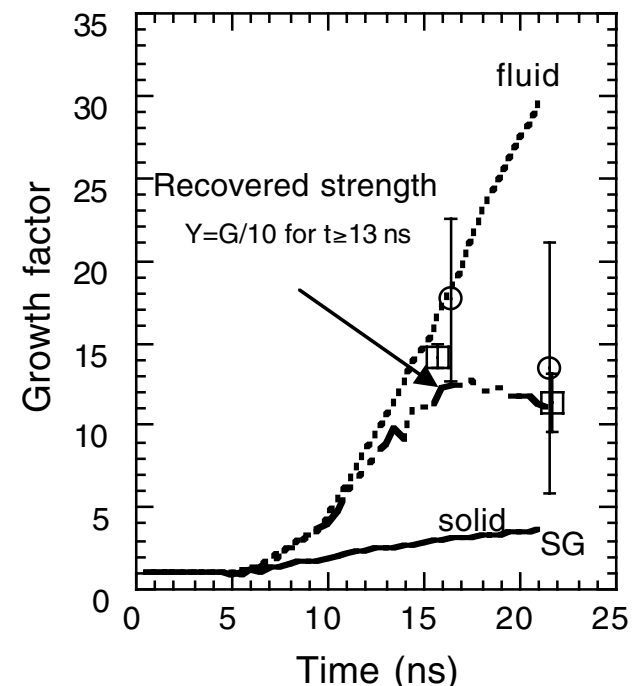
Strength in Al was inferred from measurements of Rayleigh-Taylor instability growth on Nova at ~ 1.4 MBars



Growth rates with strength are expected to be reduced from classical (fluid)

Grain structure can be measured

Data does not agree with predictions

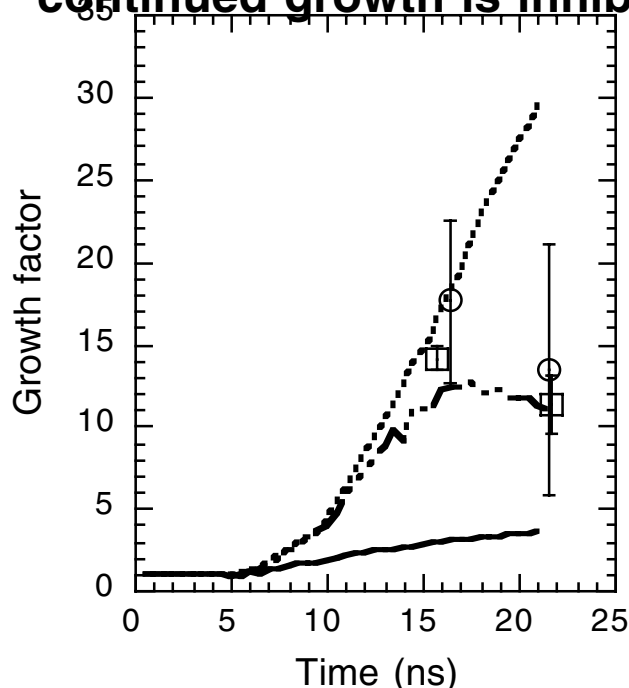




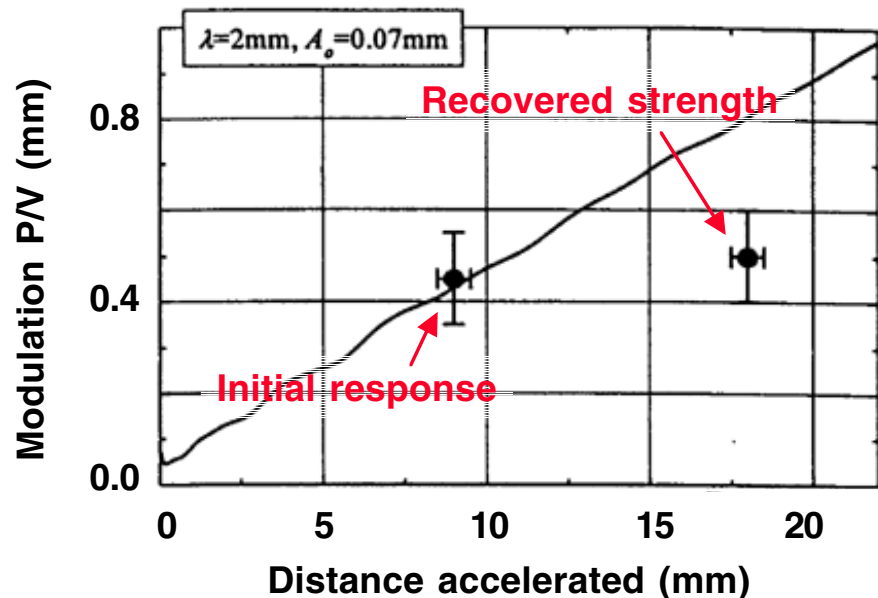
# The RT growth is nearly fluid at early times, but it is suppressed at later times



- Suggestive of data by Rayevsky/Lebedev
- High pressure strain causes localized heating and softening in shear bands; bulk Al flows as fluid due to deformation in these localized regions
- As heat dissipates the metal regains bulk solid strength and continued growth is inhibited

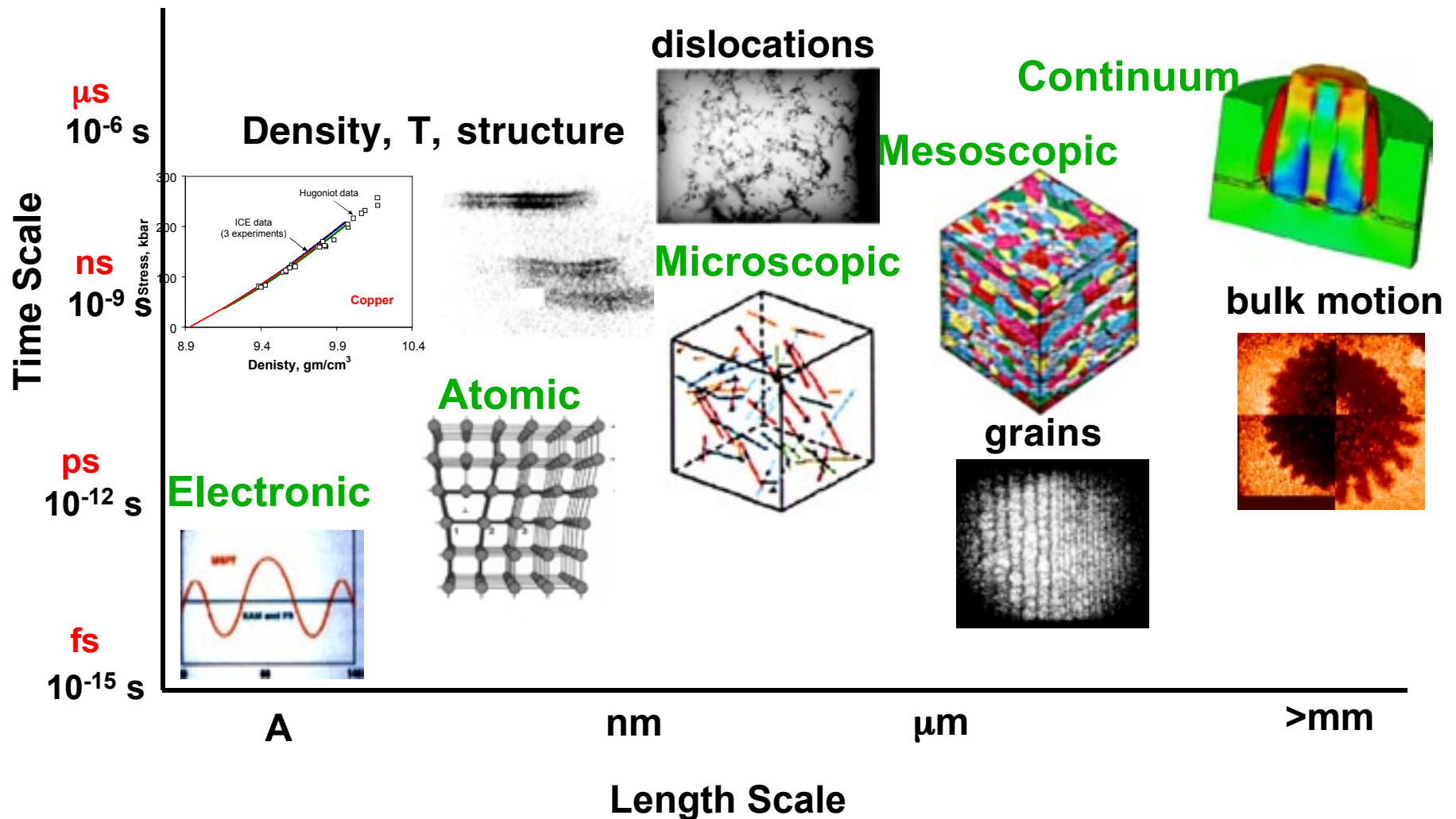


New data in new regimes => surprises

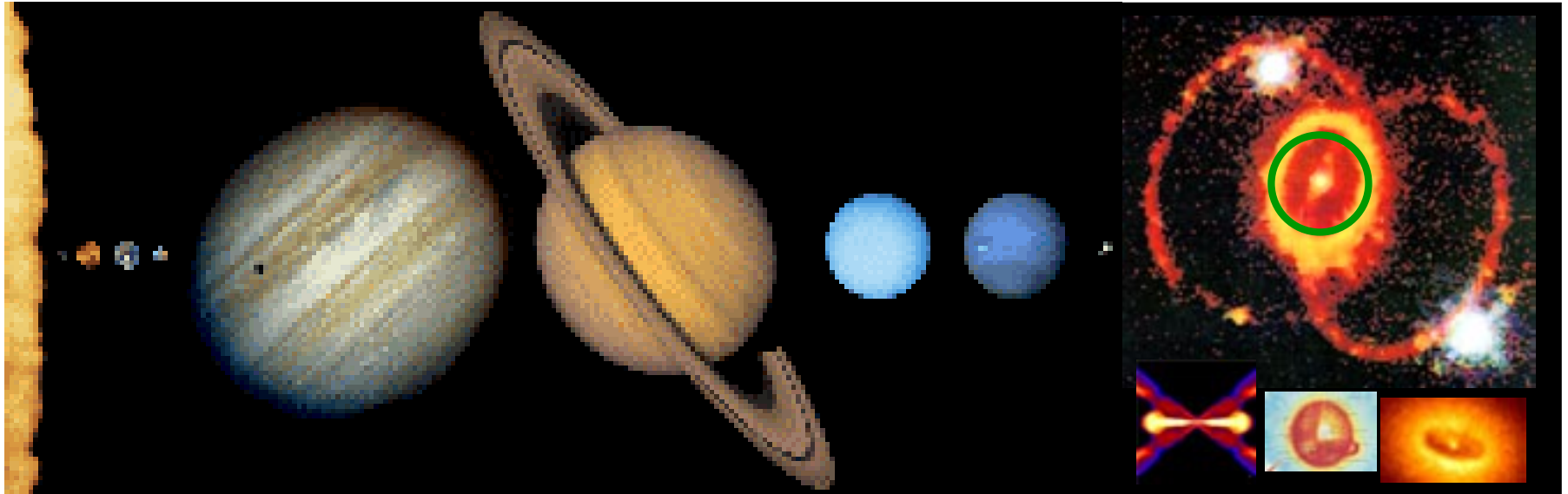


Rayevsky *et al*, IWPCTM, 1999.

# The tools are in place to perform quantitative experiments at high isentropic pressures



# Outline

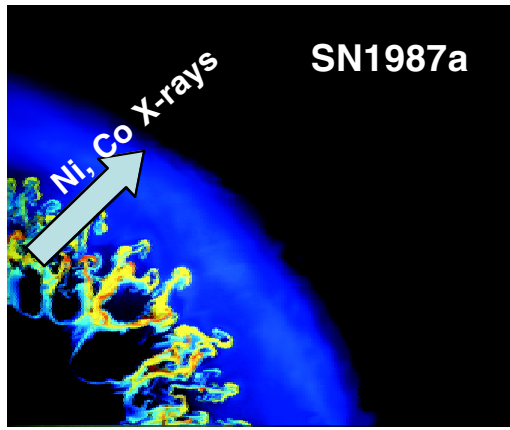


- **Significant advances in high energy density physics**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport and atomic physics*
- **Future directions**

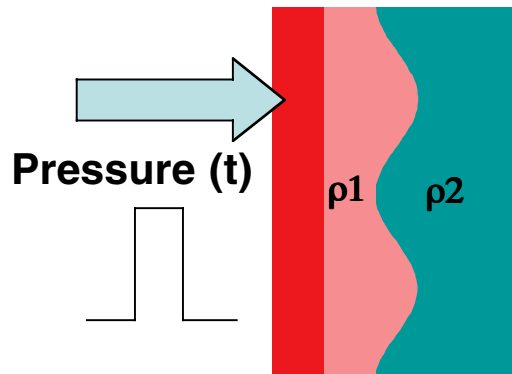
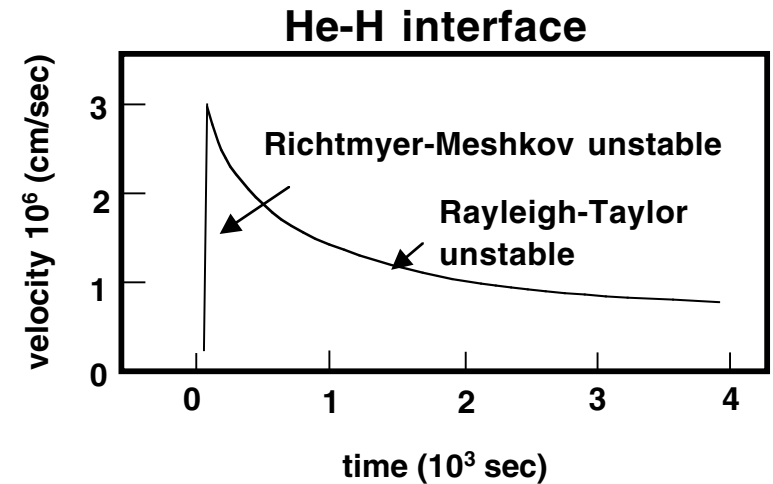
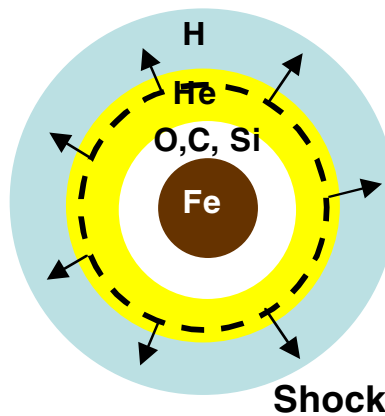
# Hydrodynamic instabilities may play a role in understanding data from supernova 1987a



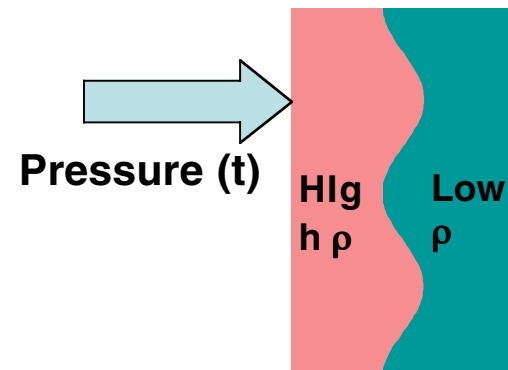
Numerical simulations of SN 1987A do not reproduce measurements



“Onion skin” model



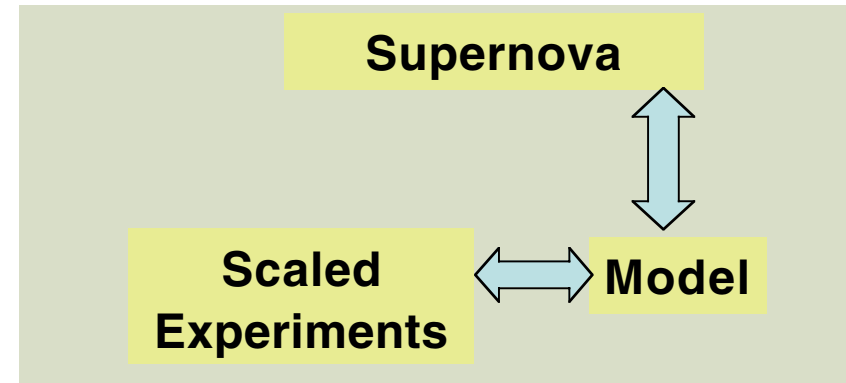
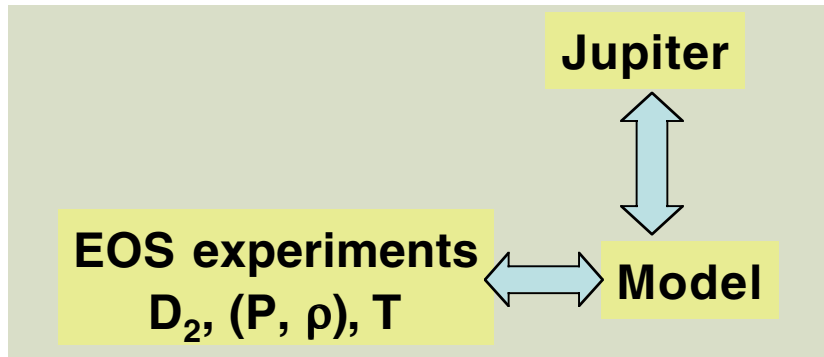
The Richtmyer-Meshkov instability occurs at an interface impulsively accelerated by a shock



The Rayleigh-Taylor instability occurs at an interface when a heavier fluid decelerates against a lighter one

How does one test complex dynamical phenomena?

# Hydrodynamics can be tested in dynamical models with scaled experimental testbeds



- The dynamical behavior of a system described by Eulers' Equations

$$\frac{\partial \mathbf{P}}{\partial t} - \gamma \frac{\mathbf{P}}{\rho} \frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{P} - \gamma \frac{\mathbf{P}}{\rho} \frac{\partial \rho}{\partial t} \mathbf{v} \cdot \nabla \rho = 0$$

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = - \frac{\nabla \mathbf{P}}{\rho} \quad \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

is invariant under any scale transformation that preserves

$$(h/\tau) \sqrt{(\rho/P)} \sim \text{Mach \#}$$

provided viscosity and thermal and radiative transport can be neglected

Hydrodynamical evolution is scaled for scale sizes  $\sim \rho / \nabla \rho$

# Practical implementation of scaling requires precise control over experimental conditions



1. Tailored pressure history

$$(h/\tau)\sqrt{(\rho/P)}$$

2. Scaled initial conditions

=> Solid materials easier

Geometry

3. Compressible materials

=> Ionized



Laser Beams

=> Pressure (t )

# Practical implementation of scaling requires precise control over experimental conditions

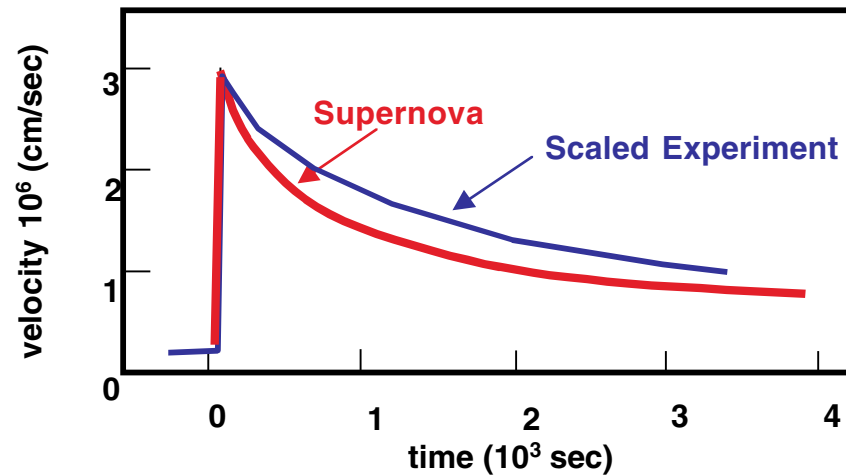
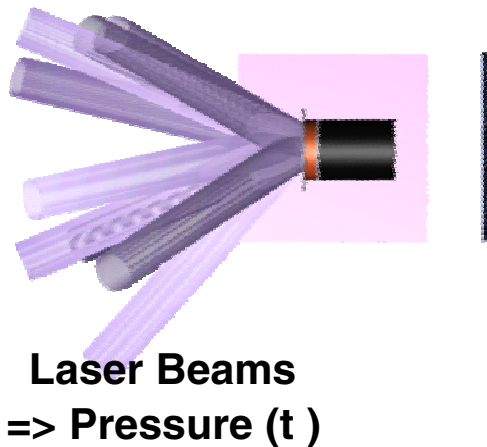


## 1. Tailored pressure history

2. Scaled initial conditions  
=> Solid materials easier  
Geometry

$$(h/\tau)\sqrt{\rho/P}$$

3. Compressible materials  
=> Ionized



# Practical implementation of scaling requires precise control over experimental conditions

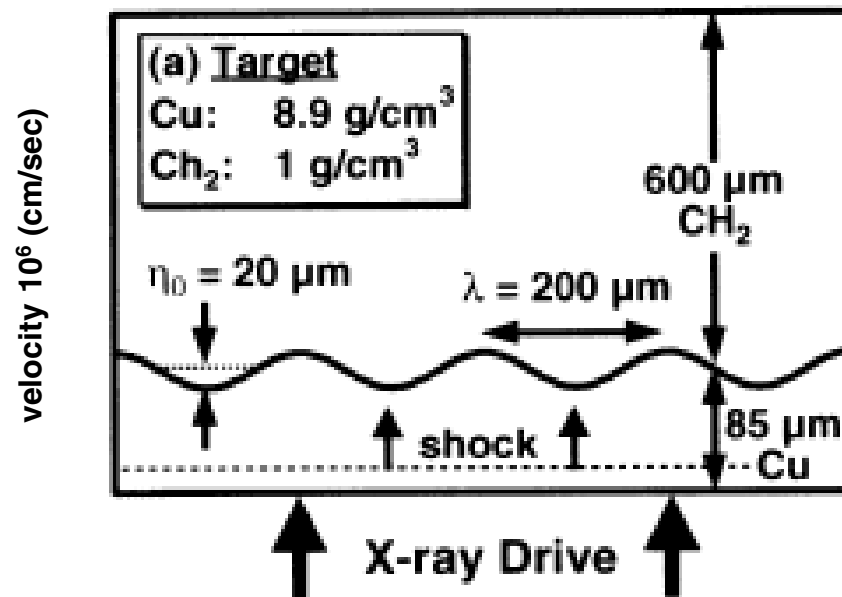
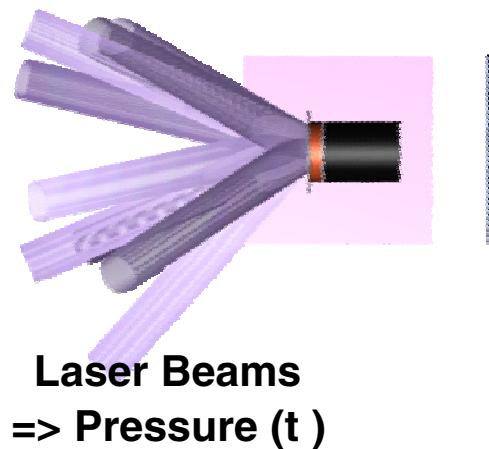


## 1. Tailored pressure history

2. Scaled initial conditions  
=> Solid materials easier  
Geometry

$$(h/\tau)\sqrt{(\rho/P)}$$

3. Compressible materials  
=> Ionized





# Practical implementation of scaling requires precise control over experimental conditions

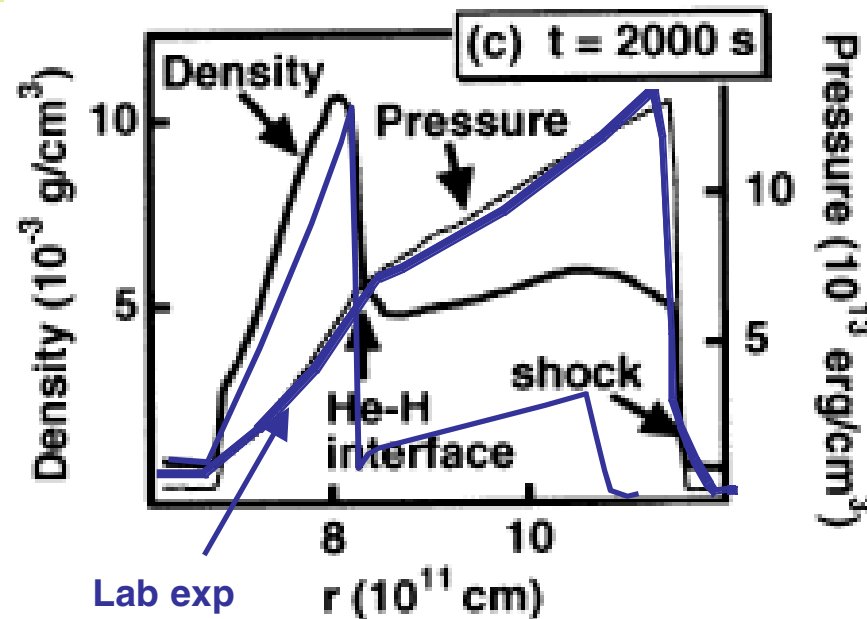
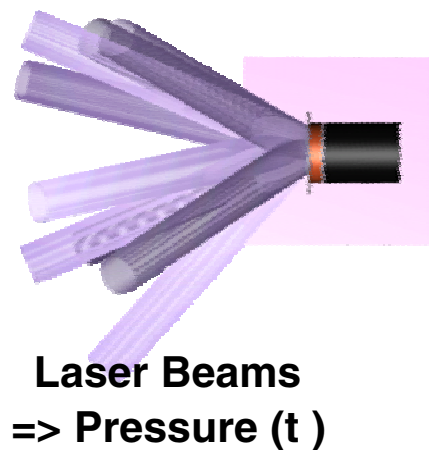


## 1. Tailored pressure history

2. Scaled initial conditions  
=> Solid materials easier  
Geometry

$$(h/\tau)\sqrt{\rho/P}$$

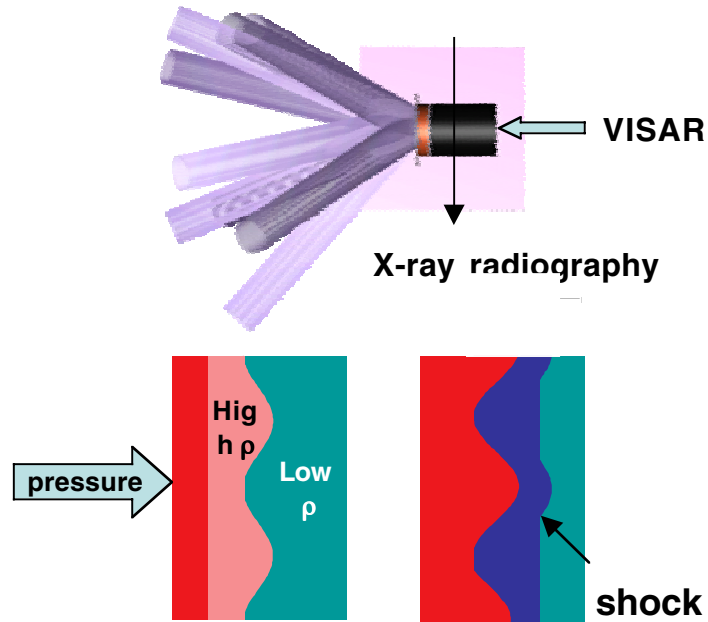
3. Compressible materials  
=> Ionized



# High Mach number RM experiments measured reduced growth due to shock proximity



Laser driven shock tube Mach  $\sim 10$

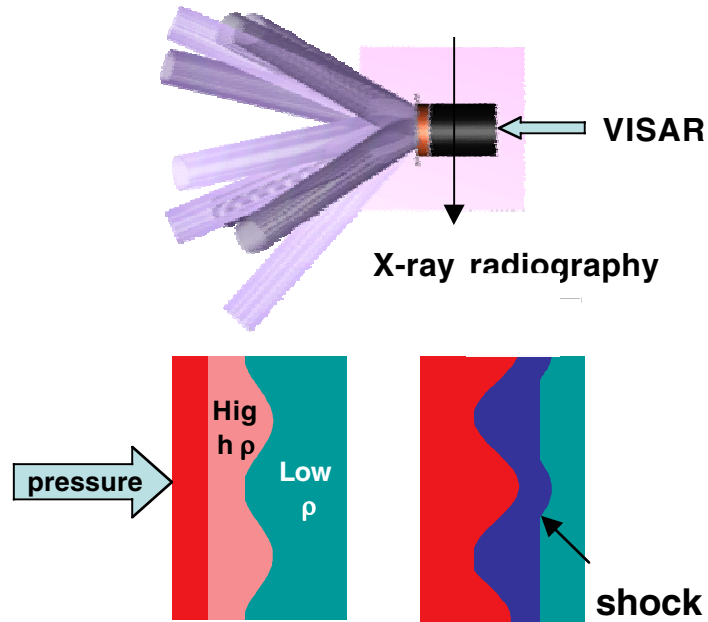


- Incompressible models (Sadot, 1998) predict that spike tip moves faster than shock
- Precise experiments can be done on laser driven shock tubes and compared to analytical models

# High Mach number RM experiments measured reduced growth due to shock proximity



Laser driven shock tube Mach  $\sim 10$

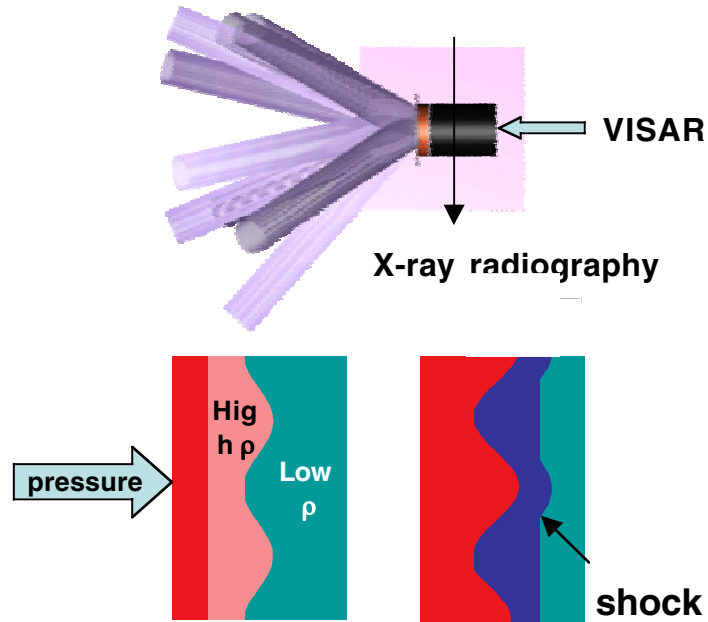


- Incompressible models (Sadot, 1998) predict that spike tip moves faster than shock
- Precise experiments can be done on laser driven shock tubes and compared to analytical models

# High Mach number RM experiments measured reduced growth due to shock proximity



Laser driven shock tube Mach  $\sim 10$

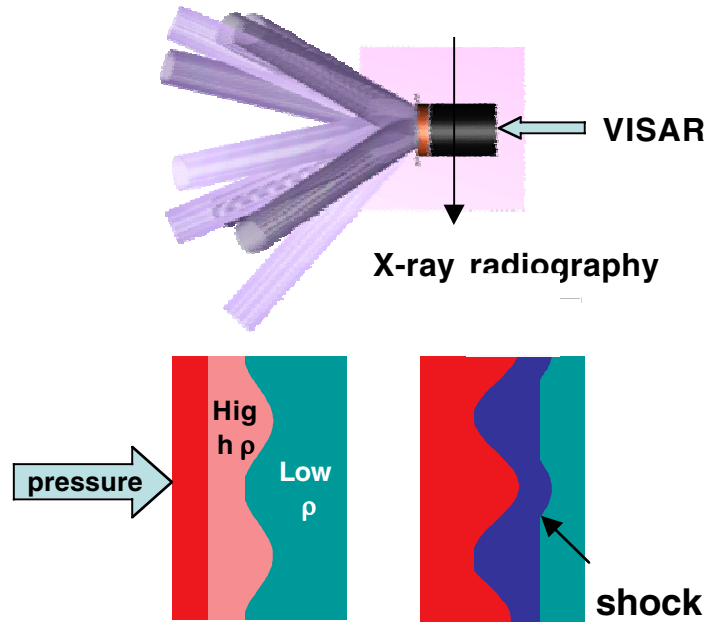


- Incompressible models (Sadot, 1998) predict that spike tip moves faster than shock
- Precise experiments can be done on laser driven shock tubes and compared to analytical models

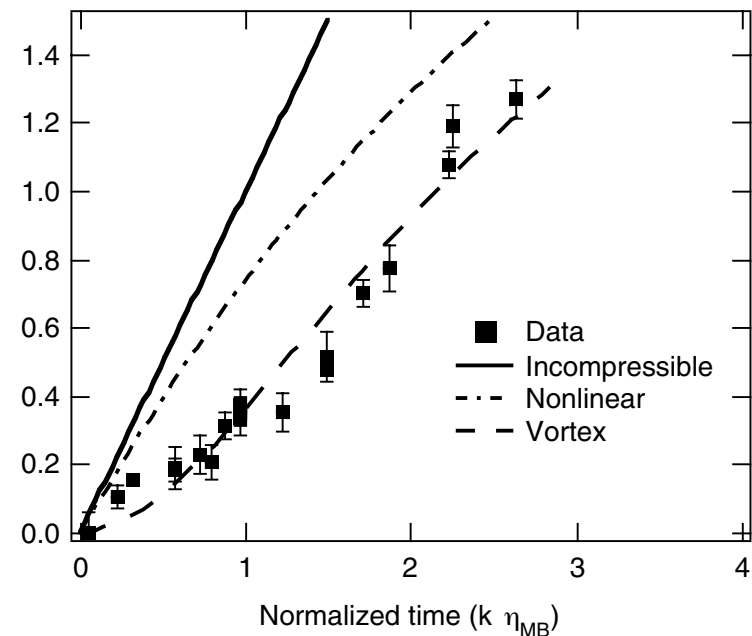
# High Mach number RM experiments measured reduced growth due to shock proximity



Laser driven shock tube Mach  $\sim 10$



- Incompressible models (Sadot, 1998) predict that spike tip moves faster than shock
- Precise experiments can be done on laser driven shock tubes and compared to analytical models

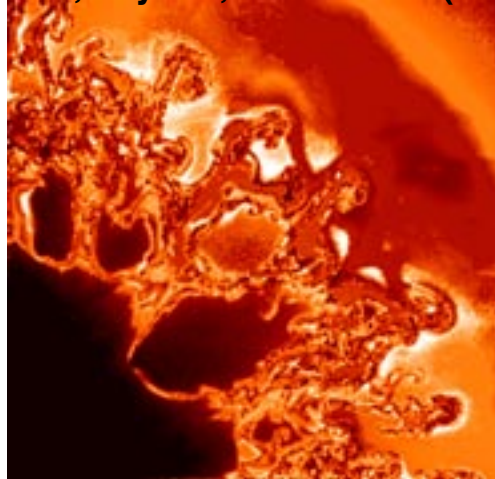


# Hydrodynamically scaled SN instability experiments have been performed on Omega

---



2D simulation of SN1987A  
Muller, Fryxell, and Arnett (1991)



**Physics was piecewise separated and scaled experiments were compared to modeling**

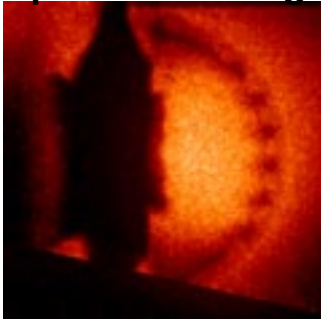
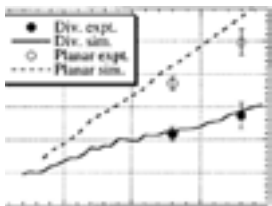
Remington, B. A., R. P. Drake, et al. (2000). Physics of Plasmas 7(5 PT2): 1641-1652.

Robey, H. F., et al. (2001). Physics of Plasmas 8(5 PT2): 2446-2453. Kane, J., et al. (2000). ApJS 127(2): 365-369.

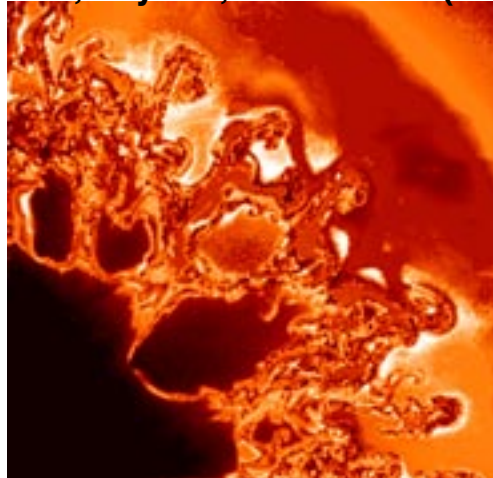
# Hydrodynamically scaled SN instability experiments have been performed on Omega



Spherical divergence



2D simulation of SN1987A  
Muller, Fryxell, and Arnett (1991)



**Physics was piecewise separated and scaled experiments were compared to modeling**

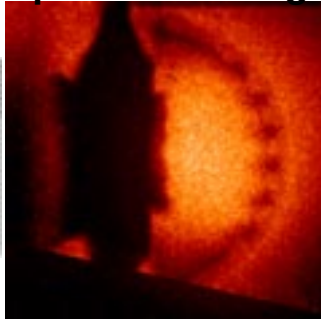
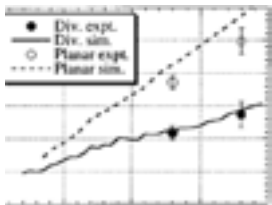
Remington, B. A., R. P. Drake, et al. (2000). *Physics of Plasmas* 7(5 PT2): 1641-1652.

Robey, H. F., et al. (2001). *Physics of Plasmas* 8(5 PT2): 2446-2453. Kane, J., et al. (2000). *ApJS* 127(2): 365-369.

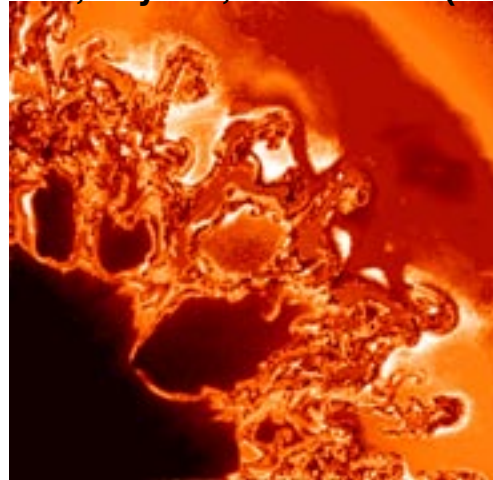
# Hydrodynamically scaled SN instability experiments have been performed on Omega



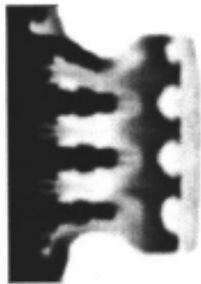
Spherical divergence



2D simulation of SN1987A  
Muller, Fryxell, and Arnett (1991)



Multi-interface coupling



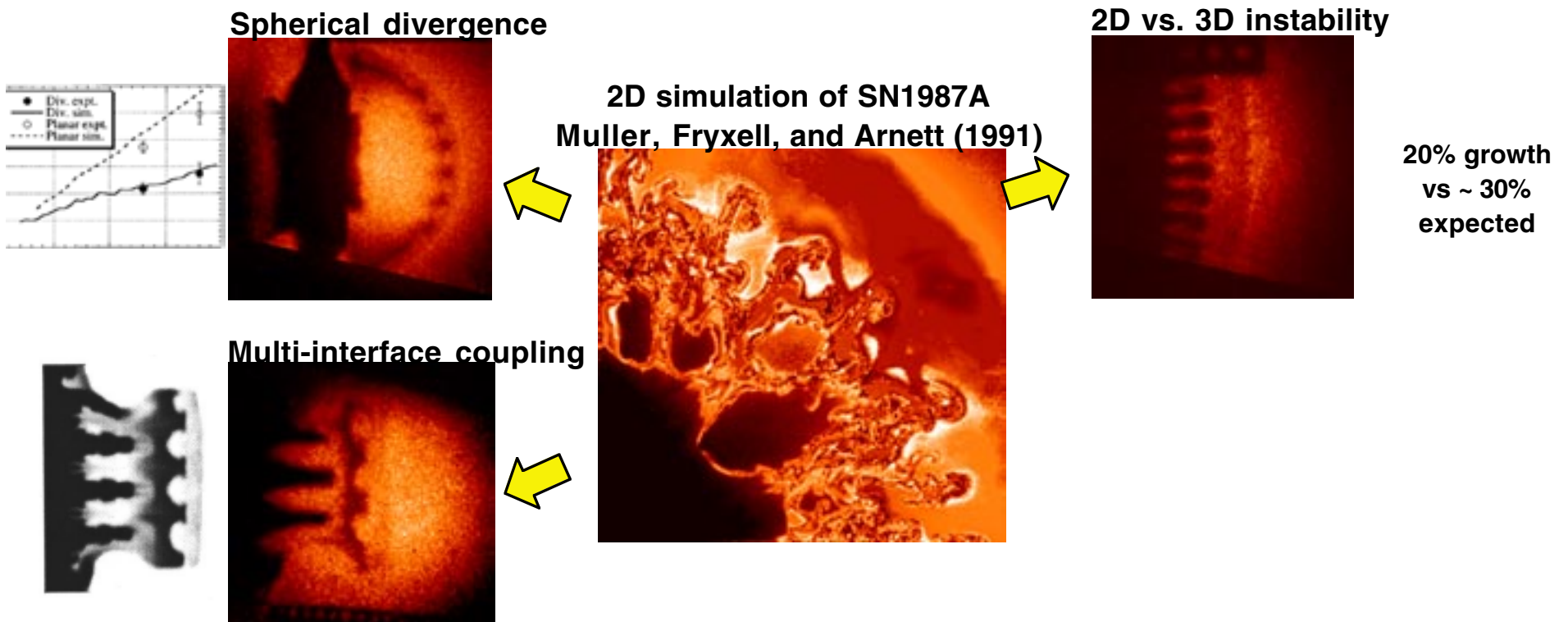
**Physics was piecewise separated and scaled experiments were compared to modeling**

Remington, B. A., R. P. Drake, et al. (2000). *Physics of Plasmas* 7(5 PT2): 1641-1652.

Robey, H. F., et al. (2001). *Physics of Plasmas* 8(5 PT2): 2446-2453. Kane, J., et al. (2000). *ApJS* 127(2): 365-369.



# Hydrodynamically scaled SN instability experiments have been performed on Omega

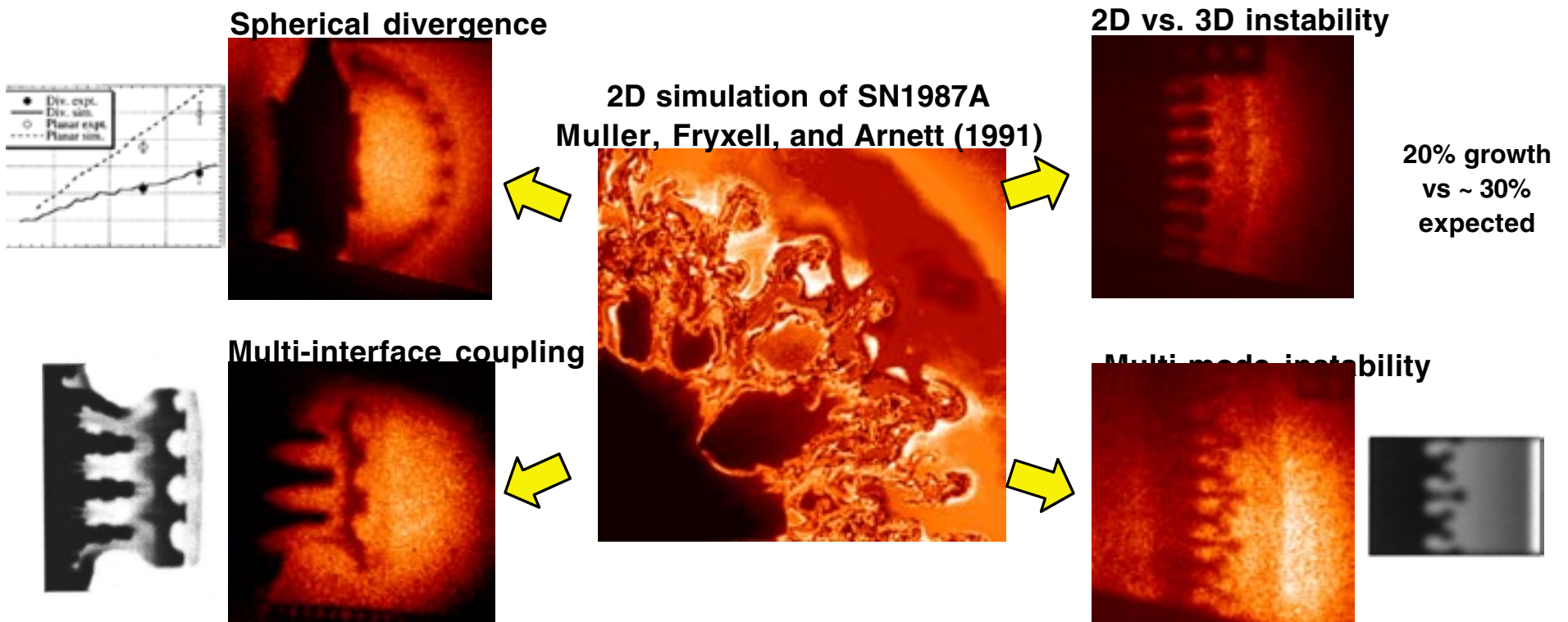


**Physics was piecewise separated and scaled experiments were compared to modeling**

Remington, B. A., R. P. Drake, et al. (2000). *Physics of Plasmas* 7(5 PT2): 1641-1652.

Robey, H. F., et al. (2001). *Physics of Plasmas* 8(5 PT2): 2446-2453. Kane, J., et al. (2000). *ApJS* 127(2): 365-369.

# Hydrodynamically scaled SN instability experiments have been performed on Omega



**Physics was piecewise separated and scaled experiments were compared to modeling**

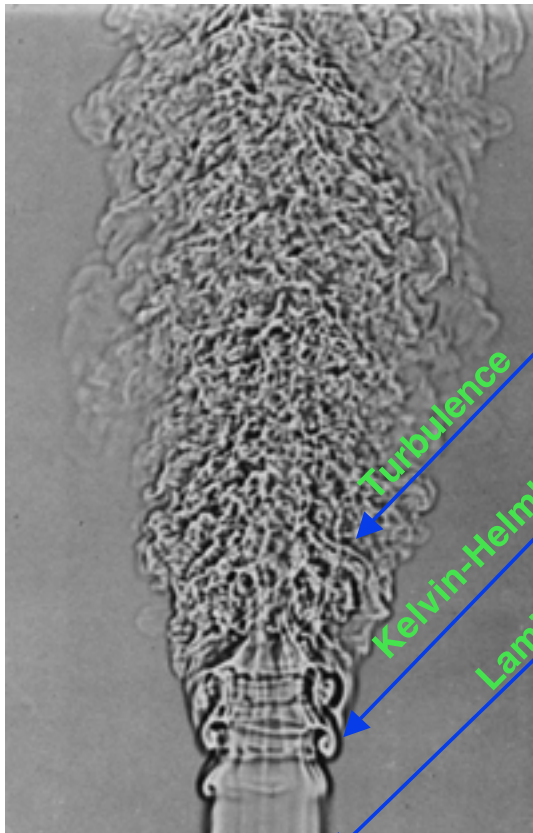
Remington, B. A., R. P. Drake, et al. (2000). *Physics of Plasmas* 7(5 PT2): 1641-1652.

Robey, H. F., et al. (2001). *Physics of Plasmas* 8(5 PT2): 2446-2453. Kane, J., et al. (2000). *ApJS* 127(2): 365-369.

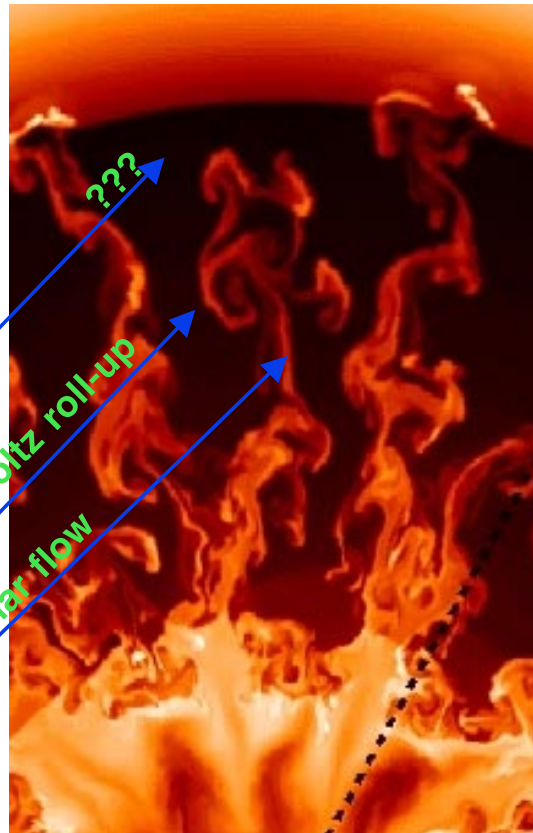
# Large scale physics can be tested, but what about small scales?



Simulations of supernova explosions do not appear to be turbulent



Data of a  
 $Re = 10^4$  flow



SN simulations of  
 $Re = 10^{10}$  hydrodynamics

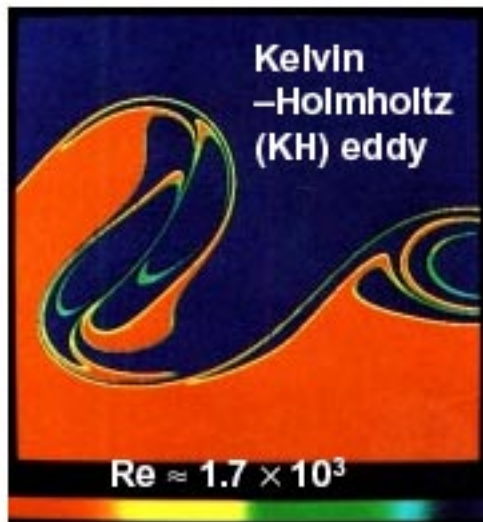
$Re = \text{inertial/viscous} = Lv/\text{viscosity}$

- Turbulence will affect mixing
  - Linear  $RM \sim t$
  - Turbulent  $RM \sim t^{5-1}$

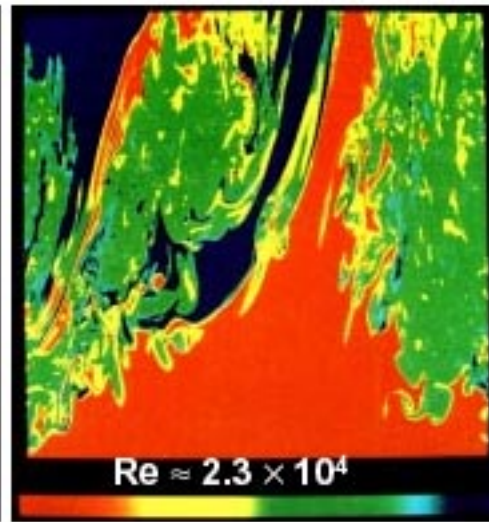
**A universal “mixing transition” at Reynolds number  $\sim 2 \times 10^4$  was proposed - physics is same afterwards**



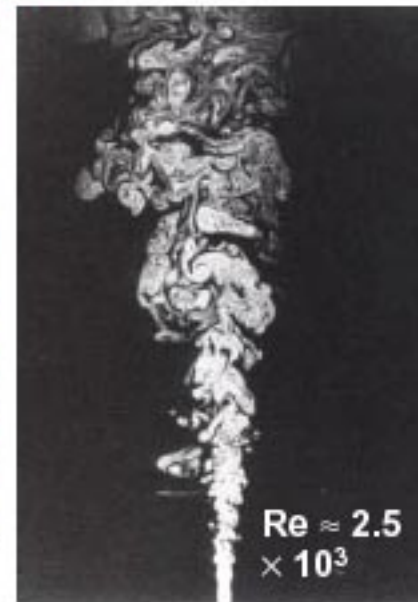
**“Steady” State**



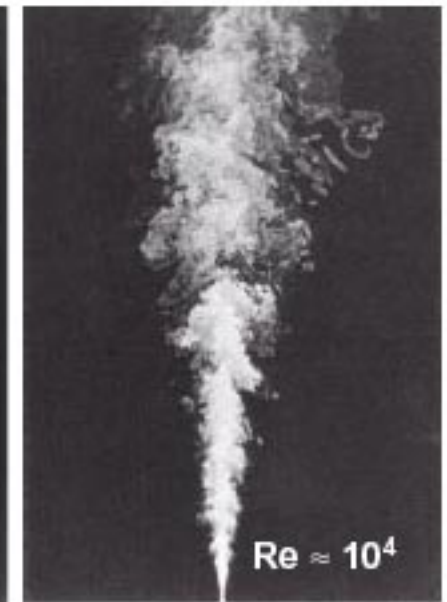
Very nonlinear, but not atomically mixed



Turbulent, atomically mixed



Turbulent, but not atomically mixed



Turbulent, atomically mixed

**For dynamical systems, it is proposed to generate the scales needed,  
time  $> (L/v) \times 100 \text{ Re}^{-.5}$**

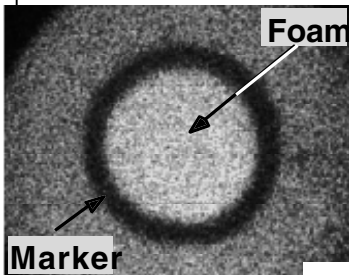
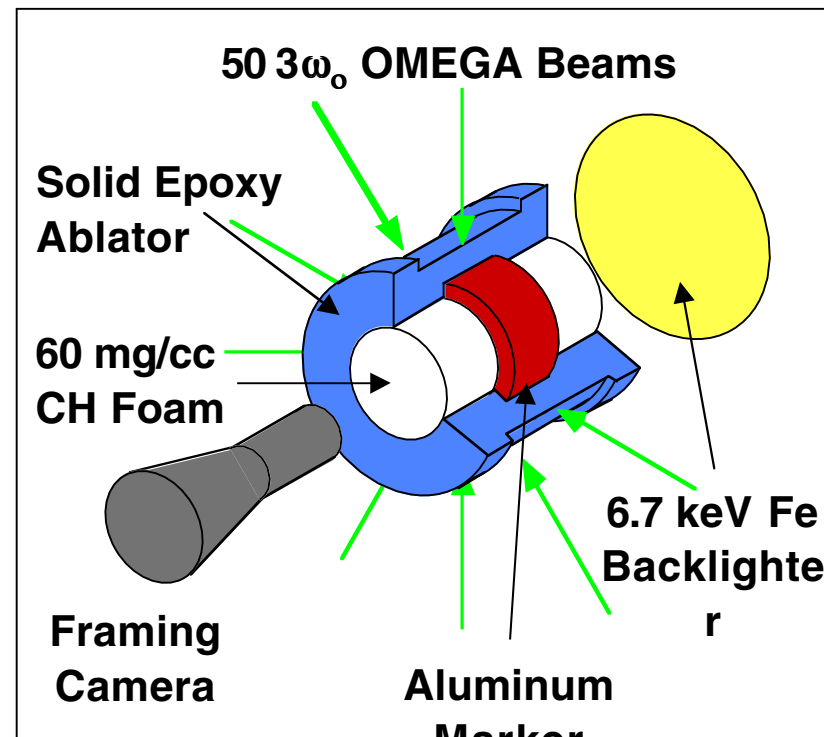
**Driving, diagnosing and modeling the transition to turbulence  
in compressible high Reynolds number flows is the next  
challenge**



# Measurements of turbulent instability growth have been performed in planar and convergent geometries



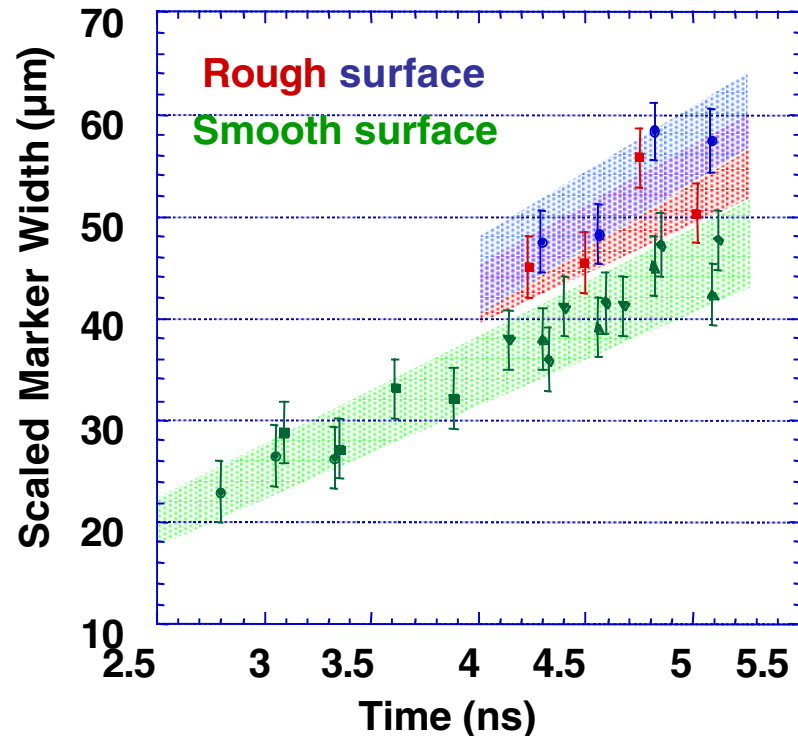
- Convergent experiments in cylindrical geometry



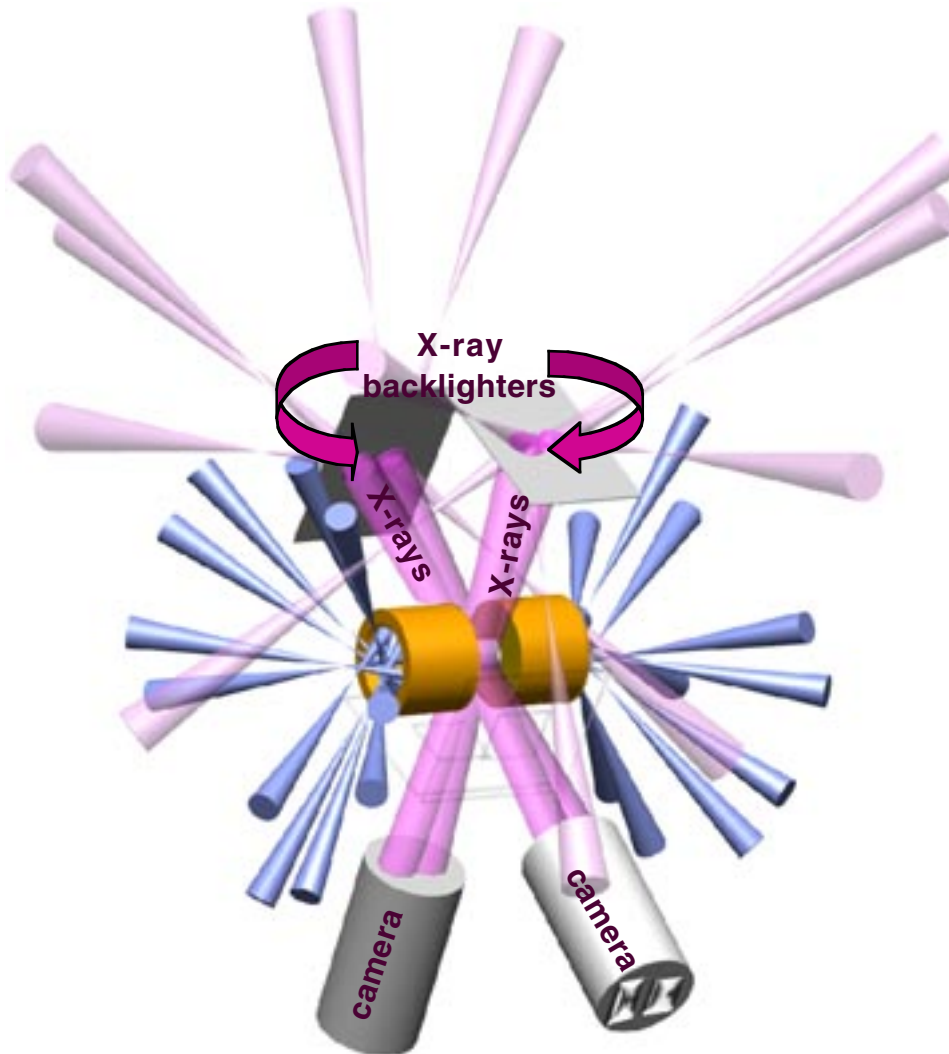
Planar geometry:

Turbulent RT, RM experiments: accuracy to 20 - 50% in model parameters

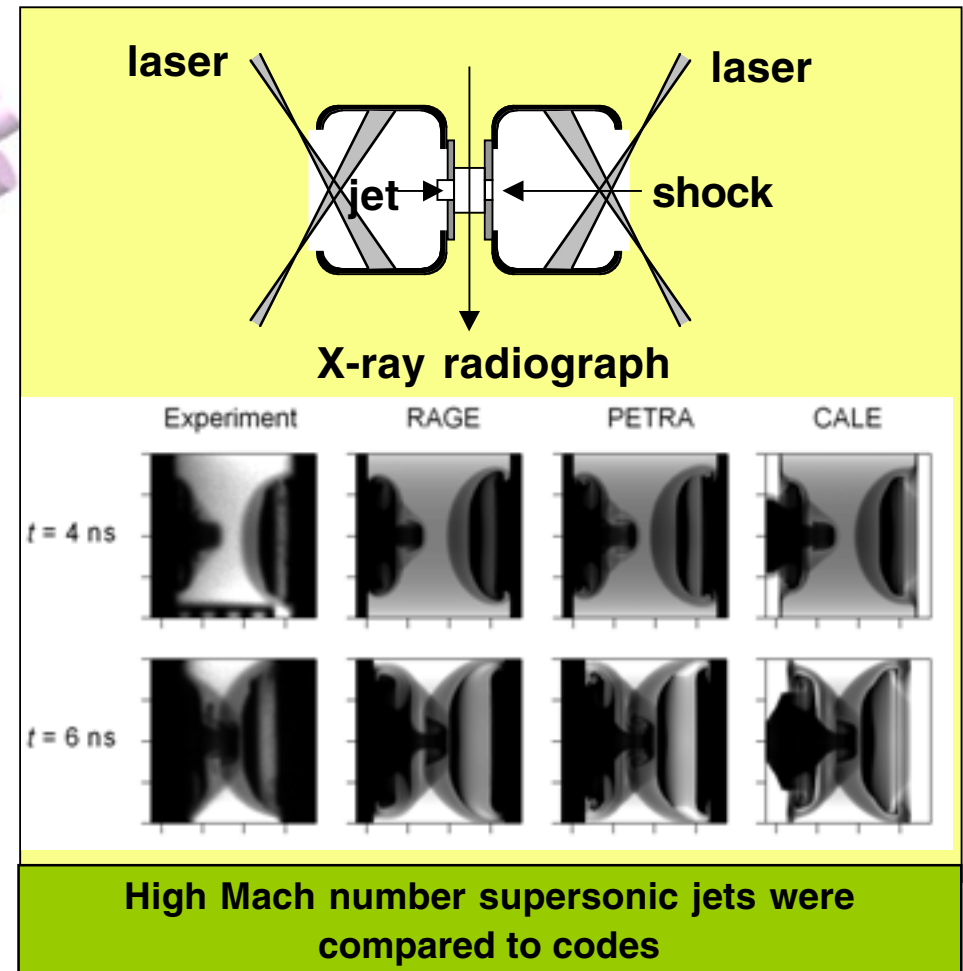
**New facilities will allow higher accuracies and more detailed measurements of turbulent growth**



# These testbeds for large and mid-scale size hydrodynamics are well established



Radiographic images at different views in space and time



Experiments have started on Z, where target sizes are  $O(10\times)$  larger)

# Outline



- **Significant advances in high energy density physics**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport and atomic physics*
- **Future directions**

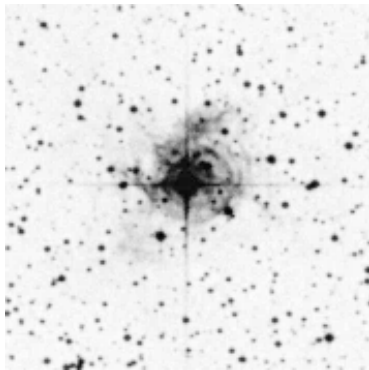
## Radiation

**Interaction of intense soft X-rays with matter is important in a host of astrophysical phenomena**

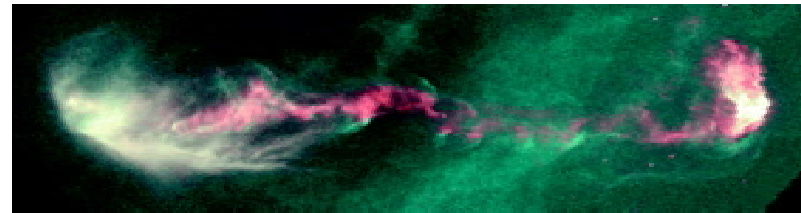
---



**Radiation penetration from a supernova into interstellar medium - ionization fronts**



**Cepheid variables - opacities of stellar envelopes**



**Formation of astrophysical jets - radiative cooling**



**Radiation from accretion material in X-ray binaries - photoionized plasmas**



# When is radiation important at high energy density conditions?



## Local Energy Density:

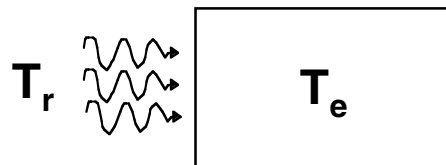
$$\text{Radiation energy density} \sim \frac{4\sigma T_R^4}{c} \quad 1 \text{ Mbar} \Rightarrow T_R = 300 \text{ eV}$$

$$\text{But material energy dominates} \quad \frac{\text{Radiation Energy Density}}{\text{Material Energy Density}} \sim \frac{T^3}{\rho (Z/A)}$$

$$\text{At } T_R = 300 \text{ eV, } 1\text{g/cc} \Rightarrow \sim 10^{-3}$$

$$\text{Or } T \sim 2 \text{ keV}$$

## Energy Transport:

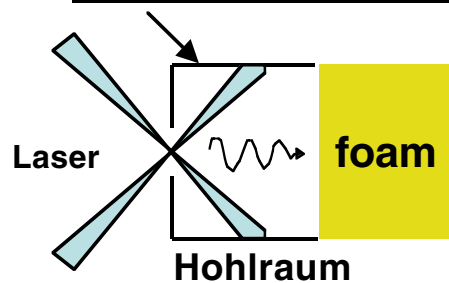


$$\frac{\text{Radiation Energy Flux}}{\text{Material Energy Flux}} = \frac{\text{Radiation Energy Density}}{\text{Material Energy Density}} \frac{c}{v}$$

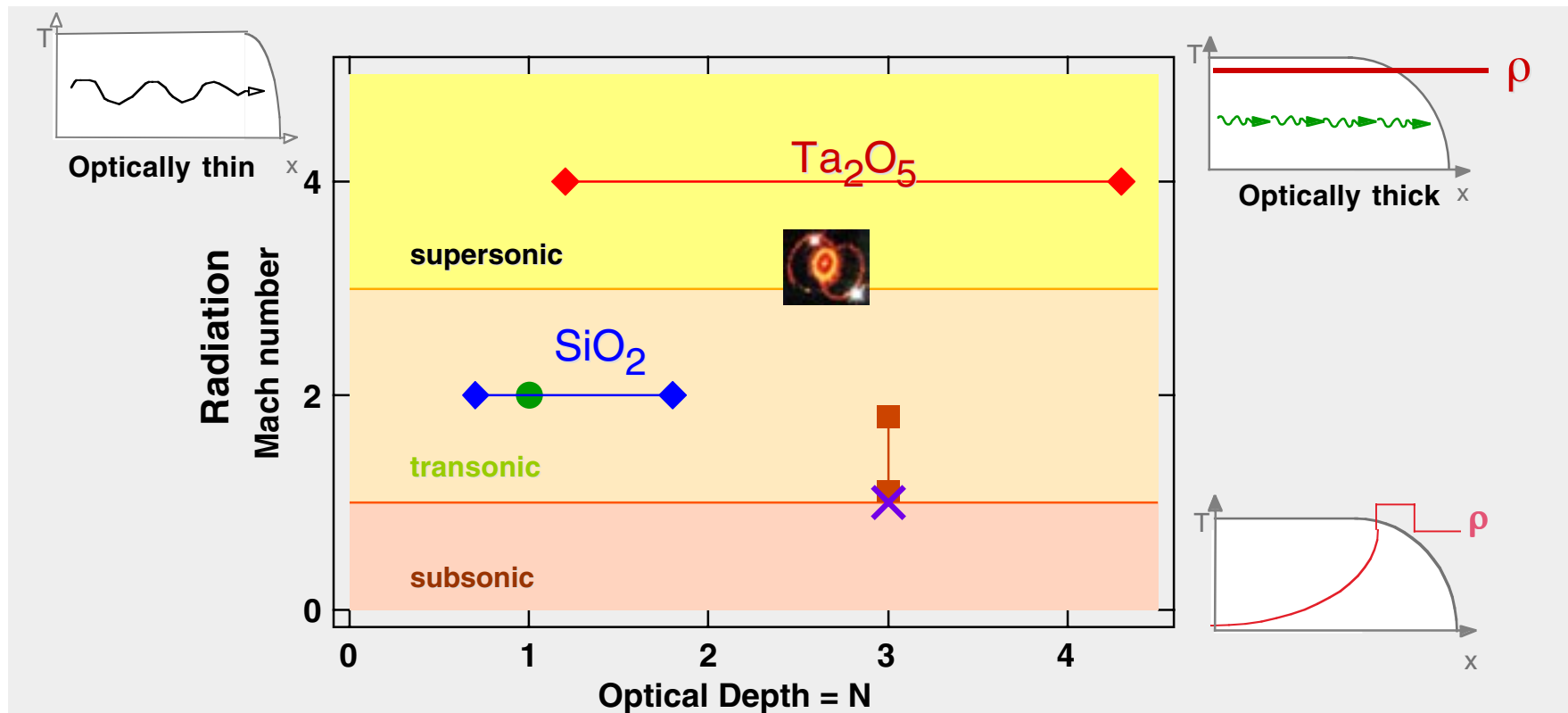
Important in most stellar atmospheres: sun  $\sim 5 \times 10^4$

$$\text{At } T = 300 \text{ eV, } 1\text{g/cc, } v = \text{sound speed} \Rightarrow \sim 2$$

# Supersonic ionization fronts have been created in the laboratory



$$\frac{\text{Radiative energy flux}}{\text{Material energy flux}} = \frac{\sigma T^4}{\varepsilon \rho c_s} \propto \left( \frac{x}{mfp_R} \right) \left( \frac{x}{c_s} \right) = NM$$



Back, C. A., et al. (2000). *Physics of Plasmas* 7(5 PT2): 2126-2134.

Afsharrad, T., (1994). *Physical Review Letters* 73(1): 74-77.

Hoarty, D., et al. (1999). *Physics of Plasmas* 6(5 PT2): 2171-2177.

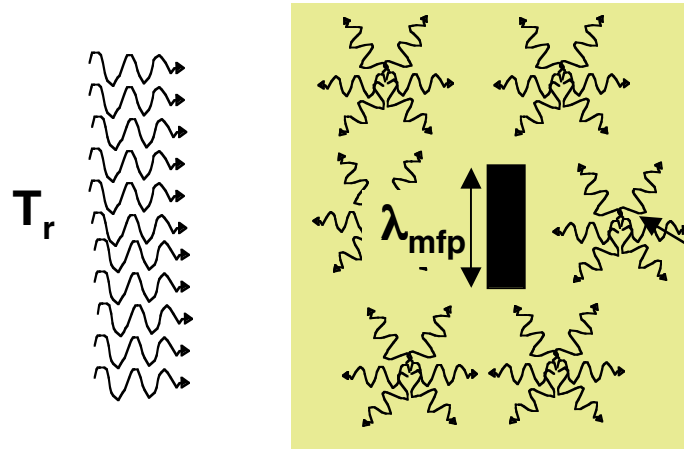
Massen, J., (1994). *Physical Review E* 50(6): 5130-5133.

Bozier, J.C. et. al., (1986) *Phys Rev Lett*, 57: 1304.

# Several techniques are available to solve radiation transport problems



- **Diffusion** assumes radiation  $\sim$  isotropic

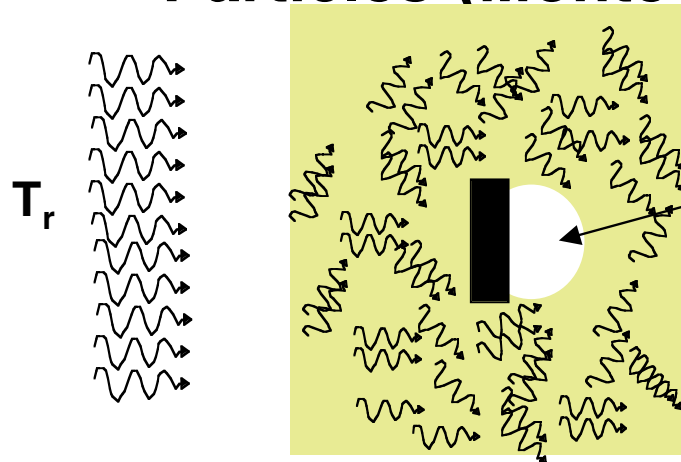


requires large optical depth  
e.g deep inside stars

Radiation from a far away Supernova  
illuminating a protostar is non-isotropic

Radiation transports behind a wall

- **Particles (Monte Carlo)**



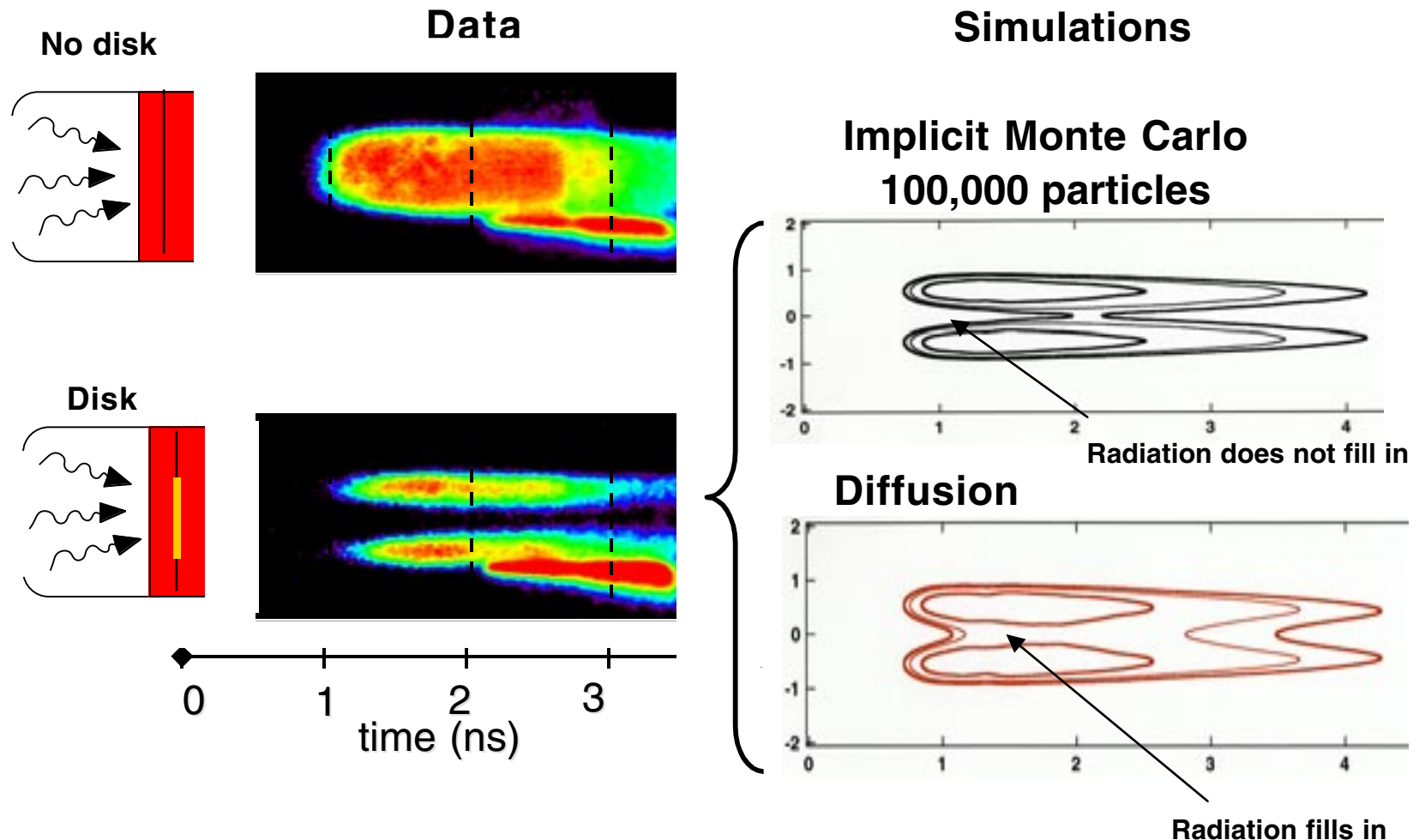
Photons are tracked as particles

Radiation does not transport behind a wall

Signal/Noise  $\sim$  ( $\#$  of particles/cell) $^{1/2}$

=> big computers

# Data from experiments can now quantitatively evaluate radiation transport methods

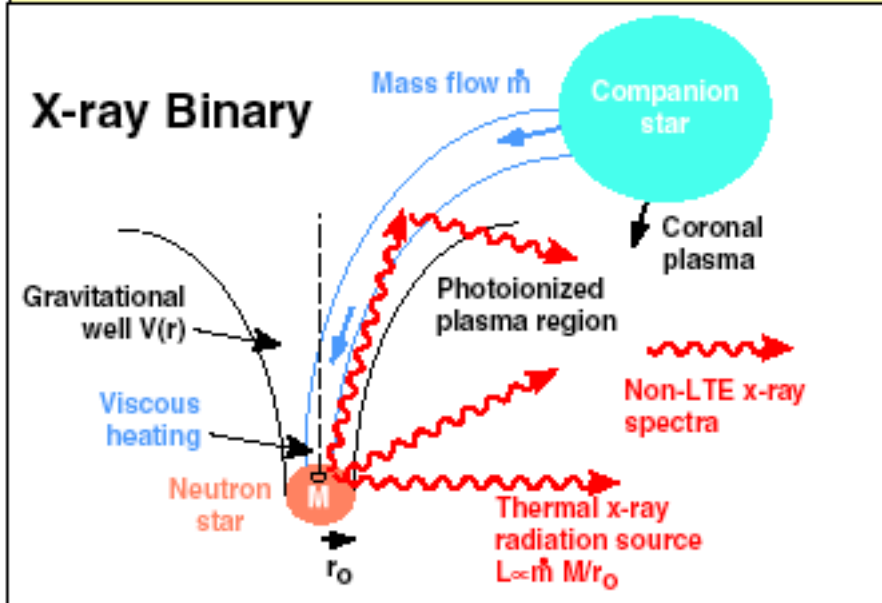


**Advanced radiation transport algorithms like Implicit Monte Carlo agrees better than diffusion in Omega experiments**

# X-ray astrophysics now needs a detailed understanding of photoionized plasmas



Gravitational energy of infalling matter  $\Rightarrow$  X-rays

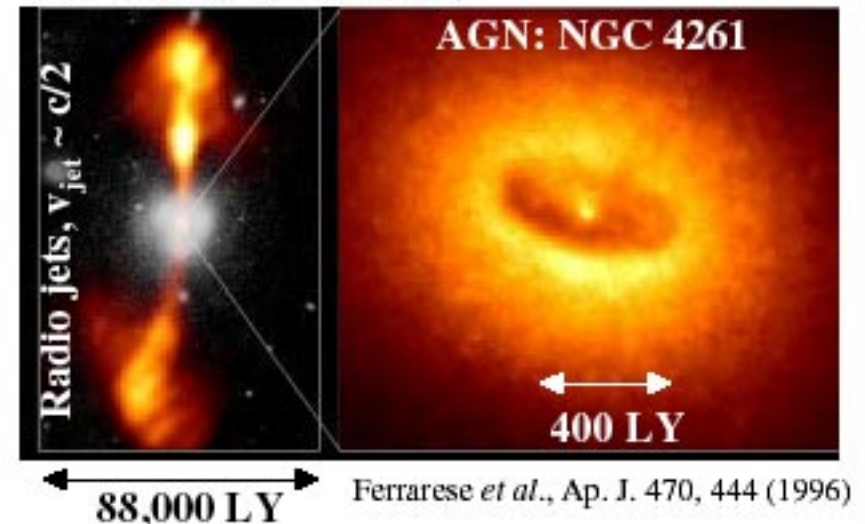


Half of all X-ray sources are accretion-powered:

- Active Galactic Nuclei (AGNs)
- Cataclysmic Variables (CVs)
- Black Hole or Neutron Star X-ray Binaries

Copious X-rays photoionize the accretion disk

Piner *et al.*, A.J. 122, 2954 (2001)

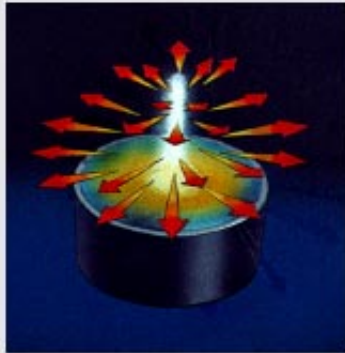


High radiation fields, low density so 3 body recombination

$$\xi \sim 4\pi I/n \sim \text{Irradiance/electron density}$$

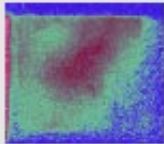
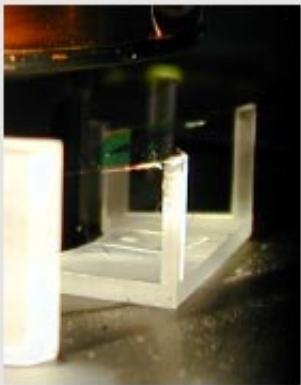
$\xi \gg 1$  for regimes of interest

# **Radiation** Understanding atomic physics in photoionized plasmas allows interpretation of astrophysical spectra



Z is Radiation Source

80 TW, 160 eV  
Quasi-blackbody w/tail  
Cylinder 1.0 cm high,  
0.15-0.2 cm diameter



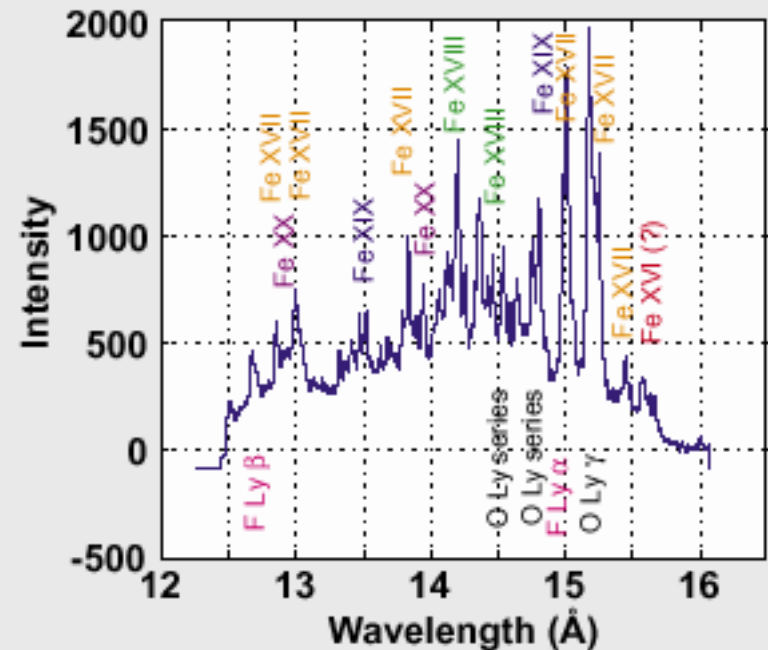
## Photoionized Sample

1.6 cm from center of source

3 Fe : 2 NaF mixture

$$n_{\text{Fe}} = 8 \times 10^{17} / \text{cm}^3$$
$$n_e \approx 29 n_{Fe} \approx 2.5 \times 10^{19} / \text{cm}^3$$

### Emission Spectrum of a Plasma at Conditions of a Photoionized Accretion Disk



RSI 72, 1227 (2001)

**Te = 30 eV**

**Z = 16.5 ± .5 in comparison with  
6 models**

**The first experimental benchmark for X-ray photoionized plasma models is available,  $\xi \sim 20$**

# Universities worldwide are engaged in research in high energy density physics

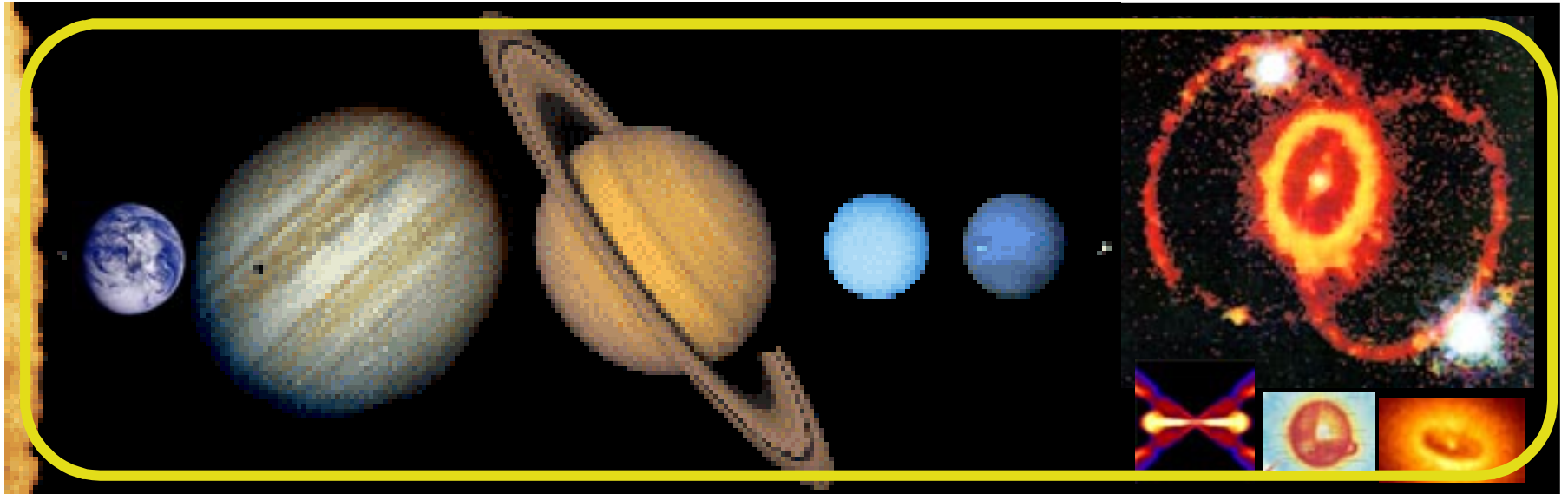


University	Research area
Berkeley	Precompressed EOS, X-ray diffraction
UCSD	Materials dynamics
U of Michigan	Supernova remnants, radiation transport
Cornell	High density plasmas
Univ of Rochester	Astrophysical jets, EOS
U of Colorado	X-ray interferometry, laboratory astrophysics
U of Reno	Z-pinch plasmas
U of Arizona	Laboratory astrophysics
Rice University	Radiative blast waves
U of Maryland	Laboratory astrophysics
UC Davis	Radiative blast waves
MIT	Nuclear physics
George Mason	Supernova hydrodynamics
U of Texas	Material Dynamics
U of Wisconsin	Atomic physics

University	Research area
Ecole Polytechnique	Opacity, EOS of water
Univ of British Columbia	EOS of foams
Osaka University	Laboratory astrophysics
Imperial College	Radiation transport, Z-pinches
University of Milan-Biocca	EOS of Au
Univ of Essex	EOS of foams
Oxford University	Material Dynamics
Univ Milan	EOS water, foams



# Outline

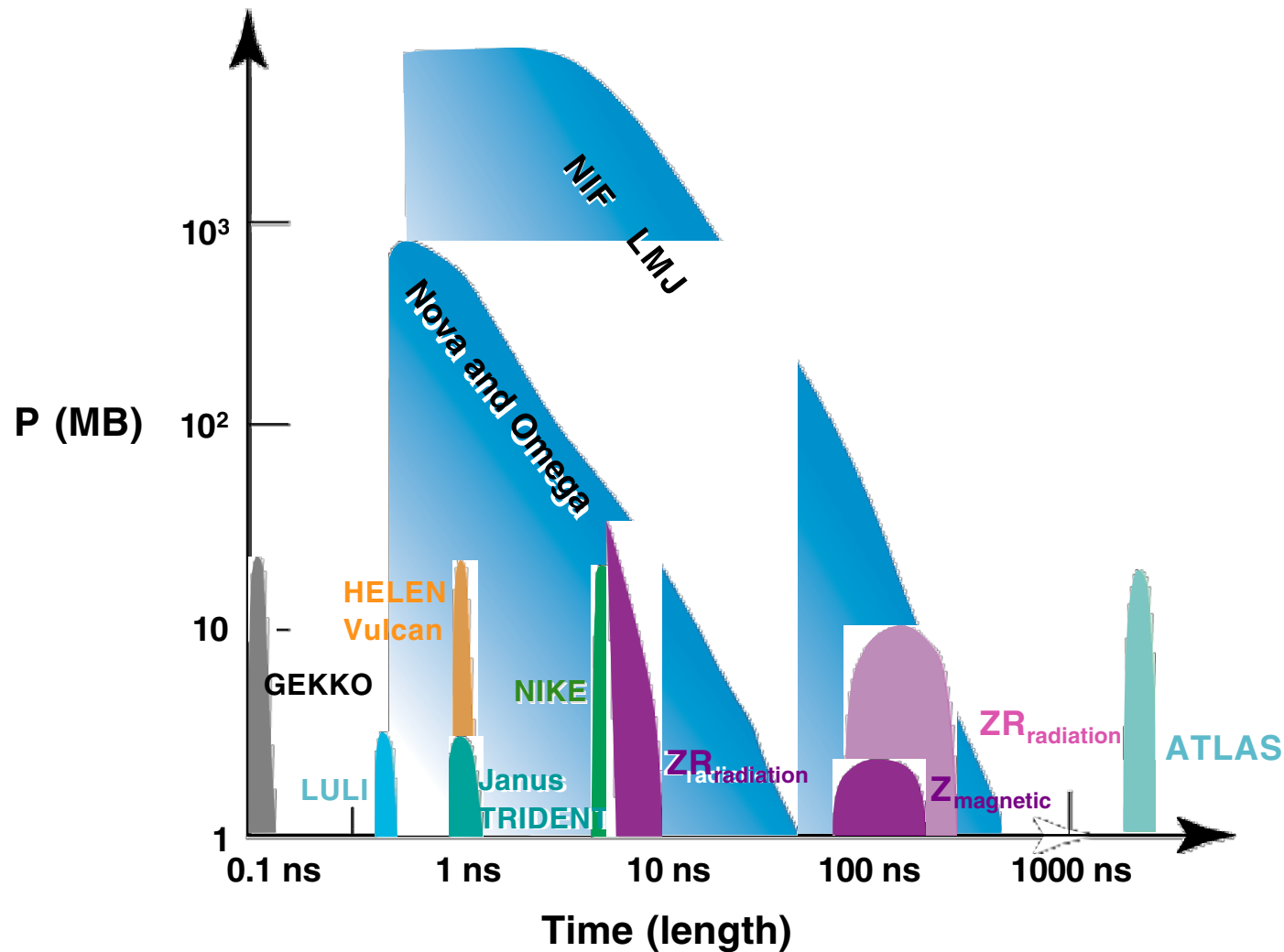


- **Significant advances in high energy density physics**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport and atomic physics*
- **Future directions**



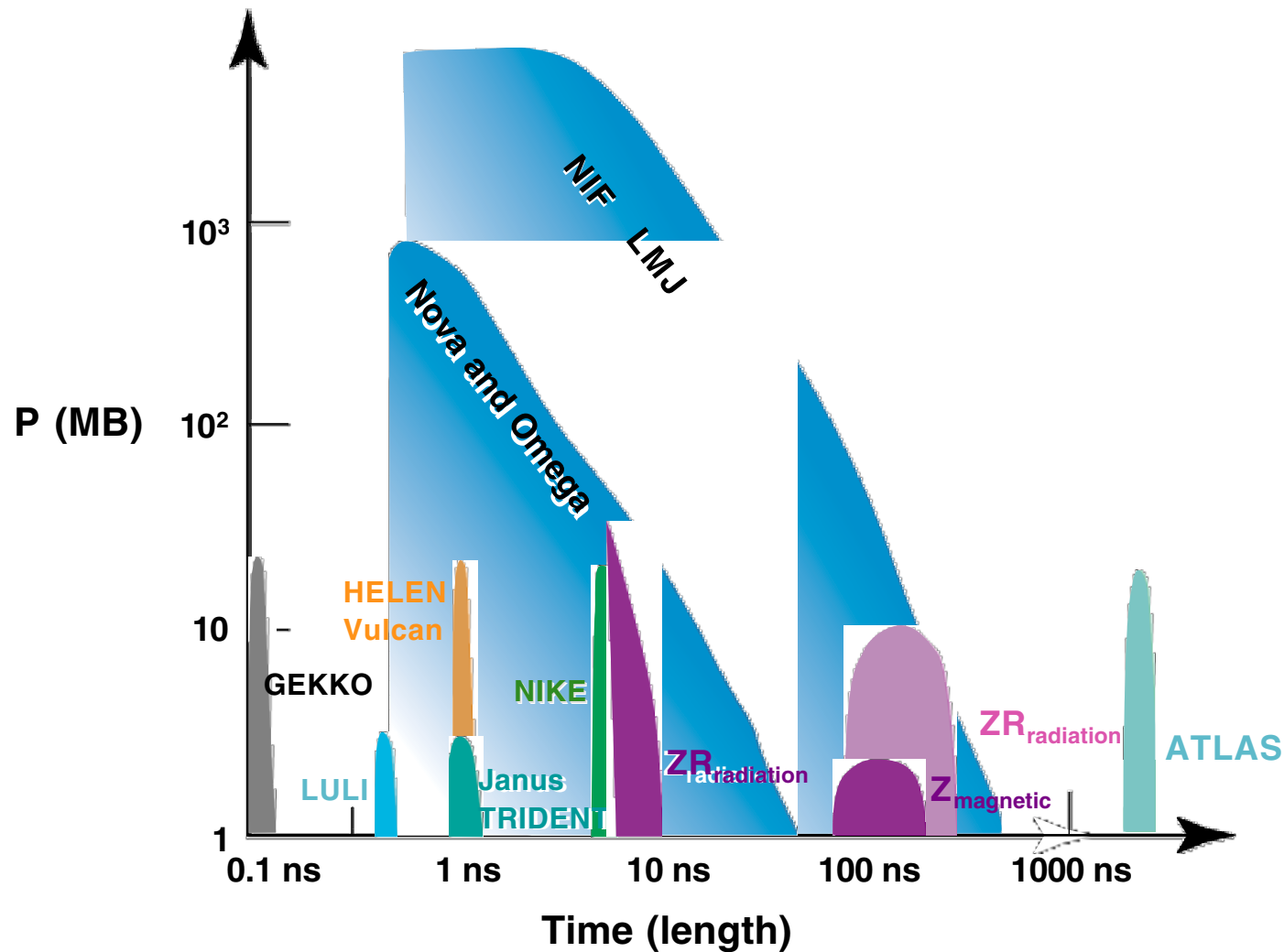
Future

# New facilities will expand high energy density research



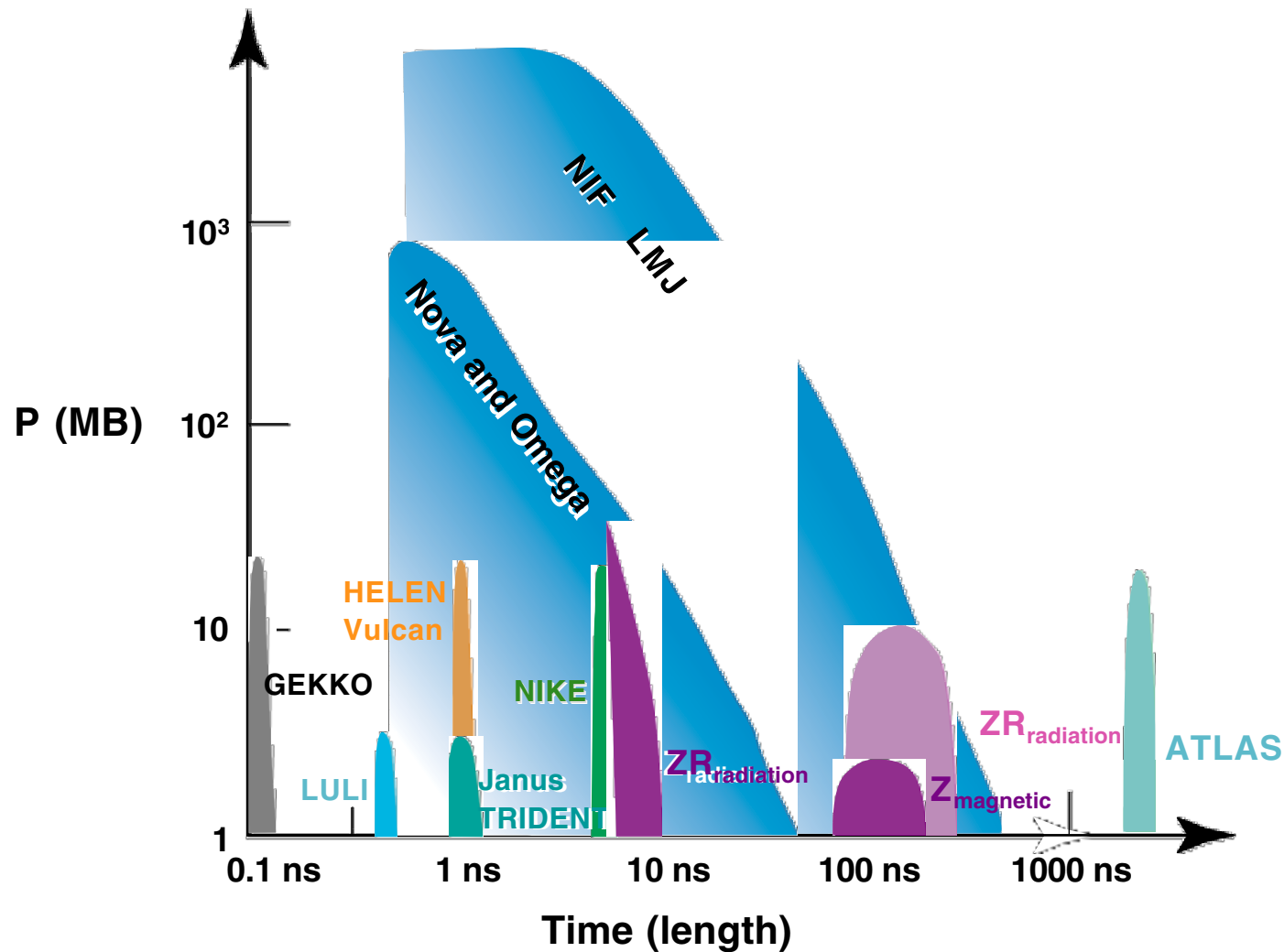
Future

# New facilities will expand high energy density research



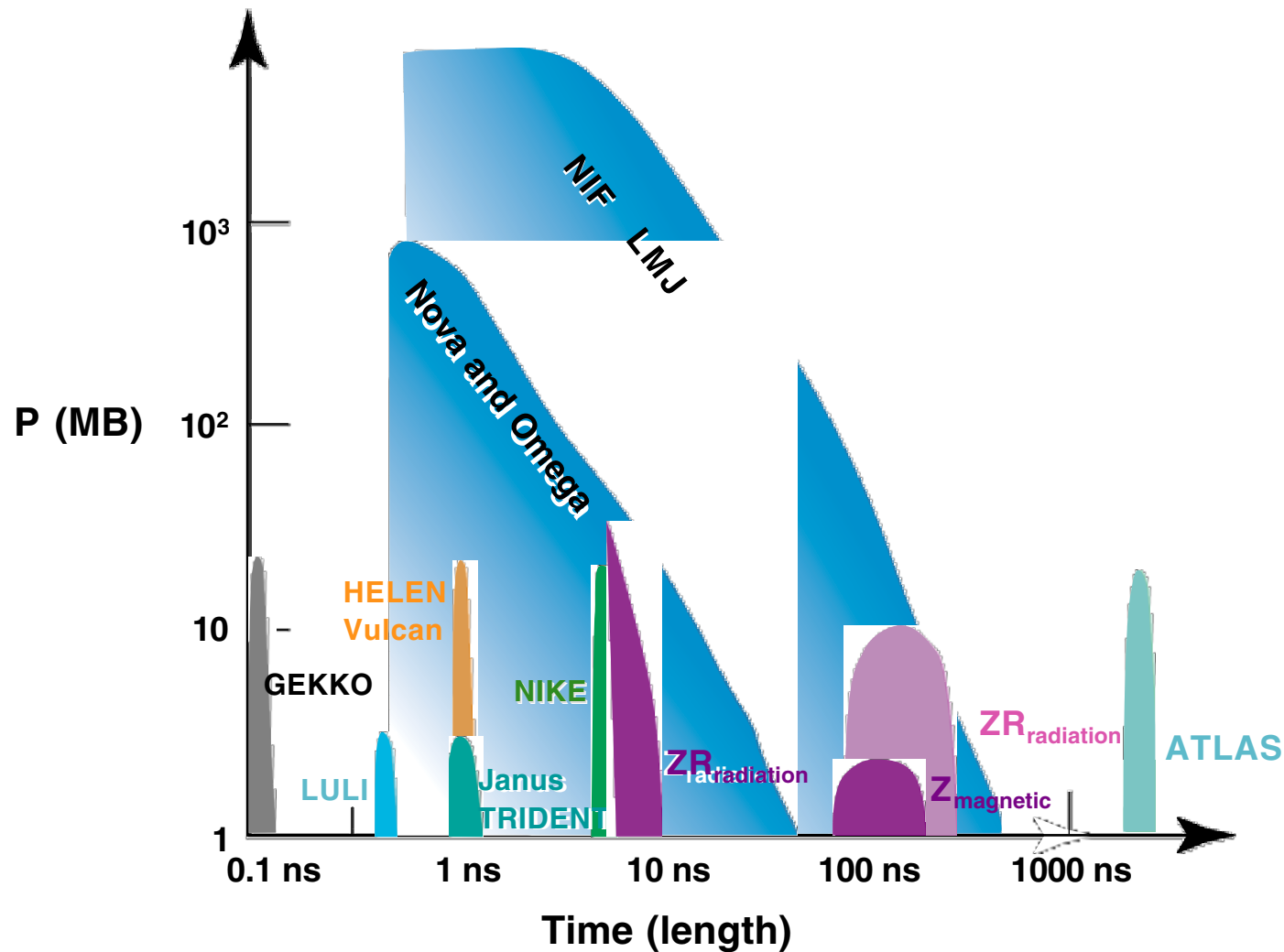
Future

# New facilities will expand high energy density research



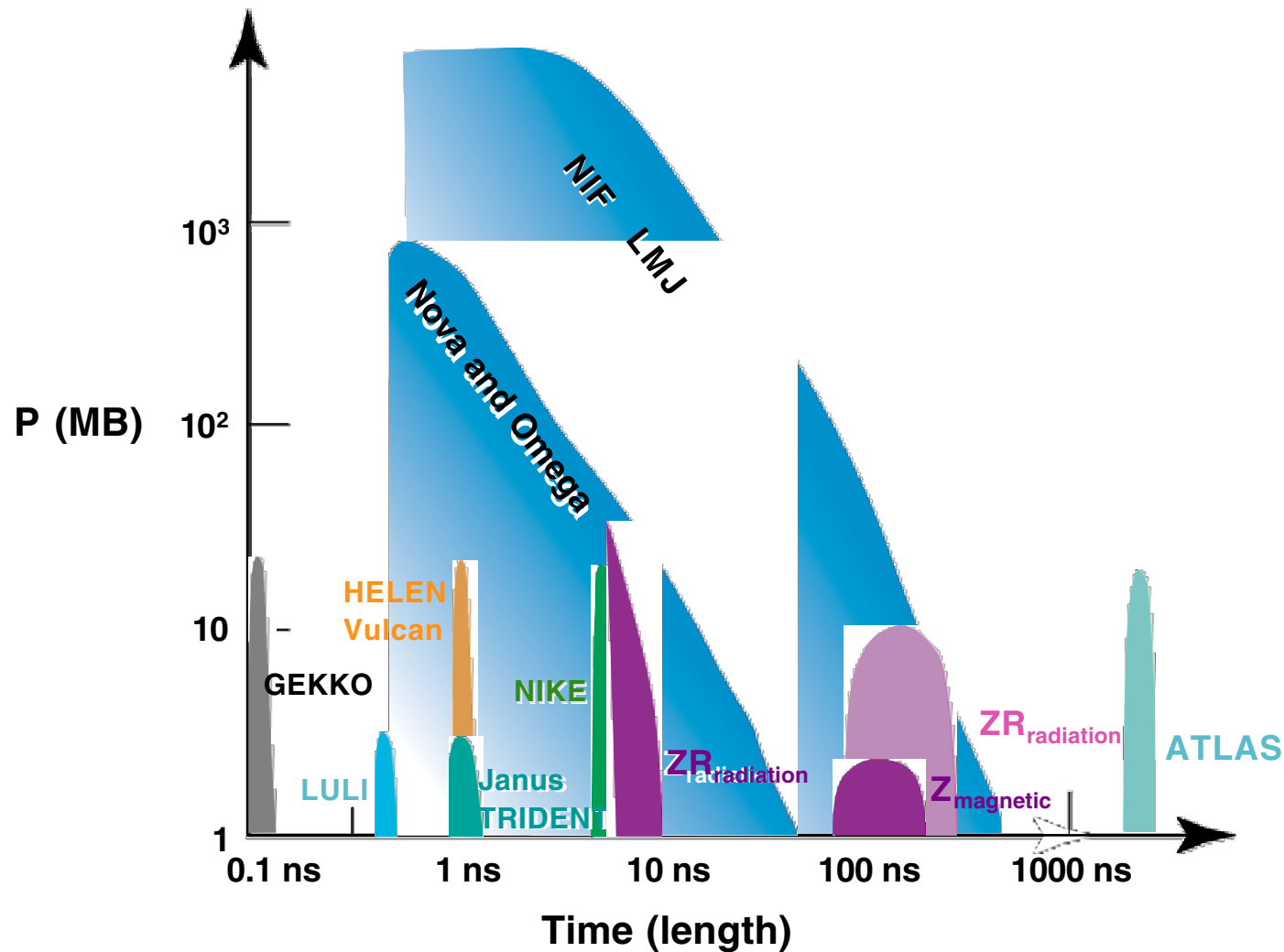
Future

# New facilities will expand high energy density research



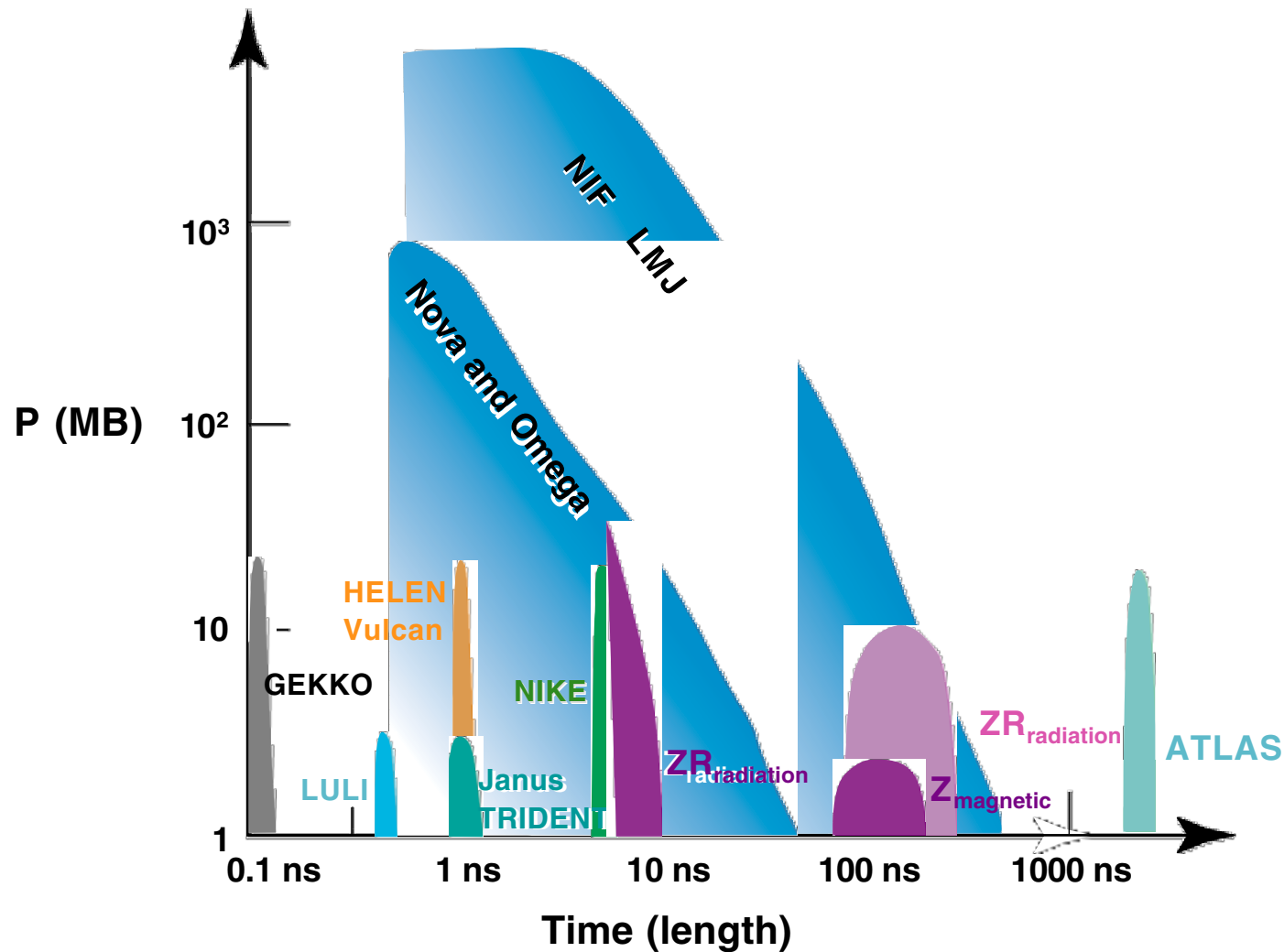
Future

# New facilities will expand high energy density research



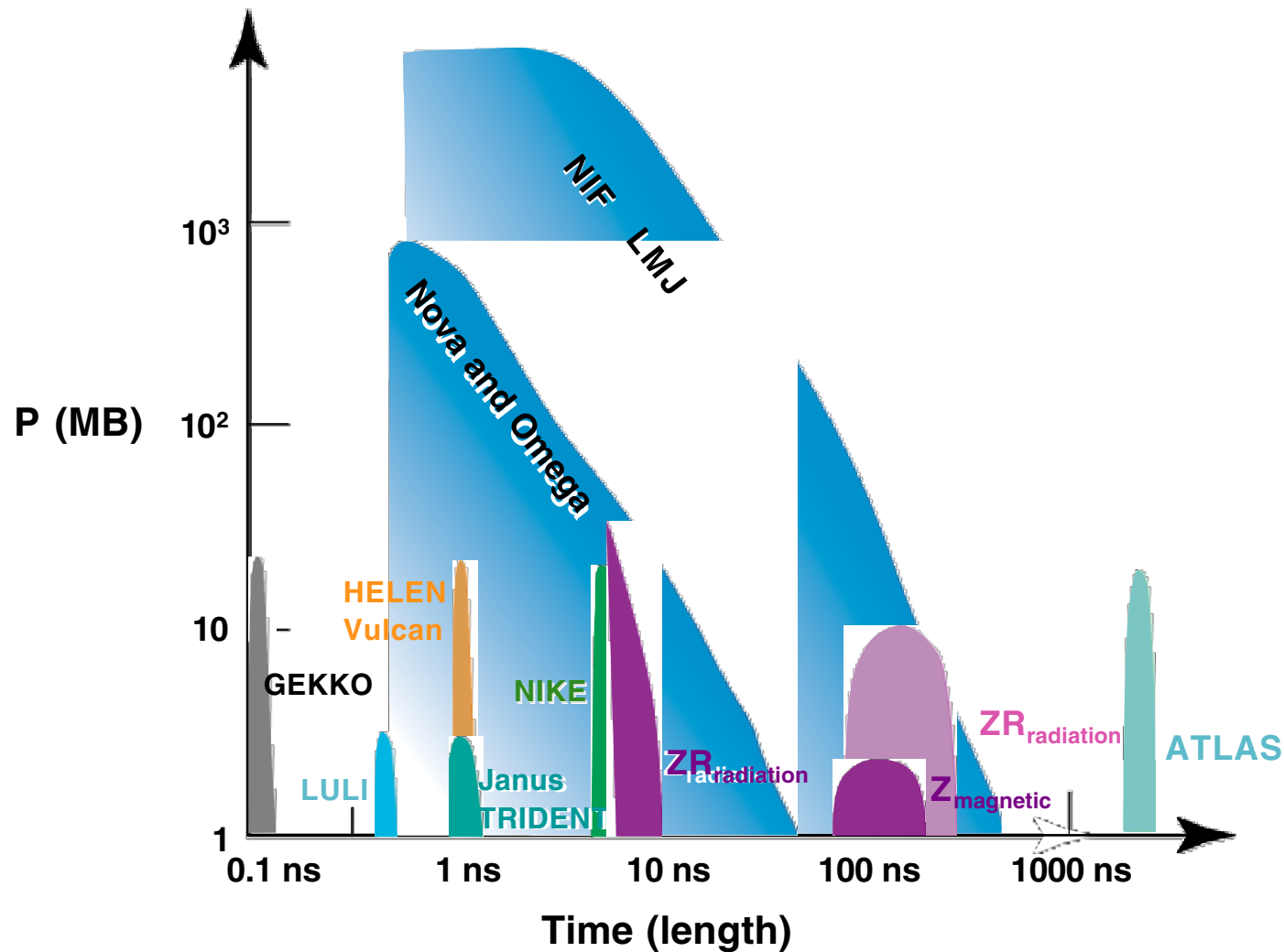
Future

# New facilities will expand high energy density research



Future

# New facilities will expand high energy density research

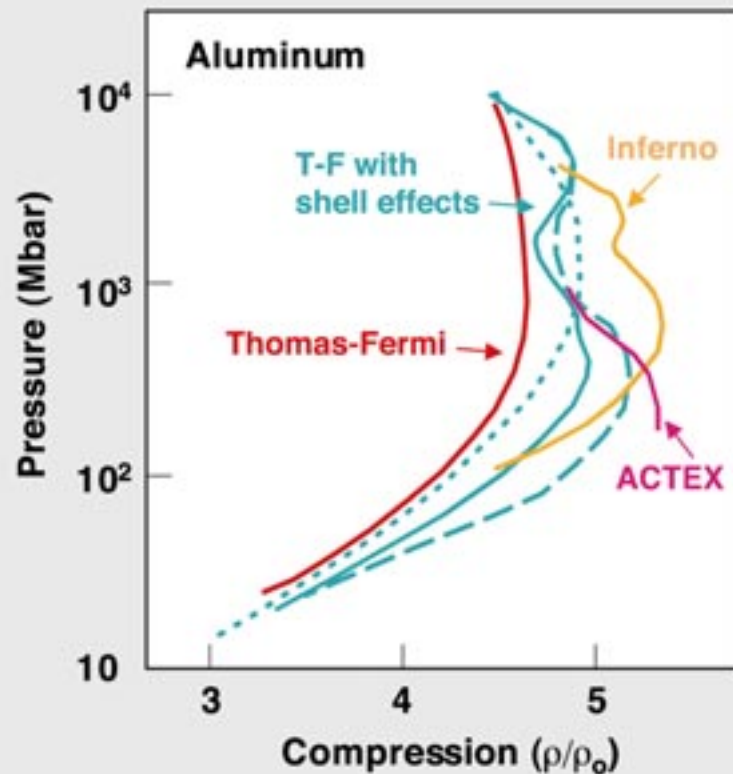


Future

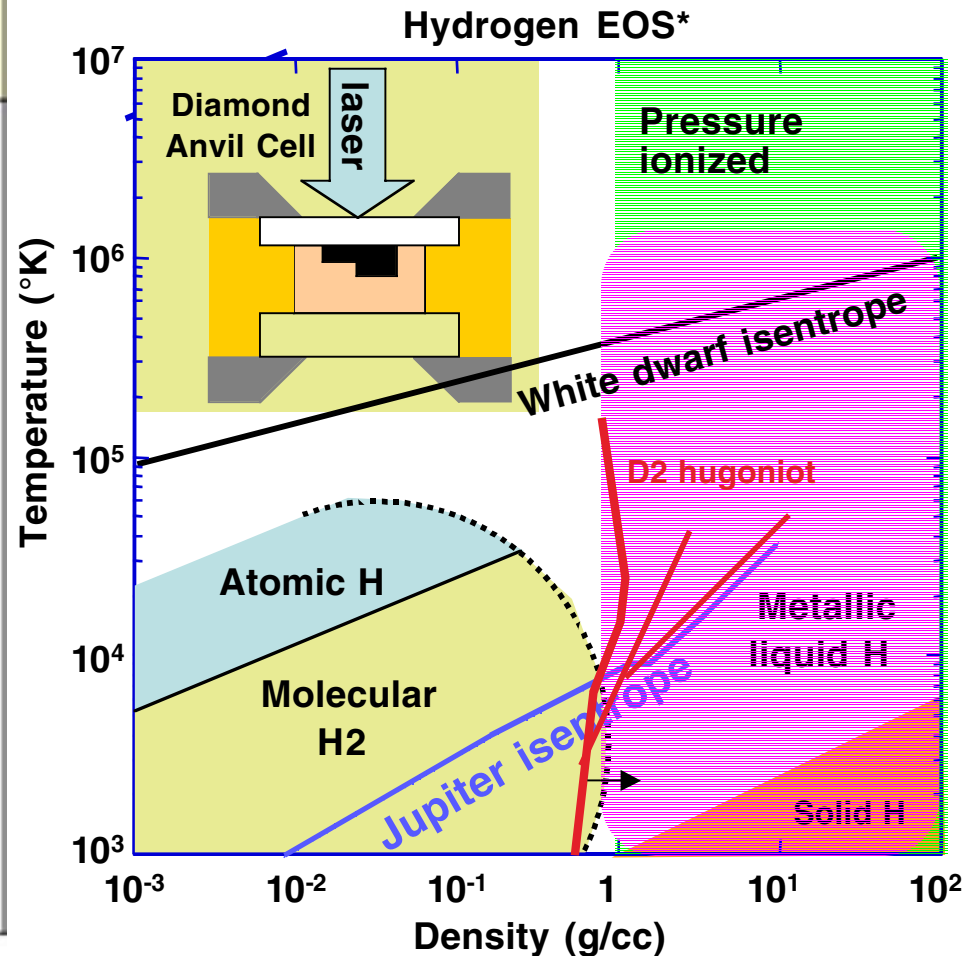
# Access to new regimes of EOS space will allow investigations of new physics



## Theories Disagree at High Pressure



At high ~ Gbar Hugoniot pressures available on NIF, ionization of shells affects predictions



Multishocks, precompression or isentropic compression will access new regimes

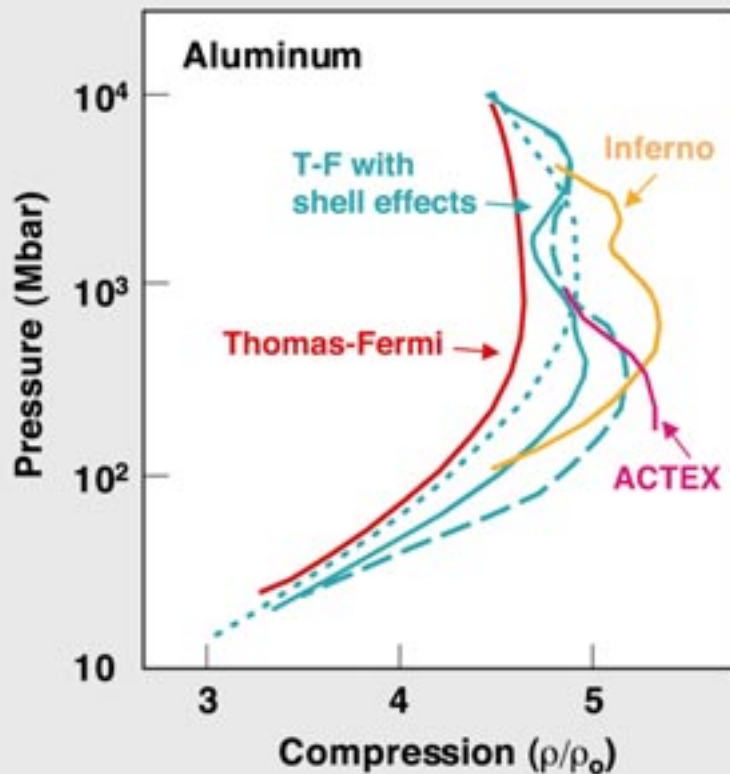


Future

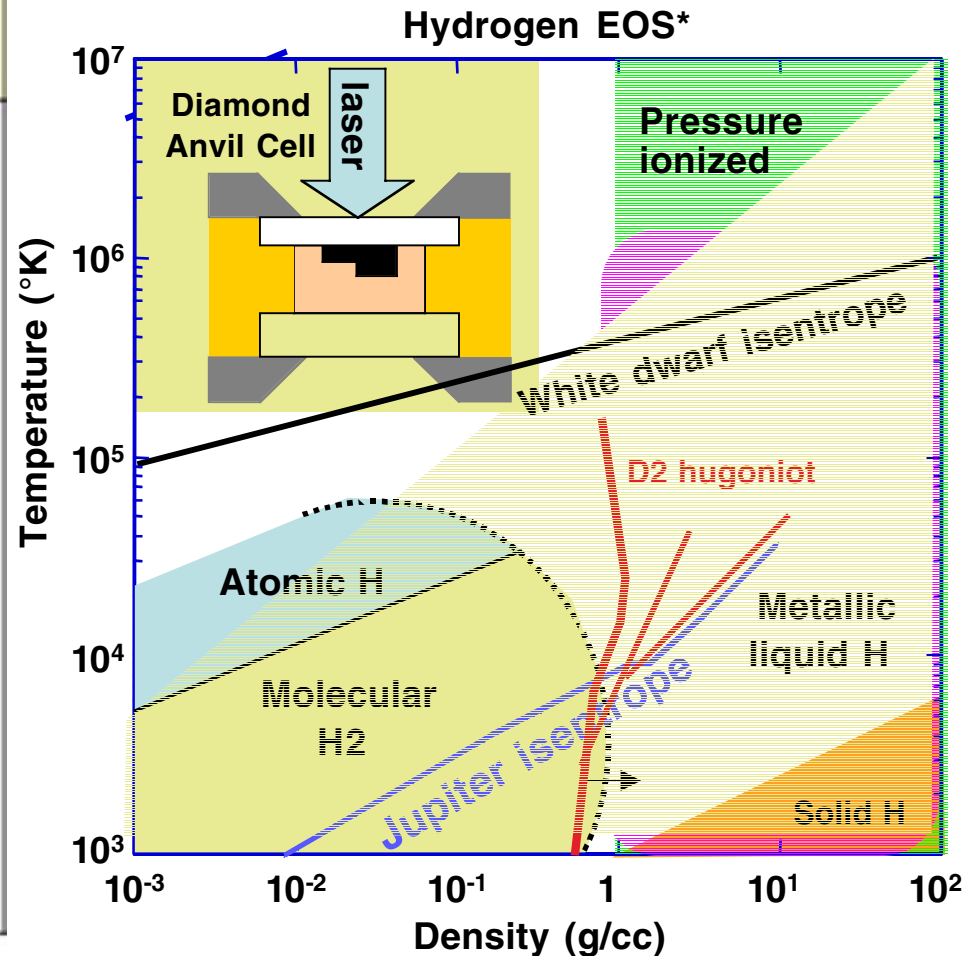
# Access to new regimes of EOS space will allow investigations of new physics



## Theories Disagree at High Pressure



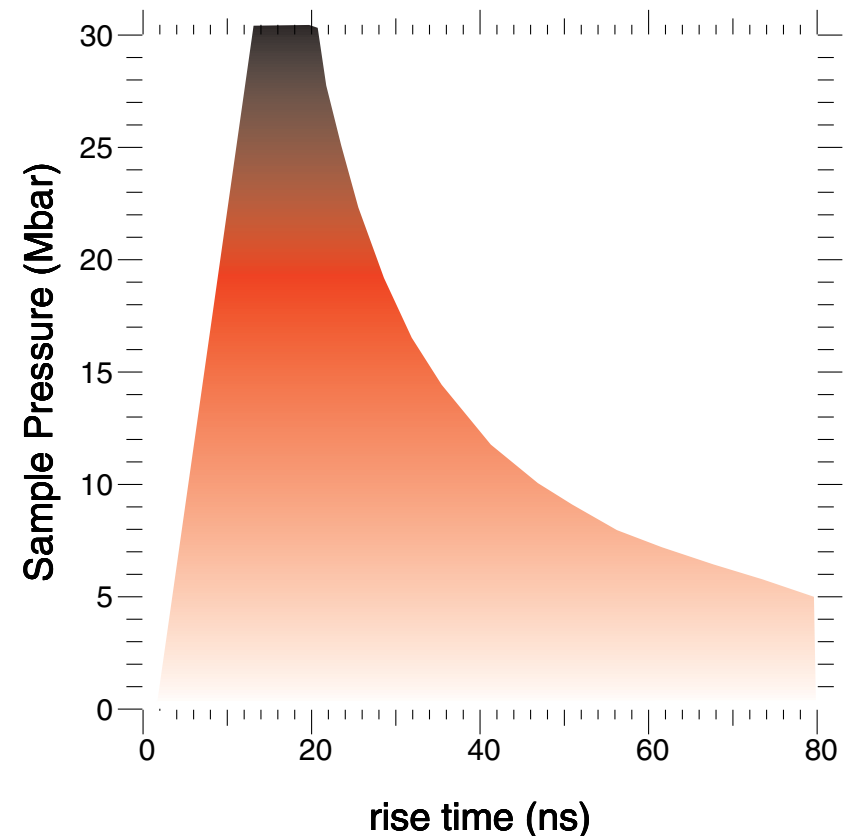
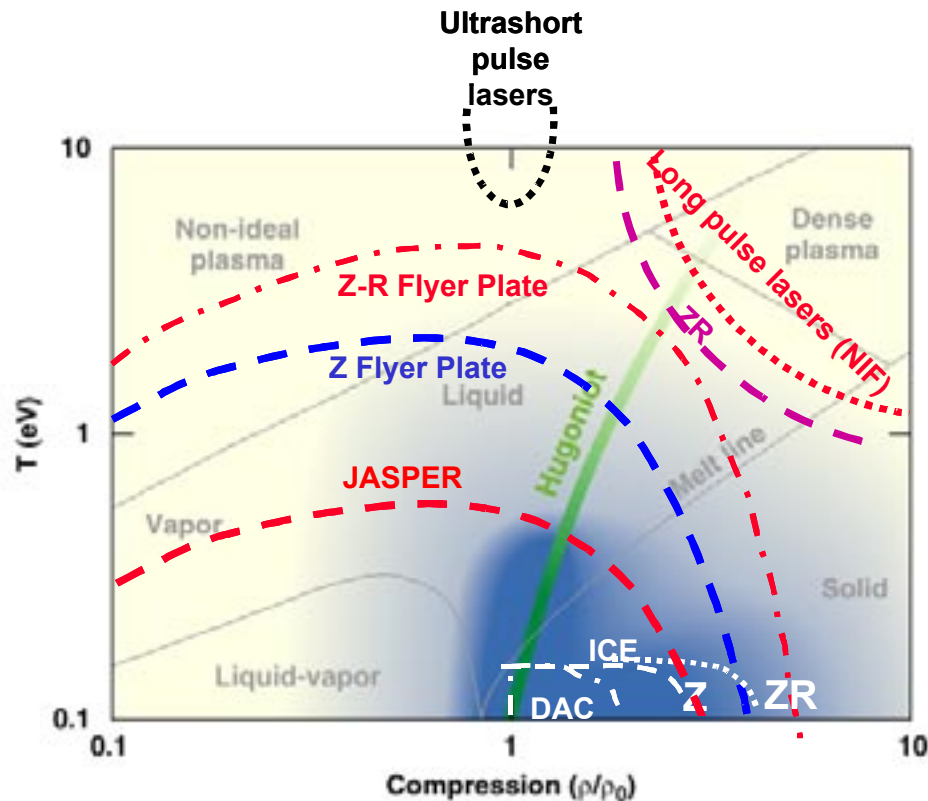
At high ~ Gbar Hugoniot pressures available on NIF, ionization of shells affects predictions



Multishocks, precompression or isentropic compression will access new regimes

Future

Very high pressures will be able to be accessed isentropically for materials science studies



- Isentropic pressures exceeding 10 Mbar for condensed matter studies on ZR
- Flyer plate impact pressures of several tens of Mbar for precise Hugoniot experiments

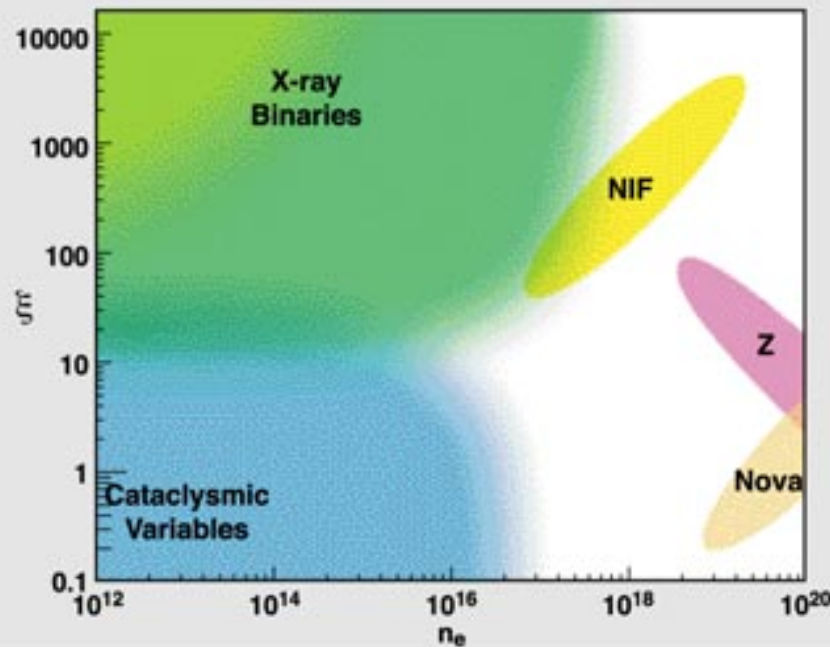
NIF will reach > 20 MBars and measure strain rate and grain size effects

Future

# New regimes will allow measurements closer to astrophysical conditions

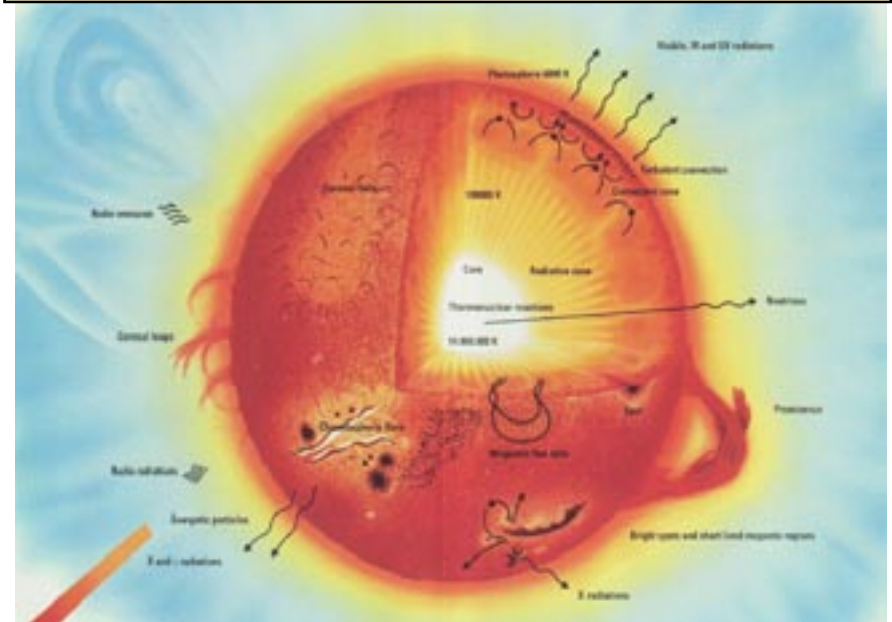


## Photo-ionized Nebulae



NIF & LMJ extends experiments closer to X-ray binary conditions

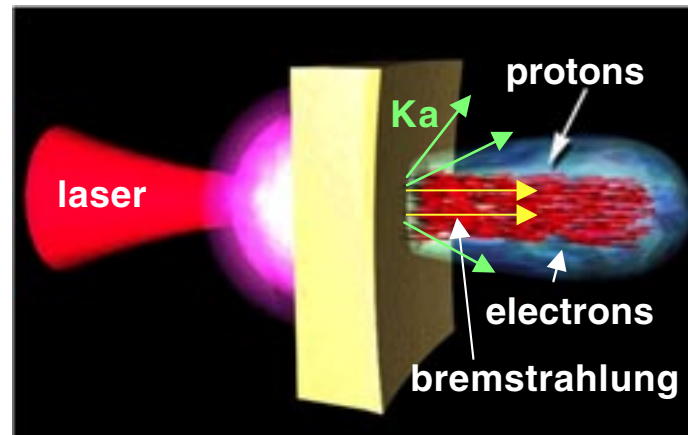
## Fe opacity impacts Solar models



NIF and LMJ extends range of temperatures for opacity measurements to  $> 300$  eV



# New class of petawatt lasers have potential for accessing and probing high energy density conditions



Hot electrons,  
protons, Ka, MeV  
bremstrahlung are  
generated

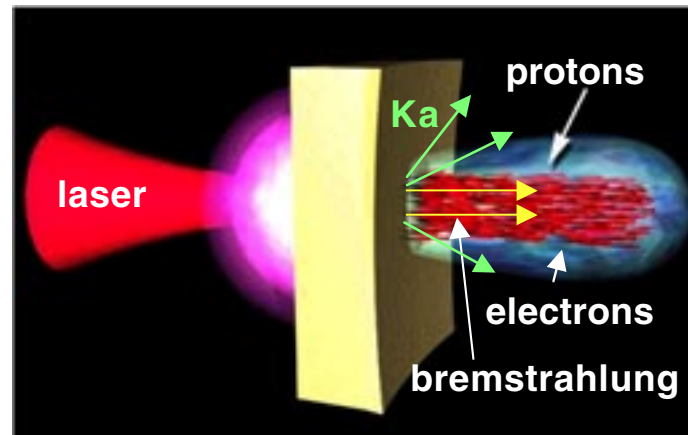
- High intensity electric and magnetic fields are generated

$$\frac{\epsilon E^2}{2} = 1 \text{ Mbar}$$

$$E \sim 10^{11} \text{ W/cm}^2 \sim e/r^2$$

electric field in Bohr atom

# New class of petawatt lasers have potential for accessing and probing high energy density conditions



Hot electrons,  
protons, Ka, MeV  
bremstrahlung are  
generated

- High intensity electric and magnetic fields are generated

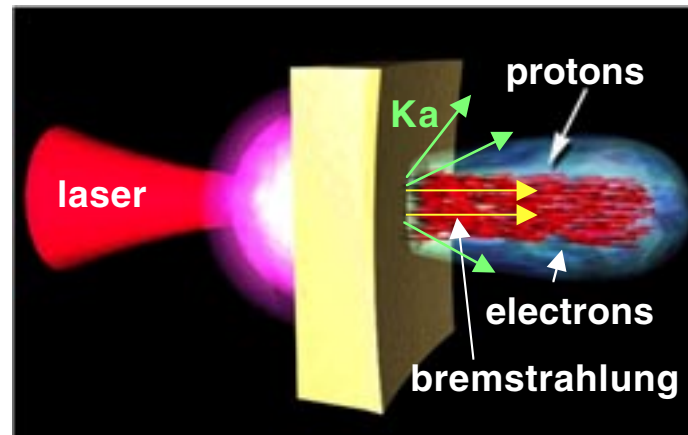
$$\frac{\epsilon E^2}{2} = 1 \text{ Mbar}$$

$$E \sim 10^{11} \text{ W/cm}^2 \sim e/r^2$$

electric field in Bohr atom

$I \sim 3 \times 10^{15} \text{ W/cm}^2 \longrightarrow$  Hot electrons  $\longrightarrow$  Ka X-rays

# New class of petawatt lasers have potential for accessing and probing high energy density conditions



Hot electrons,  
protons, Ka, MeV  
bremstrahlung are  
generated

- High intensity electric and magnetic fields are generated

$$\frac{\epsilon E^2}{2} = 1 \text{ Mbar}$$

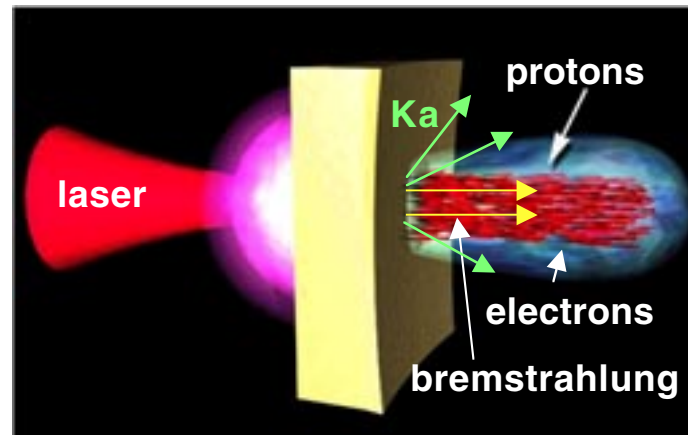
$$E \sim 10^{11} \text{ W/cm}^2 \sim e/r^2$$

electric field in Bohr atom

$$I \sim 3 \times 10^{15} \text{ W/cm}^2 \longrightarrow \text{Hot electrons} \longrightarrow \text{Ka X-rays}$$

$$I \sim 10^{18} \text{ W/cm}^2 \longrightarrow \frac{\text{quiver momentum}}{m_e c} = 1$$

# New class of petawatt lasers have potential for accessing and probing high energy density conditions



Hot electrons,  
protons, Ka, MeV  
bremstrahlung are  
generated

- High intensity electric and magnetic fields are generated

$$\frac{\epsilon E^2}{2} = 1 \text{ Mbar}$$

$$E \sim 10^{11} \text{ W/cm}^2 \sim e/r^2$$

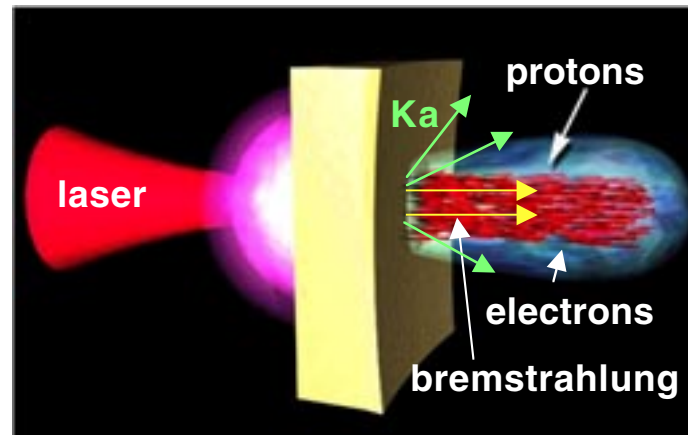
electric field in Bohr atom

$$I \sim 3 \times 10^{15} \text{ W/cm}^2 \longrightarrow \text{Hot electrons} \longrightarrow \text{Ka X-rays}$$

$$I \sim 10^{18} \text{ W/cm}^2 \longrightarrow \frac{\text{quiver momentum}}{m_e c} = 1$$

$$I \sim 10^{19} \text{ W/cm}^2 \longrightarrow \text{Mev Bremstrahlung}$$

# New class of petawatt lasers have potential for accessing and probing high energy density conditions



Hot electrons,  
protons, Ka, MeV  
bremstrahlung are  
generated

- High intensity electric and magnetic fields are generated

$$\frac{\epsilon E^2}{2} = 1 \text{ Mbar}$$

$$E \sim 10^{11} \text{ W/cm}^2 \sim e/r^2$$

electric field in Bohr atom

$$I \sim 3 \times 10^{15} \text{ W/cm}^2 \longrightarrow \text{Hot electrons} \longrightarrow \text{Ka X-rays}$$

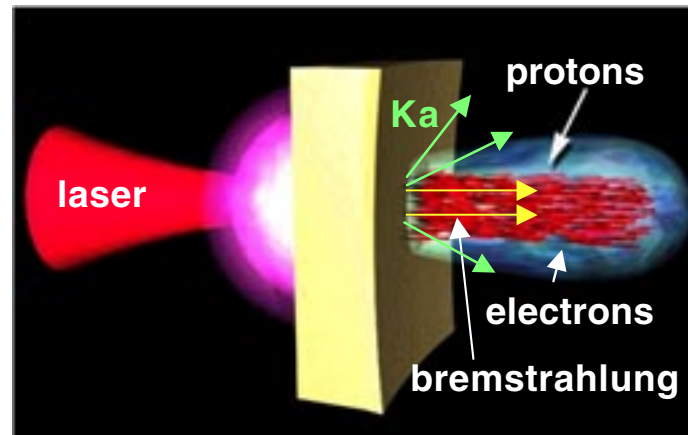
$$I \sim 10^{18} \text{ W/cm}^2 \longrightarrow \frac{\text{quiver momentum}}{m_e c} = 1$$

$$I \sim 10^{19} \text{ W/cm}^2 \longrightarrow \text{Mev Bremstrahlung}$$

$$I \sim 10^{20} \text{ W/cm}^2 \longrightarrow \text{Mev protons, 400 MG, pair production}$$



# New class of petawatt lasers have potential for accessing and probing high energy density conditions



Hot electrons,  
protons, Ka, MeV  
bremstrahlung are  
generated

- High intensity electric and magnetic fields are generated

$$\frac{\epsilon E^2}{2} = 1 \text{ Mbar}$$

$$E \sim 10^{11} \text{ W/cm}^2 \sim e/r^2$$

electric field in Bohr atom

$$I \sim 3 \times 10^{15} \text{ W/cm}^2 \longrightarrow \text{Hot electrons} \longrightarrow \text{Ka X-rays}$$

$$I \sim 10^{18} \text{ W/cm}^2 \longrightarrow \frac{\text{quiver momentum}}{m_e c} = 1$$

$$I \sim 10^{19} \text{ W/cm}^2 \longrightarrow \text{Mev Bremstrahlung}$$

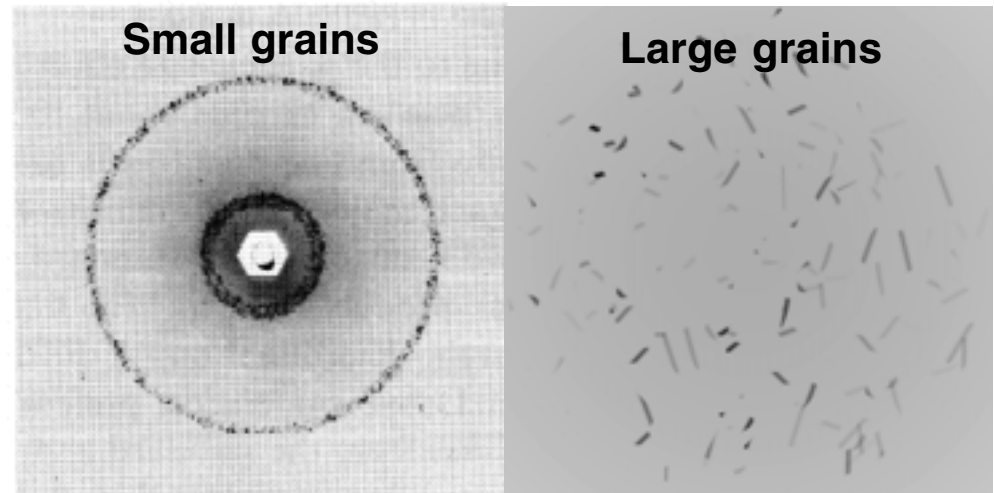
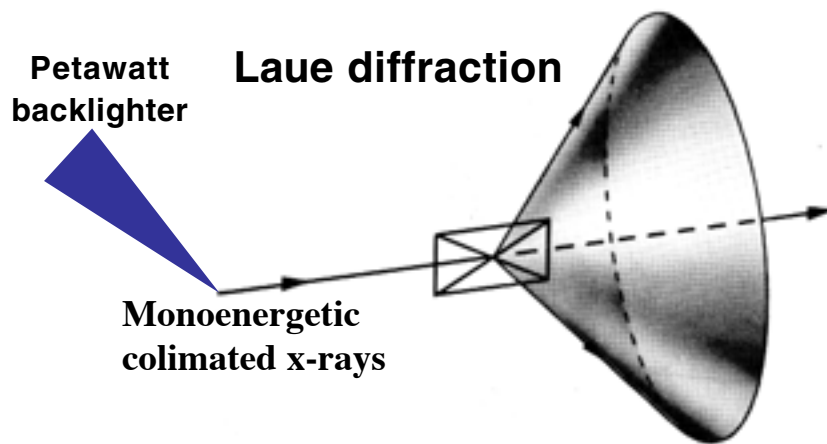
$$I \sim 10^{20} \text{ W/cm}^2 \longrightarrow \text{Mev protons, 400 MG, pair production}$$

$$I \sim 10^{24} \text{ W/cm}^2 \longrightarrow \text{Nuclear reactions}$$

# High photon energy and proton probing of high energy density conditions is possible



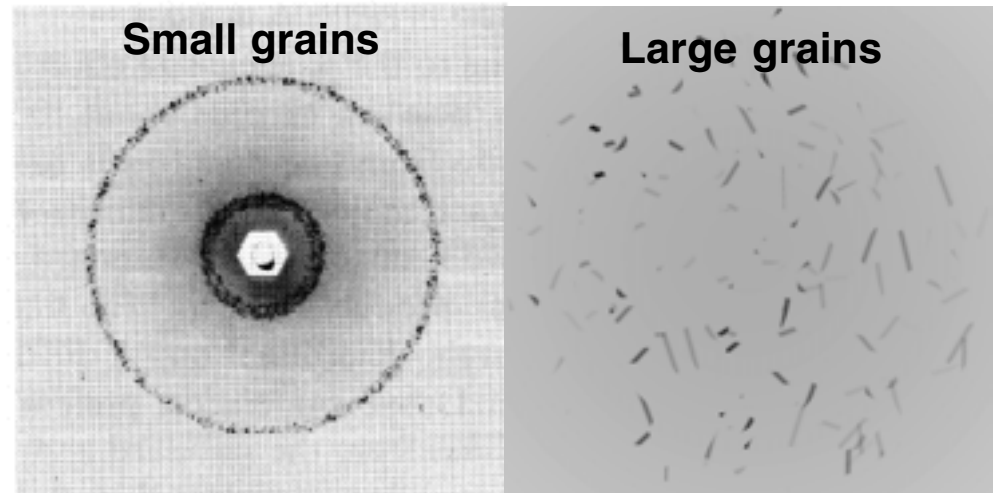
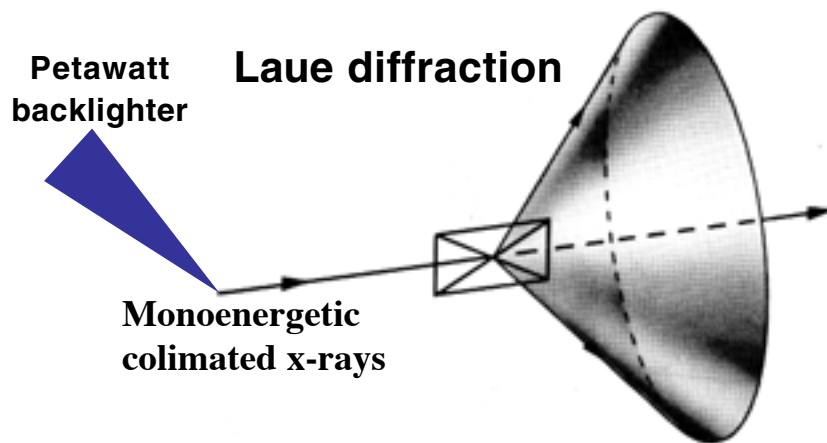
- 20-100 keV Ka can be used to measure grain size and phase of solid material at high compression



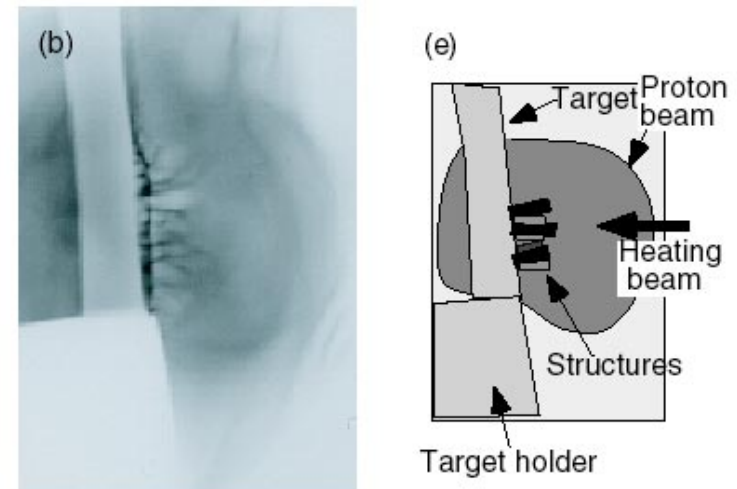
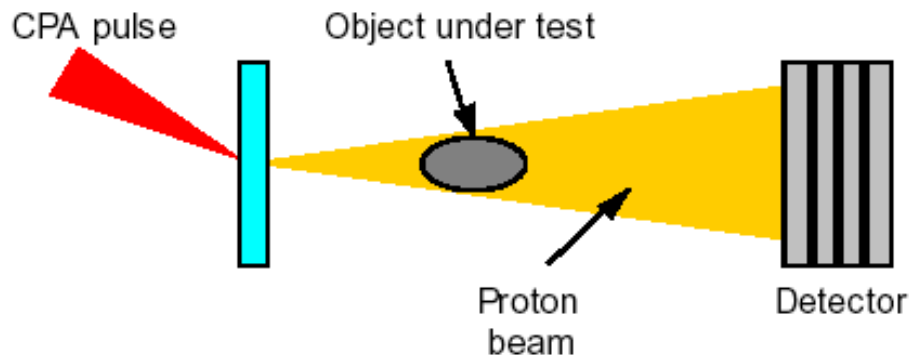
# High photon energy and proton probing of high energy density conditions is possible



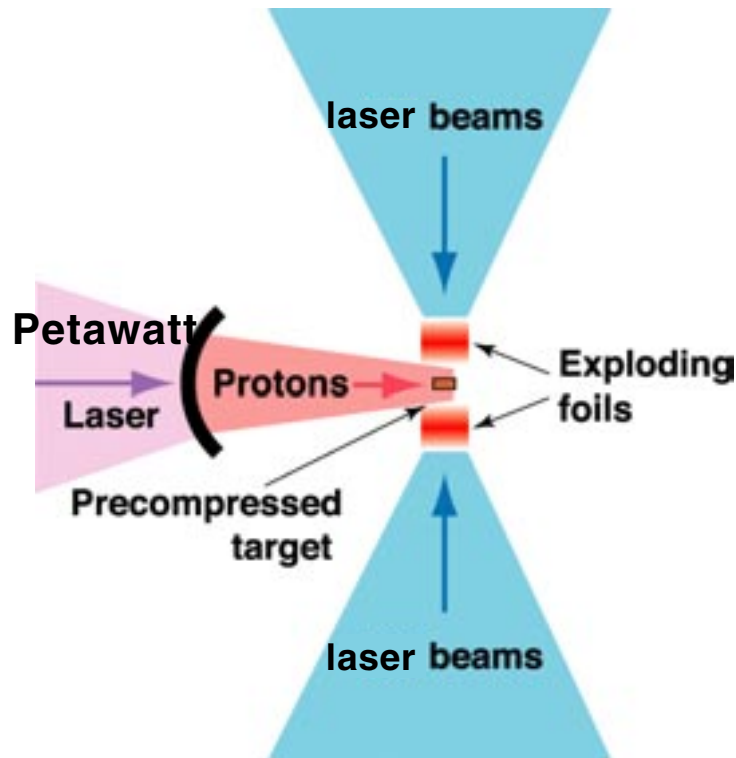
- 20-100 keV Ka can be used to measure grain size and phase of solid material at high compression



## •Proton imaging:

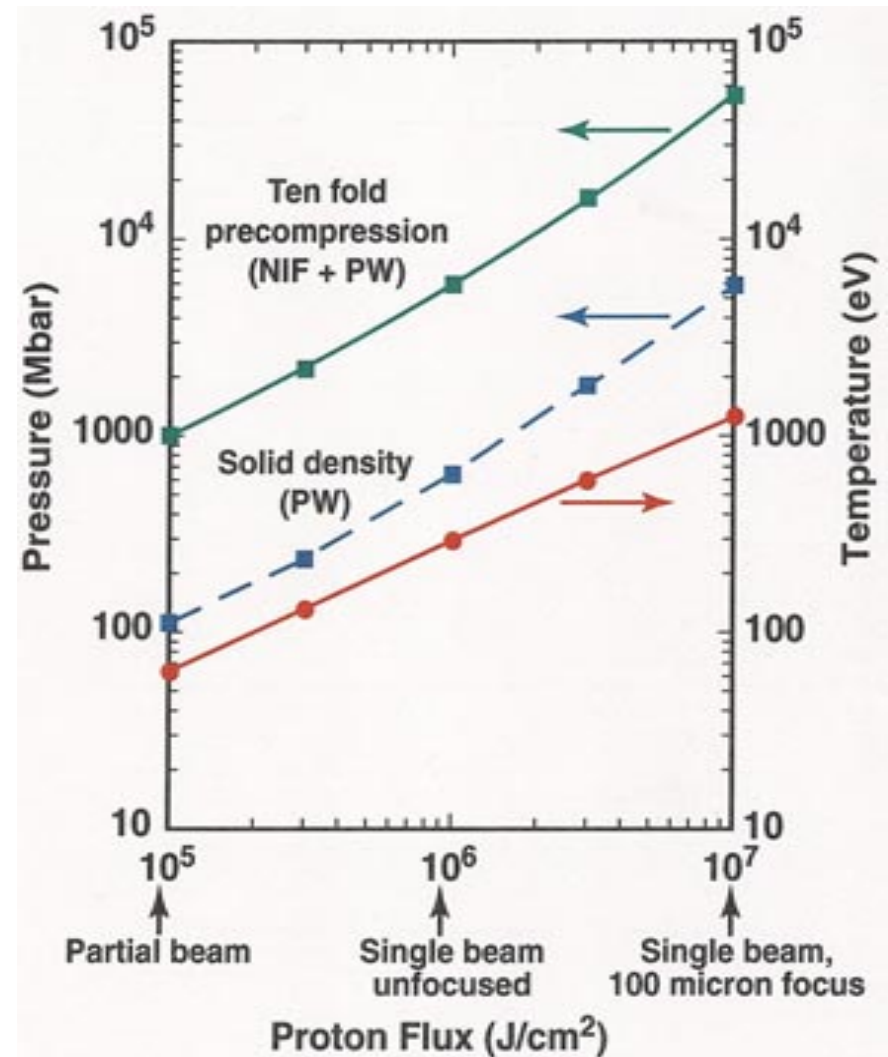


# The capability to independently compress and heat can expand the regimes for EOS and opacity measurements



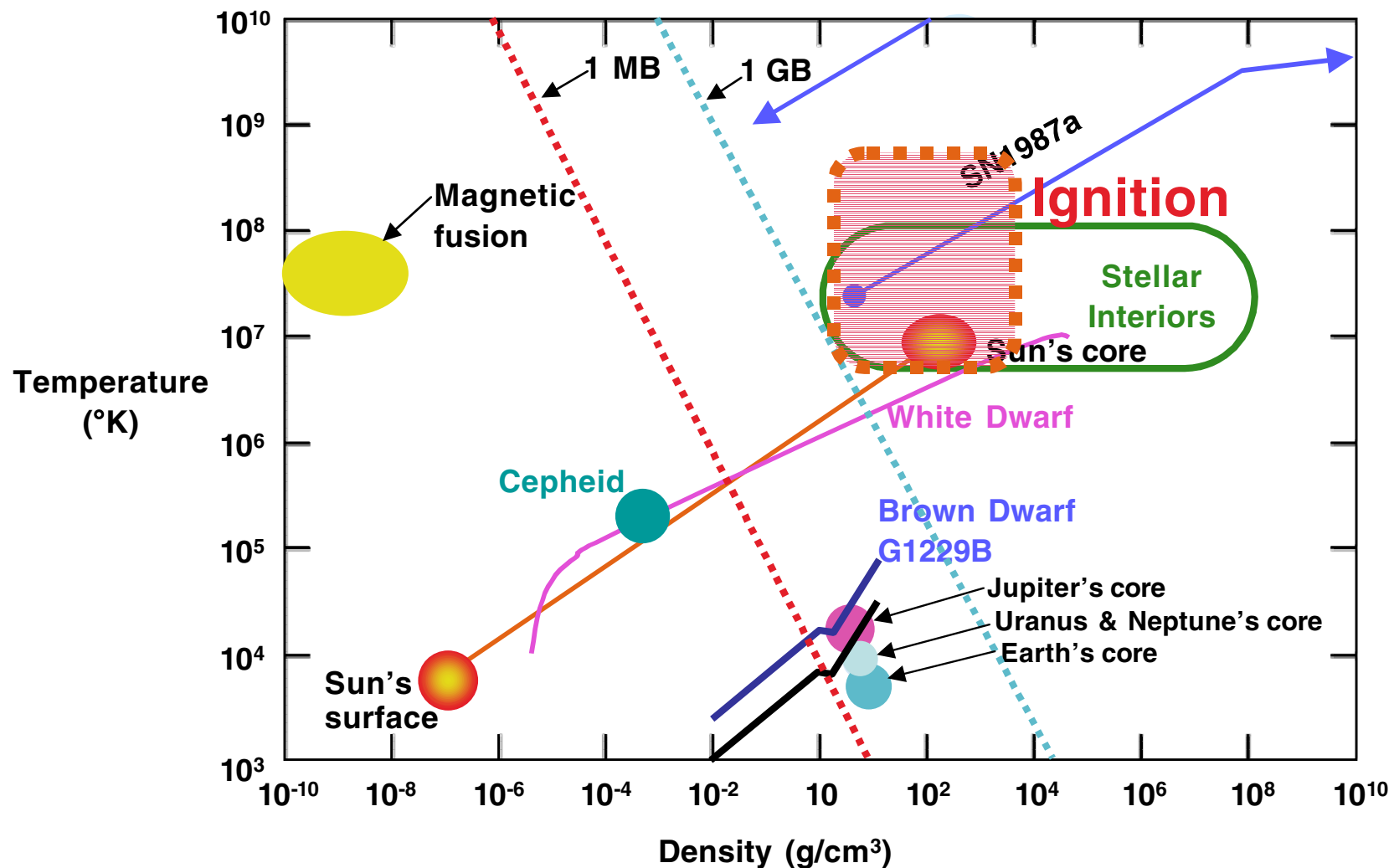
Long pulse lasers are used to precompress a material

Ion and electron beams can potentially heat materials uniformly to high energy density



Future

In burning capsules, thermonuclear reaction rates in stars may be studied



# Summary

---



- **Significant advances in high energy density physics have occurred over the last six years**
- **The ability to make precise measurements in new regimes allows comparion with models**
  - *Hugoniot equation-of-state*
  - *Materials science at high pressure*
  - *Hydrodynamics*
  - *Radiation transport*
- **New facilities will expand access to high energy density regimes**

## **Recommendation from NRC report**

### **Connecting Quarks with the Cosmos: Eleven Science Questions for the New Century**

**“Discern the physical principles that govern extreme astrophysical environments through the laboratory study of high energy density physics... The field is in its infancy...”**